

## Dynamic design of high-speed railway bridges

Simplifications and guidelines
Master's Thesis in the International Master's Programme Structural Engineering and Building Performance Design

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CHALMERS UNIVERSITY OF TECHNOLOGY
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Design curves for a two-span bridge in 2D, Figure 4.51.
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#### Abstract

In 2006 the bridge code was updated to include a demand of dynamic analysis of railway bridges subjected to high-speed trains. This analysis includes several complex loads to be examined for a large range of train velocities and is therefore complicated and time consuming. The dynamic analysis' complexity also makes it hard to establish guidelines for the early bridge design even with increased experience. There is hence a need today to introduce both guidelines for the early bridge design and simplifications for the advanced dynamic analysis of railway bridges. This master thesis provides an extensive study of railway bridges' dynamic behavior and presents


 a method of guidance for the dynamic design.As a continuation of previous master theses, carried out at Reinertsen Sverige AB, a method for transforming a railway bridge into a single-degree-of-freedom (SDOF) system is examined in this thesis. It is shown that a railway bridge of several spans can be successfully transformed into a SDOF system that provides accurate results compared to a finite element analysis.

A parameter study is performed in this thesis with the purpose of studying the individual bridge parameters' effect on the dynamic response. It is shown that material parameters exert an describable effect, while the effect of geometric parameters is more complicated. The increased understandings concerning bridge parameters' effect on the dynamic response and resonance phenomenon in railway bridges are used in the creation of a graphical calculation tool, referred to as a design curve. A design curve is a comprehensible presentation of an extensive number of finite element analyses from which it possible to obtain the designing vertical acceleration, independent of the choice of train velocity, material parameters and cross-section geometry.
A comparison between 2D and 3D analyses of railway bridges is performed in this thesis. It is shown how the designing acceleration is altered as 3D geometry and hence bridge torsion and load eccentricity is considered. A discussion is presented on the possibility of creating guidelines and simplifications of railway bridges' dynamic design with consideration to 3D effects.

Keywords: Dynamics, Railway bridges, high-speed trains, HSLM, SDOF, eccentric loading, resonance, vibrations

Dynamisk design av järnvägsbroar utsatta för höghastighetståg
Förenklingar och riktlinjer
Examensarbete inom det Internationella mastersprogrammet Structural Engineering and Building Performance Design
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## SAMMANFATTNING

År 2006 uppdaterades bronormen till att inkludera krav på den dynamiska analysen av järnvägsbroar som är utsatta för höghastighetståg. En sådan analys inkluderar flera komplexa lastfall som ska undersökas för ett stort spann av hastigheter, vilket gör analysen komplicerad och tidskrävande. Komplexiteten i den dynamiska analysen gör det svårt att upprätta riktlinjer för brokonstruktionen i ett tidigt skede, även med ökad erfarenhet i ämnet. Idag finns det därför ett behov av riktlinjer för den tidiga designen och förenklingar av den avancerade dynamiska analysen av järnvägsbroar. Detta examensarbete består av en omfattande studie av järnvägsbroars dynamiska beteende och presenterar en metod med riktlinjer för dynamisk design.

Som fortsättning på tidigare examensarbeten, utförda på Reinertsen Sverige AB, har en metod för att omvandla en järnvägsbro till ett enfrihetsgradssystem (SDOF) undersöks. Det visas att en järnvägsbro i flera span framgångsrikt kan förvandlas till ett SDOF system som ger tillförlitliga resultat jämfört med finita element analyser.

En parameterstudie har genomförts med syfte att studera individuella broparametrars inverkan på den dynamiska responsen. Det visas att materialparametrar har beskrivbar inverkan medan effekten av de geometriska parametrarna är mer komplicerad. Den ökade förståelsen för hur olika broparametrar påverkar den dynamiska responsen och resonansfenomen i järnvägsbroar används för att skapa ett grafiskt beräkningsverktyg som kallas designkurvor. En designkurva är en jämförbar presentation av ett stort antal finita element analyser från vilka det är möjligt att erhålla den vertikala dimensionerande accelerationen oberoende av val av tåghastighet, materialparametrar och sektionens geometri.

En jämförelse mellan 2D och 3D analyser av järnvägsbroar utförs i detta examensarbete. Det visas hur den dimensionerande accelerationen förändras med olika geometrier, varför vridning och lastexcentricitet har undersökts. Slutligen presenteras en diskussion kring möjligheten att skapa riktlinjer och förenklingar för den dynamiska dimensioneringen av järnvägsbroar med hänsyn till 3D-effekter.

Nyckelord: Dynamik, Järnvägsbroar, höghastighetståg, HSLM, enfrihetsgradssystem, excentrisklast, resonans, vibrationer

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## Preface

The dynamic design of high-speed railway bridges have been studied in this master thesis with the aim of producing guidelines and simplifications in the design process. This thesis has been performed at Reinertsen Sverige AB in co-operation with the Department of Civil and Environmental Engineering at Chalmers University of Technology.

Mario Plos, PhD, at Concrete Structures at Chalmers University of Technology was the examiner and Ginko Gueorguiev, MSc, at Reinertsen Sverige AB, was the supervisor of this thesis.

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Patrik Eriksson and Emanuel Trolin

## Notations

Roman upper case letters

| $A$ | Cross-section area $\left[\mathrm{m}^{2}\right]$ |
| :--- | :--- |
| 2D | Two-dimensional |
| 3D | Three-dimensional |
| $C$ | Viscosity of the dashpot [Ns $/ \mathrm{m}]$ |
| $\boldsymbol{C}$ | Damping matrix [Ns $/ \mathrm{m}]$ |
| $\widehat{\boldsymbol{C}}$ | Modal damping matrix |
| $\widehat{\boldsymbol{C}}_{\text {red }}$ | Reduced modal damping matrix |
| $D(r)$ | Dynamic magnification factor [-] |

D Coach length [m]

D Dynamical matrix
DOF Degrees of freedom
$E \quad$ Young's Modulus $\left[\mathrm{N} / \mathrm{m}^{2}\right]$
FE Finite Element
FEM Finite Element Method
HSLM High-speed load model
$I \quad$ Moment of inertia $\left[\mathrm{m}^{4}\right]$
$K \quad$ Stiffness of the SDOF model spring [ $\mathrm{N} / \mathrm{m}$ ]
$\boldsymbol{K} \quad$ Stiffness matrix $[\mathrm{N} / \mathrm{m}]$
$\widehat{\boldsymbol{K}} \quad$ Modal stiffness matrix
$\widehat{\boldsymbol{K}}_{\text {red }} \quad$ Reduced modal stiffness matrix
$L \quad$ Beam/bridge length [m]
$L_{\text {tot }} \quad$ Beam/bridge total length [m]
$M \quad$ Mass [kg], mass of the SDOF model $[\mathrm{kg}]$
$M \quad$ Mass matrix [kg]
$\widehat{\boldsymbol{M}} \quad$ Modal mass matrix
$\widehat{\boldsymbol{M}}_{\text {red }} \quad$ Reduced modal mass matrix
$\mathrm{N} \quad$ Number of intermediate coaches [-]
P Point force [kN]
$\mathrm{R}^{2} \quad$ Least square values [-]
SDOF Single Degree Of Freedom
U Amplitude [m]
$\boldsymbol{U} \quad$ Amplitude vector [m]
$U_{0} \quad$ Static amplitude [m]
$U(r) \quad$ Dynamic amplitude [m]

## Roman lower case letters

| $a$ | Design acceleration |
| :--- | :--- |
| b | Width of bridge $[\mathrm{m}]$ |
| $c_{c r}$ | Critical damping coefficient |
| $c_{i}$ | Viscosity of the dashpot $[\mathrm{Ns} / \mathrm{m}]$ |
| d | $[\mathrm{m}]$ |
| e | Eccentricity $[\mathrm{m}]$ |
| $f_{i}$ | Forces $[\mathrm{N}]$ |
| $f_{i}$ | Eigenfrequency $[\mathrm{Hz}]$ |
| $f_{i, \text { single }}$ | Eigenfrequency for single-span bridge $[\mathrm{Hz}]$ |
| $g_{i}(\mu)$ | Eigenfrequency function |


| h | Height [m] |
| :---: | :---: |
| 1 | Length of beam/bridge [m] |
| $m_{i}$ | Mass [kg] |
| $n$ | Number of eigenvalue |
| $p(t)$ | Time dependent load [ N ] |
| $p_{y}(x, t)$ | Load per unit length in location $x$ at time $t[\mathrm{~N} / \mathrm{m}]$ |
| $\boldsymbol{p}(t)$ | Time dependent load vector [ N ] |
| $p_{0}$ | [N] |
| $r$ | Frequency ratio [-] |
| $u$ | Displacement/deflection [m] |
| $p(t)_{\text {sdof }}$ | Load acting on SDOF model [ N ] |
| $\boldsymbol{u}$ | Displacement vector [m] |
| $u_{s}$ | Eigenvector for iteration s |
| $\dot{u}$ | First derivative of $u$ with respect to time t , velocity [ $\mathrm{m} / \mathrm{s}$ ] |
| $\dot{\boldsymbol{u}}$ | First derivative of $\boldsymbol{u}$, velocity vector [ $\mathrm{m} / \mathrm{s}$ ] |
| $\ddot{u}$ | Second derivative of $u$ with respect to time t , acceleration $\left[\mathrm{m} / \mathrm{s}^{2}\right]$ |
| $\boldsymbol{u}$ | Second derivative of $\boldsymbol{u}$, acceleration vector $\left[\mathrm{m} / \mathrm{s}^{2}\right]$ |
| $u(t)$ | Total displacement in time [m] |
| $u_{c}(t)$ | Homogenous solution of the general equation of motion [m] |
| $u_{p}(t)$ | Particular displacement solution of the general equation of motion [m] |
| $u(t)_{\text {sdof }}$ | Displacement of SDOF system [m] |
| $u(t)_{\text {bridge }}$ | Critical displacement in a bridge [m] |
| $\dot{u}_{p}$ | First derivative of $u_{p}$ with respect to time t , velocity [ $\mathrm{m} / \mathrm{s}$ ] |
| $\ddot{u}_{p}$ | Second derivative of $u_{p}$ with respect to time t , acceleration $\left[\mathrm{m} / \mathrm{s}^{2}\right]$ |
| x | Space coordinate |
| $\bar{s}$ | Root to the equation of motion for a SDOF system |
| $v$ | Velocity [m/s], Poisson's ratio [-] |
| $v\left(x_{e}, t\right)$ | Displacements of a beam |
| $v_{s}$ | Un-scaled eigenfrequency for iteration s |
| $v_{\text {resonance }}$ | Train velocity at which resonance effects occur [ $\mathrm{m} / \mathrm{s}$ ] |
| vcr | Critical velocity [ $\mathrm{m} / \mathrm{s}$ ] |
| $\mathrm{v}_{\mathrm{cr}}$ | Critical velocity [ $\mathrm{m} / \mathrm{s}$ ] |
| Greek upper case letters |  |
| $\Omega$ | Load frequency [Hz] |
| $\Delta t$ | Time step [s] |
| $\boldsymbol{\Phi}$ | Modal matrix [-] |
| $\boldsymbol{\Phi}_{\text {red }}$ | Reduced modal matrix [-] |
| Greek lower case letters |  |
| $\alpha$ | Phase angle [-] |
| $\beta$ |  |
| $\beta^{*}$ |  |
| $\eta$ | Describes the span relation of a three-span bridge |
| $\eta_{i}$ | Modal coordinate of mode i |
| $\boldsymbol{\eta}$ | Modal coordinates |
| $\boldsymbol{\eta}_{\text {red }}$ | Reduced modal coordinates |
| $\dot{\boldsymbol{\eta}}$ | Modal velocity vector |


| $\ddot{\boldsymbol{\eta}}$ | Modal acceleration vector |
| :--- | :--- |
| $\kappa$ | Describes the geometry of a three-span bridge |
| $\lambda_{\mathrm{i}}$ | Eigenvalue |
| $\zeta$ | Viscous damping factor [\%] |
| $\mu$ | Span ratio for a two span bridge |
| $\pi$ | Mathematical constant $3.14159 \ldots$ |
| $\rho$ | Density of material $\left[\mathrm{kg} / \mathrm{m}^{3}\right]$ |
| $\tau$ |  |
| $\omega_{\mathrm{n}}$ | Undamped natural frequency $[\mathrm{rad} / \mathrm{s}]$ |
| $\omega_{\mathrm{i}}$ | Natural frequency or Eigenfrequency $[\mathrm{rad} / \mathrm{s}]$ |
| $\omega_{\mathrm{d}}$ | Damped natural frequency $[\mathrm{rad} / \mathrm{s}]$ |
| $\omega_{\text {sdof }}$ |  |
| $\boldsymbol{\phi}_{\boldsymbol{i}}$ | Natural mode or eigenvector $[-]$ |

## 1 Introduction

### 1.1 Background

The interest in design of railway bridges that allows travelling of high-speed trains has increased in recent years. In 2006 the Swedish railway administration updated their bridge code regarding the dynamic design of these railway bridges. Before this update the dynamic aspects were only considered through a dynamic amplification factor in static analysis. The new requirements though demand an additional extensive dynamic analysis to be performed for all bridges subjected to high-speed trains travelling in a speed faster than $200 \mathrm{~km} / \mathrm{h}$. The demand is in line with the requirements and recommendations from the annex used in connection with Eurocode.

The additional requirement of dynamic analyses means that both static and dynamic analyses are required for the design of railway bridges today. It is possible that bridges that are structural sound when statically loaded are not acceptable when considering its dynamic response. If the dynamic analysis proves unsatisfactory both static and dynamic analyses have to be re-made making the design process time consuming and expensive.

Dynamic analysis is very time consuming to execute since the code specifies several complex train loads to be used for a range of train velocities. Hence, there is a need for simplifications and guidelines for the dynamic analysis to make it more time efficient. There is also a need for guidelines in the design process that considers the dynamic aspects, to be used parallel to static analysis and to prevent complications that arise when dynamic criteria proves unfulfilled.

### 1.2 Previous work

This master thesis is a continuation from three projects carried out at Reinertsen Sweden AB during 2007 and 2008. In these theses it has been shown that it is possible to simplify a single-span bridge structure with a single-degree-of-freedom system (SDOF) and attempts have been made to create guidelines for the dynamic analysis of railway bridges through FE-analyses, graphical presentation and the SDOF model.

In Ekström and Kieri (2007) focus was put on the comparison of displacements and accelerations for a SDOF system subjected to various types of loading. It was shown that a SDOF system can be used to approximate the acceleration response by comparing the result with a FE-model. The continuation of this thesis, performed by De Leon and Lasn (2008), treated the transformation procedure of railway bridges into a SDOF model by using an alternative procedure called the force scaling approach. This approach is more compatible with complex structures and was also successfully applied on portal frames. De Leon and Lasn (2008) also made an attempt to establish a graphical guideline with its basis in the SDOF model. In both Ekström and Kieri (2007) and De Leon and Lasn (2008) calculations were made to increase the understanding of the dynamic response behaviour of railway bridges.

In Gustafsson (2008) an attempt to build a 2D FE-model in Matlab was successfully made. It was stated that the FE program reduces the calculation time required to make
a 2D analysis with a certain set of predefined geometrical and material parameters and hence could be used to simplify the dynamic calculations. Different methods of modelling the train load were also examined.

### 1.3 Aim

The overall aim of this master thesis was to find guidelines and simplifications for the dynamic response for railway bridges subjected to high speed trains. This includes easier and faster calculations procedures, guidance to suitable parameter combinations and tools for verification of the advanced calculations.

This thesis aims to investigate the compatibility of the previous developed SDOF model on multi-span bridges. The investigation includes how such a simplification shall be carried out and what possibilities and limitations the model has.

The thesis aims also to increase the understanding of the dynamic response in railway bridges and the differences between 2D and 3D analysis.

### 1.4 Method

An extensive literature study has been performed in the beginning of the master thesis to increase the knowledge about design and dynamic behaviour of railway bridges. Several master theses have been written at the technical universities of Chalmers, KTH and Lund about dynamic design of railway bridges during the last years which all have been studied. Further, Swedish and European design codes have been studied to increase the knowledge of demands and requirements for the dynamic behaviour of railway bridges subjected to high-speed trains.

The advances made in creating a SDOF system by Ekström and Kieri (2007) and De Leon and Lasn (2009) was used as a basis for the development of a SDOF model that is compatible with multi-span bridges.

Examination of the individual influence of bridge parameters on the dynamic response was made to increase the understanding of the dynamic behaviour in railway bridges. This increased understanding together with accumulated knowledge of simplified FE modelling, the possibilities in SDOF systems and graphical presentation was used in the development of possible guidelines and simplifications of the dynamic calculations.

Differences between the dynamic response in 2D and 3D were studied by examining single-span bridges subjected to eccentric loading and with variation in width.

Matlab was used to create programs for the required SDOF and FE calculations. The commercial software ADINA was used to verify the created Matlab programs.

### 1.5 Limitations

To manage the thesis over the limited time schedule several limitations have been made. Some of these limitations are defined by the national bridge code BV Bro Banverket (2006) as this thesis is limited to the required calculation procedure described in this code.

The train load was assumed to be the HSLM model described in BV Bro. This model treats the train load as moving point loads and consequently the train mass, rail irregularities and interaction between train and rail have not been considered.

In the bridge code there are demands on the dynamic response of railway bridges concerning vertical deflection, vertical acceleration, horizontal transverse deflection, twist of the bridge deck and uplifting of the bearings. Here the studied response was the vertical acceleration as this design criterion often governs the dynamic aspect of railway bridge design.

Throughout the thesis the material parameters are chosen to resemble concrete. A linear viscoelastic material model is consequently used. The section is hence treated as uncracked and the reinforcement in the bridge is neglected.

### 1.6 General layout

The outline of the report consists of three major Chapters. Chapter 2 and 3 are together an introduction to the field of structural dynamics and the railway bridge code relevant to this thesis.

In Chapter 4 dynamic analyses in 2D for railway bridges are investigated extensively. The Chapter treats the determination of eigenfrequencies for multi-span bridges, the transformation of multi-span bridges to a SDOF model, individual examinations of bridge parameters effect on the acceleration response and finally an attempt to create guidelines for the dynamic design of railway bridges considering 2D analyses.

Chapter 5 treats dynamic analysis of railway bridges in 3D. The Chapter shows the additional complications that arise when 3D geometry is considered.

More extensive presentations of the layouts in Chapter 4 and 5 are presented in the introduction of each Chapter. A summary, discussion and suggestions for further studies are presented in the end of the thesis.

## 2 Basic Dynamics

This chapter treats fundamental theory of structural dynamics that is vital for the examinations made in this thesis. Theory behind discrete parameter models described by one and multiple degrees of freedom is presented, together with application of the so called mode superposition method which is very important for model reduction with consideration to railway bridge dynamics. Some short comments on numerical integration and continuous models are also presented and the dynamic resonance phenomenon is explained.

### 2.1 SDOF systems

A system that consists of a single-degree-of-freedom is called a SDOF system. The SDOF system is a good application for understanding the fields of dynamics. Since mode superposition (which is introduced in section 2.5) basically is a way to transform larger systems into independent SDOF systems, its solution is also vital for the fields of dynamics. A mass-spring-dashpot model is a good example of a SDOF system and will be used in the following section, see Figure 2.1.


Figure 2.1 SDOF system in the form of a mass-spring-dashpot model.
Newton's second law is required for the derivation of the equation of motion for the SDOF system and is defined as

$$
\begin{equation*}
\sum F=M \ddot{u} \tag{2.1}
\end{equation*}
$$

Figure 2.2 shows all forces acting on the mass in Figure 2.1.


Figure 2.2 Forces acting on the mass in the mass-spring-dashpot model
The forces in Figure 2.2 and Newton's second law in equation (2.1) give together the equation of motion for a single degree of freedom system as

$$
\begin{equation*}
M \ddot{u}+C \dot{u}+K u=p(t) \tag{2.2}
\end{equation*}
$$

Solving the equation of motion can be done either analytically or numerically. Often it is not possible to solve more complex systems analytically. However it is possible for a SDOF system subjected to a sinusoidal load, and this analytical solution will for the purpose of explaining the fundamentals in structural dynamics be presented in the upcoming section.

### 2.2 Analytical solution of an SDOF

The analytical solution to a SDOF system subjected to a sinusoidal load differs whether damping is considered or not. The derivation procedure is also somewhat different and therefore independent sub-sections will treat the undamped and damped case respectively in this section.

### 2.2.1 Undamped SDOF system

Consider the SDOF system in Figure 2.1 without damping subjected to a sinusoidal load. This system has the equation of motion in the form

$$
\begin{equation*}
M \ddot{u}+K u=p_{0} \cos (\Omega t) \tag{2.3}
\end{equation*}
$$

Where the solution can be divided in two parts: the particular and complementary solution. In the field of structural dynamics the terms forced and natural motion are used, respectively, and they together form the total response, i.e.

$$
\begin{equation*}
u(t)=u_{p}(t)+u_{c}(t) \tag{2.4}
\end{equation*}
$$

Consider first the natural motion $u_{c}$ or, as it is also called, the free vibration. The natural motion is the general solution to the equation of motion when the loading is absent. It is convenient to rewrite equation (2.3) for solving the natural motion as

$$
\begin{equation*}
\ddot{\mathrm{u}}+\omega_{\mathrm{n}}^{2} \cdot \mathrm{u}=0 \tag{2.5}
\end{equation*}
$$

where

$$
\omega_{n}=\sqrt{\frac{K}{M}}=\text { undamped natural frequency }
$$

The general solution to equation (2.5) can be written in the form of trigonometric functions as

$$
\begin{equation*}
u_{c}(t)=A_{1} \cdot \cos \left(\omega_{n} \cdot t\right)+A_{2} \cdot \sin \left(\omega_{n} \cdot t\right) \tag{2.6}
\end{equation*}
$$

where $\mathrm{A}_{1}$ and $\mathrm{A}_{2}$ are constants that are determined by initial conditions.
Consider now the forced motion of the system created by the sinusoidal load. It is seen that the forced motion from equation (2.3) will have the form

$$
\begin{equation*}
u_{p}(t)=U \cdot \cos (\Omega t) \tag{2.7}
\end{equation*}
$$

where U is the amplitude of the forced motion. By inserting equation (2.7) into (2.3) an expression for the amplitude $U$ can be derived as

$$
\begin{equation*}
U=\frac{p_{0}}{K-M \Omega^{2}} \tag{2.8}
\end{equation*}
$$

This expression can (for a pedagogical reason of the upcoming examination of resonance effects) be rewritten as

$$
\begin{equation*}
U=\frac{U_{0}}{1-\mathrm{r}^{2}} \tag{2.9}
\end{equation*}
$$

where

$$
\begin{equation*}
U_{0}=\frac{p_{0}}{K} \tag{2.10}
\end{equation*}
$$

is the static displacement and

$$
\begin{equation*}
r=\frac{\Omega}{\omega_{n}} \tag{2.11}
\end{equation*}
$$

is the frequency ratio. The total response for the undamped SDOF system as the sum of the forced and natural motion can now be concluded to be

$$
\begin{equation*}
u(t)=\frac{U_{0}}{1-\mathrm{r}^{2}} \cdot \cos (\Omega t)+A_{1} \cdot \cos \left(\omega_{n} \cdot t\right)+A_{2} \cdot \sin \left(\omega_{n} \cdot t\right) \tag{2.12}
\end{equation*}
$$

### 2.2.2 Damped SDOF system

The derivation of the analytical solution for a damped SDOF system is more advanced than the derivation without consideration to damping in section 2.2.1. However, the approach is the same and as for the case without damping the total response is divided in forced and natural motion which solutions are solved separately.

The equation of motion for the SDOF system that considers damping and is subjected to a sinusoidal load has the appearance

$$
\begin{equation*}
M \ddot{u}+C \dot{u}+K u=p_{0} \cos (\Omega t) \tag{2.13}
\end{equation*}
$$

As for the case without damping it is convenient to rewrite this equation for the derivation of the natural motion as

$$
\begin{equation*}
\ddot{u}+2 \zeta \omega_{n} \dot{u}+\omega_{n}^{2} u=0 \tag{2.14}
\end{equation*}
$$

where

$$
\begin{aligned}
& \omega_{n}=\sqrt{\frac{K}{M}}=\text { undamped natural frequency } \\
& \zeta=\frac{c}{c_{c r}}=\text { viscous damping factor } \\
& c_{c r}=2 \sqrt{K M}=\text { critical damping coefficient }
\end{aligned}
$$

To solve the natural motion the form of the solution is assumed as

$$
\begin{equation*}
u(t)=\bar{C} e^{\bar{s} t} \tag{2.15}
\end{equation*}
$$

By inserting equation (2.15) into (2.14) we obtain the characteristic equation

$$
\begin{equation*}
\bar{s}^{2}+2 \cdot \zeta \cdot \omega_{n}+\omega_{n}^{2}=0 \tag{2.16}
\end{equation*}
$$

which roots can be solved as

$$
\begin{equation*}
\bar{s}_{1,2}=-\zeta \cdot \omega_{n} \pm \omega_{n} \sqrt{\zeta^{2}-1} \tag{2.17}
\end{equation*}
$$

The size of $\zeta$ can be divided in three cases:

- $0<\zeta<1$ : underdamped
- $\zeta=1: \quad$ critically damping
- $\zeta>1$ overdamped

The solution to the natural motion will differ depending on which case that is considered. In this thesis only the underdamped case will be treated as the viscous damping factor usually lies in the range of 0 to 0.05 for real life structures (like railway bridges). For information of the other cases, see for example Craig and Kurdila (2006).

By the use of Euler's formula the solution of the underdamped case can be derived to be

$$
\begin{equation*}
u(t)=e^{-\zeta \omega_{n} t}\left(A_{1} \cdot \cos \left(\omega_{d} t\right)+A_{2} \cdot \sin \left(\omega_{d} t\right)\right) \tag{2.18}
\end{equation*}
$$

where

$$
\begin{equation*}
\omega_{d}=\omega_{n} \sqrt{1-\zeta^{2}} \tag{2.19}
\end{equation*}
$$

is called the damped natural frequency.
Consider now the forced motion of equation (2.13). The solution can be written as

$$
\begin{equation*}
u_{p}=U \cdot \cos (\Omega t-\alpha) \tag{2.20}
\end{equation*}
$$

where $U$ is the amplitude and $\alpha$ is the phase angle. The velocity and acceleration of the forced response then become

$$
\begin{align*}
& \dot{u}_{p}=-\Omega U \cdot \sin (\Omega t-\alpha)  \tag{2.21}\\
& \ddot{u}_{p}=-\Omega^{2} U \cdot \cos (\Omega t-\alpha) \tag{2.22}
\end{align*}
$$

The expressions for the displacement, velocity and acceleration of the forced response, inserted into equation (2.13), give that

$$
\begin{align*}
& -M \Omega^{2} U \cdot \cos (\Omega t-\alpha)-C \Omega U \cdot \sin (\Omega t-\alpha)+  \tag{2.23}\\
& +K U \cdot \cos (\Omega t-\alpha)=p_{0} \cos (\Omega t)
\end{align*}
$$

from which it is obtained that

$$
\begin{equation*}
U=\frac{p_{0}}{\sqrt{\left(K-M \Omega^{2}\right)^{2}+(C \Omega)^{2}}}=\frac{U_{0}}{\sqrt{\left(1-\mathrm{r}^{2}\right)^{2}+(2 \zeta r)^{2}}} \tag{2.24}
\end{equation*}
$$

and

$$
\begin{equation*}
\tan \alpha=\frac{C \Omega}{K-M \Omega^{2}}=\frac{2 \zeta r}{1-\mathrm{r}^{2}} \tag{2.25}
\end{equation*}
$$

The total response of the damped SDOF system can now be written as

$$
\begin{align*}
& u(t)=\frac{U_{0}}{\sqrt{\left(1-\mathrm{r}^{2}\right)^{2}+(2 \zeta r)^{2}}} \cdot \cos (\Omega t-\alpha)+  \tag{2.26}\\
& +e^{-\zeta \omega_{n} t}\left(A_{1} \cdot \cos \left(\omega_{d} t\right)+A_{2} \cdot \sin \left(\omega_{d} t\right)\right)
\end{align*}
$$

### 2.3 Resonance

The reason why a dynamic analysis is required in structures subjected to frequency loading is because of the so called resonance effects. Resonance occurs when $r \approx 1$, i.e. when the loading frequency is approximately equal to the natural or damped natural frequency. Resonance can cause very large amplitudes of the response compared to the corresponding static loading. To examine this consider the dynamic magnification factor defined as the ratio between dynamic and static amplitude, i.e.

$$
\begin{equation*}
D(r)=\frac{U(r)}{U_{0}} \tag{2.27}
\end{equation*}
$$

With $U$ defined as in equations (2.9) and (2.24) the dynamic magnification factor becomes

$$
\begin{equation*}
D(r)=\frac{1}{1-\mathrm{r}^{2}} \tag{2.28}
\end{equation*}
$$

for an undamped SDOF system and

$$
\begin{equation*}
D(r)=\frac{1}{\sqrt{\left(1-\mathrm{r}^{2}\right)^{2}+(2 \zeta r)^{2}}} \tag{2.29}
\end{equation*}
$$

for a damped SDOF system. Figure 2.3 shows the dynamic magnification factor for different values of the viscous damping factor and variation of frequency ratio.


Figure $2.3 \quad$ Dynamic magnification factor for variation in damping and frequency ratio.
Higher values of damping cause large differences between the natural and damped natural frequency. It can hence be seen in Figure 2.3 that resonance occurs for lower values of the frequency ratio for higher damping. For the undamped system the dynamic magnification factor will go towards infinity as $r \rightarrow 1$.

### 2.4 MDOF systems

The dynamic analyses of most real structures are based on multiple-degree-offreedom systems, MDOF system. To show how the equation of motion for a MDOF system is assembled consider the lumped mass model shown in Figure 2.4.


Figure $2.4 \quad$ A three-degree-of-freedom system with lumped mass and viscous damping
Figure 2.5 shows all forces acting on the three masses in Figure 2.4.


By use of Newton's Second law three force equilibriums can be established as

$$
\begin{align*}
& m_{1} \ddot{u}_{1}=p_{1}(t)+f_{2}+f_{5}-f_{1}-f_{4} \\
& m_{2} \ddot{u}_{2}=p_{2}(t)+f_{3}+f_{6}-f_{2}-f_{5}  \tag{2.30}\\
& m_{3} \ddot{u}_{3}=p_{3}(t)-f_{3}-f_{6}
\end{align*}
$$

where the forces in Figure 2.5 are defined as

$$
\begin{align*}
& f_{1}=k_{1} u_{1} \\
& f_{2}=k_{2}\left(u_{2}-u_{1}\right) \\
& f_{3}=k_{3}\left(u_{3}-u_{2}\right)  \tag{2.31}\\
& f_{4}=c_{1} \dot{u}_{1} \\
& f_{5}=c_{2}\left(\dot{u}_{2}-\dot{u}_{1}\right) \\
& f_{6}=c_{3}\left(\dot{u}_{3}-\dot{u}_{2}\right)
\end{align*}
$$

By combining equations (2.30) and (2.31) the force equilibrium can be written in matrix form as

$$
\begin{align*}
& \left(\begin{array}{ccc}
m_{1} & 0 & 0 \\
0 & m_{2} & 0 \\
0 & 0 & m_{3}
\end{array}\right)\left(\begin{array}{l}
\ddot{u}_{1} \\
\ddot{u}_{2} \\
\ddot{u}_{3}
\end{array}\right)+\left(\begin{array}{ccc}
c_{1}+c_{2} & -c_{2} & 0 \\
-c_{2} & c_{2}+c_{3} & -c_{3} \\
0 & -c_{3} & c_{3}
\end{array}\right)\left(\begin{array}{l}
\dot{u}_{1} \\
\dot{u}_{2} \\
\dot{u}_{3}
\end{array}\right)+\cdots  \tag{2.32}\\
& +\left(\begin{array}{ccc}
k_{1}+k_{2} & -k_{2} & 0 \\
-k_{2} & k_{2}+k_{3} & -k_{3} \\
0 & -k_{3} & k_{3}
\end{array}\right)\left(\begin{array}{l}
u_{1} \\
u_{2} \\
u_{3}
\end{array}\right)=\left(\begin{array}{l}
p_{1}(t) \\
p_{2}(t) \\
p_{3}(t)
\end{array}\right)
\end{align*}
$$

or using matrix format as

$$
\begin{equation*}
\boldsymbol{M} \ddot{\boldsymbol{u}}+\boldsymbol{C} \dot{\boldsymbol{u}}+\boldsymbol{K} \boldsymbol{u}=\boldsymbol{p}(t) \tag{2.33}
\end{equation*}
$$

Equation (2.33) represents the equation of motion of a MDOF system. The appearance of equation (2.33) is the same for two- and three-dimensional analysis and forms the basis to the finite element method and hence also dynamic analysis of railway bridges.

### 2.4.1 Eigenfrequencies and eigenmodes

Consider now the free vibration of a MDOF system with $N$ degrees of freedom. The corresponding equation of motion for the free vibration is defined as

$$
\begin{equation*}
M \ddot{u}+K u=\mathbf{0} \tag{2.34}
\end{equation*}
$$

where $\boldsymbol{M}$ is the mass matrix and $\boldsymbol{K}$ is the stiffness matrix. Both $\boldsymbol{M}$ and $\boldsymbol{K}$ have the size $N \mathrm{x} N$. As for a SDOF system a harmonic solution to equation (2.34) is assumed, i.e.

$$
\begin{equation*}
\boldsymbol{u}(t)=\boldsymbol{U} \cos \left(\omega_{n} t-\alpha\right) \tag{2.35}
\end{equation*}
$$

which inserted in equation (2.34) gives that

$$
\begin{equation*}
\left[\boldsymbol{K}-\omega_{n}^{2} \boldsymbol{M}\right] \boldsymbol{U}=\mathbf{0} \tag{2.36}
\end{equation*}
$$

Since $\boldsymbol{U} \neq 0$ it is required that

$$
\begin{equation*}
\operatorname{det}\left(\boldsymbol{K}-\omega_{n}^{2} \boldsymbol{M}\right)=\mathbf{0} \tag{2.37}
\end{equation*}
$$

which is the characteristic equation of the MDOF system. There are $N$ roots to equation (2.37) , $N$ eigenvalues, each corresponding to a squared eigenfrequency. For each eigenvalue there is a corresponding eigenvector, or natural mode. The determination of the eigenfrequencies and natural modes for a MDOF system is most commonly made numerically.

### 2.4.2 Properties of natural modes

The solution to mode superposition, presented in section 2.5 , is based on two properties of the natural modes, scaling and orthogonality, which will be presented in this section.

### 2.4.2.1 Scaling

The natural modes' shapes are unique, i.e. for a specific value of one mode's first element the remaining elements are fixed. However the amplitudes of the natural modes are arbitrary and they can hence be scaled in any appropriate manner. A scaled mode is in this thesis referred to as $\boldsymbol{\phi}_{\boldsymbol{i}}$, where $i=1 \ldots N$. A common procedure is to scale the modes so that the modal mass, which is defined as

$$
\begin{equation*}
M_{i}=\boldsymbol{\phi}_{\boldsymbol{i}}^{T} \boldsymbol{M} \boldsymbol{\phi}_{\boldsymbol{i}} \tag{2.38}
\end{equation*}
$$

becomes equal to one. This scaling procedure has been used in this thesis and is assumed for all modes in the remainder of the thesis. If equation (2.36) is multiplied by $\boldsymbol{\phi}_{\boldsymbol{i}}{ }^{T}$ it is obtained that

$$
\begin{equation*}
\boldsymbol{\phi}_{i}^{\mathrm{T}} \boldsymbol{K} \boldsymbol{\phi}_{i}=\omega_{r}^{2}\left(\boldsymbol{\phi}_{i}^{\mathrm{T}} \boldsymbol{M} \boldsymbol{\phi}_{i}\right) \tag{2.39}
\end{equation*}
$$

and hence that the modal stiffness can be defined as

$$
\begin{equation*}
K_{i}=\boldsymbol{\phi}_{i}^{\mathrm{T}} \boldsymbol{K} \boldsymbol{\phi}_{i}=\omega_{i}^{2} M_{i}=\omega_{i}^{2} \tag{2.40}
\end{equation*}
$$

### 2.4.2.2 Orthogonality

A very important property of natural modes is that they are orthogonal to each other. It can be shown that

$$
\begin{align*}
\boldsymbol{\phi}_{i}^{\mathrm{T}} \boldsymbol{K} \boldsymbol{\phi}_{j} & =0  \tag{2.41}\\
\boldsymbol{\phi}_{i}^{\mathrm{T}} \boldsymbol{M} \boldsymbol{\phi}_{j} & =0
\end{align*}
$$

where $\boldsymbol{\phi}_{i}$ and $\boldsymbol{\phi}_{j}$ are two different natural modes.

### 2.5 Mode superposition

Mode superposition is extensively used in dynamic analysis, and can be used for both solving dynamic problems analytically and simplifying numerical calculations. To employ mode superposition the modal matrix must first be defined as a matrix which columns represent the scaled natural modes, i.e.

$$
\Phi=\left[\begin{array}{llll}
\phi_{1} & \phi_{2} & \ldots & \phi_{N} \tag{2.42}
\end{array}\right]
$$

The concept of mode superposition lies in the possibility of representing the true response as

$$
\begin{equation*}
\boldsymbol{u}(t)=\sum_{i=1}^{N} \boldsymbol{\phi}_{i} \eta_{i}(t)=\boldsymbol{\Phi} \boldsymbol{\eta}(t) \tag{2.43}
\end{equation*}
$$

Where $\eta_{i}$ is called the modal coordinate of mode $i$. If the equation of motion is multiplied with $\boldsymbol{\Phi}^{\mathrm{T}}$ and combined with equation (2.43) it is obtained that

$$
\begin{equation*}
\widehat{\boldsymbol{M}} \ddot{\boldsymbol{\eta}}+\widehat{\boldsymbol{C}} \dot{\boldsymbol{\eta}}+\widehat{\boldsymbol{K}} \boldsymbol{\eta}=\boldsymbol{\Phi}^{\mathrm{T}} \boldsymbol{p}(t) \tag{2.44}
\end{equation*}
$$

where

$$
\begin{align*}
& \widehat{\boldsymbol{M}}=\boldsymbol{\Phi}^{\mathbf{T}} \boldsymbol{M} \boldsymbol{\Phi}=\operatorname{diag}\left(M_{1}, M_{2}, \ldots, M_{N}\right)=\boldsymbol{I} \\
& \widehat{\boldsymbol{K}}=\boldsymbol{\Phi}^{\mathbf{T}} \boldsymbol{K} \boldsymbol{\Phi}=\operatorname{diag}\left(K_{1}, K_{2}, \ldots, K_{N}\right)=\operatorname{diag}\left(\omega_{1}^{2}, \omega_{2}^{2}, \ldots, \omega_{N}^{2}\right)  \tag{2.45}\\
& \widehat{\boldsymbol{C}}=\boldsymbol{\Phi}^{\mathbf{T}} \boldsymbol{C} \boldsymbol{\Phi}
\end{align*}
$$

are denoted the modal mass matrix, modal stiffness matrix and modal damping matrix respectively. It should be note that the modal mass and stiffness matrix are diagonals because of the orthogonality of the natural modes. This means that the modal coordinates are uncoupled with regard to stiffness and mass.

Whether the modal coordinates are uncoupled with regard to damping depends on the damping's definition. In several design cases it is not possible to establish a diagonal modal damping matrix, for example when there is a presence of energy absorbers or viscous damping. However for many structures (including railway bridges) it is not possible to define such a damping mechanisms. In those cases it is common to use damping definitions that lead to a diagonal modal damping matrix. Two such methods will be presented in this thesis and are referred to as Raleigh damping and modal damping.

### 2.5.1 Rayleigh damping

In Rayleigh damping the damping matrix is defined as a linear combination of the stiffness and mass matrix, i.e.

$$
\begin{equation*}
\boldsymbol{C}=a_{o} \boldsymbol{M}+a_{1} \boldsymbol{K} \tag{2.46}
\end{equation*}
$$

where $a_{o}$ and $a_{1}$ are constants.
The modal damping matrix then becomes

$$
\begin{equation*}
\widehat{\boldsymbol{C}}=\boldsymbol{\Phi}^{\mathbf{T}} \boldsymbol{C} \boldsymbol{\Phi}=\operatorname{diag}\left(\left(a_{0}+a_{1} \omega_{i}^{2}\right) M_{i}\right)=\operatorname{diag}\left(2 \zeta_{i} \omega_{i} M_{i}\right) \tag{2.47}
\end{equation*}
$$

where $\zeta_{i}$ is called the modal damping factor and corresponds to the viscous damping factor defined for the SDOF system. For the last step in equation (2.47) to be true it is required that

$$
\begin{equation*}
\zeta_{i}=\frac{1}{2}\left(\frac{a_{0}}{\omega_{i}}+a_{1} \omega_{i}\right) \tag{2.48}
\end{equation*}
$$

Basically Rayleigh damping is utilized by choosing two modal damping factors and calculating $\mathrm{a}_{0}$ and $\mathrm{a}_{1}$ based on these values.

### 2.5.2 Modal damping

Rayleigh damping has the obvious disadvantage of only having the possibility of specifying the modal damping factor for two natural modes. However the modal damping method allows specification of all modal damping factors. A disadvantage compared to Rayleigh damping is the somewhat more complicated calculation of the damping matrix. However, when using mode superposition the damping matrix is commonly not required but only the modal damping matrix. In the modal damping method the modal damping matrix is basically assumed to satisfy

$$
\begin{equation*}
\widehat{\boldsymbol{C}}=\boldsymbol{\Phi}^{\mathbf{T}} \boldsymbol{C} \boldsymbol{\Phi}=\operatorname{diag}\left(2 \zeta_{i} \omega_{i} M_{i}\right) \tag{2.49}
\end{equation*}
$$

It can be shown that the damping matrix based on this assumption can be calculated as

$$
\begin{equation*}
\boldsymbol{C}=\sum_{\mathrm{i}=1}^{\mathrm{N}} \frac{2 \zeta_{\mathrm{i}} \omega_{\mathrm{i}}}{\mathrm{M}_{\mathrm{i}}}\left(\boldsymbol{M} \boldsymbol{\phi}_{r}\right)\left(\boldsymbol{M} \boldsymbol{\phi}_{r}\right)^{\mathrm{T}} \tag{2.50}
\end{equation*}
$$

### 2.5.3 Advantages with mode superposition

The modal coordinates become completely uncoupled in the use of mode superposition if a definition of damping that gives a diagonal modal damping matrix is used. Basically $N$ SDOF systems are obtained in the form

$$
\begin{equation*}
M_{i} \ddot{u}_{i}+2 \zeta_{i} \omega_{i} M_{i} \dot{u}_{i}+K_{i} u_{i}=\boldsymbol{\phi}_{i}^{\mathrm{T}} \boldsymbol{p}(t) \tag{2.51}
\end{equation*}
$$

If a simple load acts on the system it may be possible to solve the dynamic problem analytically. As an example consider the case where a sinusoidal load acts on the system. The solution then becomes similar to that of equation (2.26). However in many situations the load is so complex that numerical integration procedures still are preferred, see section 2.6.

However, even when the calculations are solved numerically there is a great use in mode superposition in the form of model reduction. Model reduction using mode superposition is made by only including a specific number of eigenmodes. It is assumed that

$$
\begin{equation*}
\boldsymbol{u}(t)=\sum_{i=1}^{M} \boldsymbol{\phi}_{i} \eta_{i}(t)=\boldsymbol{\Phi}_{\text {red }} \boldsymbol{\eta}_{\text {red }}(t) \tag{2.52}
\end{equation*}
$$

where $M<N$ and

$$
\Phi_{\text {red }}=\left[\begin{array}{llll}
\boldsymbol{\phi}_{1} & \phi_{2} & \ldots & \boldsymbol{\phi}_{M} \tag{2.53}
\end{array}\right]
$$

is the reduced modal matrix and

$$
\boldsymbol{\eta}_{\text {red }}=\left[\begin{array}{llll}
\eta_{1} & \eta_{2} & \ldots & \eta_{M} \tag{2.54}
\end{array}\right]^{\mathrm{T}}
$$

contain the reduced modal matrix's corresponding modal coordinates. If the equation of motion is multiplied with $\boldsymbol{\Phi}_{\text {red }}{ }^{\mathrm{T}}$ and combined with equation (2.52) it is obtained that

$$
\begin{equation*}
\widehat{\boldsymbol{M}}_{r e d} \ddot{\boldsymbol{\eta}}_{\text {red }}+\widehat{\boldsymbol{C}}_{\text {red }} \dot{\boldsymbol{\eta}}_{\text {red }}+\widehat{\boldsymbol{K}}_{\text {red }} \boldsymbol{\eta}_{\text {red }}=\boldsymbol{\Phi}_{\text {red }}^{\mathrm{T}} \boldsymbol{p}(t) \tag{2.55}
\end{equation*}
$$

where

$$
\begin{align*}
& \widehat{\boldsymbol{M}}_{\text {red }}=\boldsymbol{\Phi}_{\text {red }}^{\mathrm{T}} \boldsymbol{M} \boldsymbol{\Phi}_{\text {red }}=\operatorname{diag}\left(M_{1}, \ldots, M_{M}\right)=\text { unit matrix }=\boldsymbol{I} \\
& \widehat{\boldsymbol{K}}_{\text {red }}=\boldsymbol{\Phi}_{\text {red }}^{\mathrm{T}} \boldsymbol{K} \boldsymbol{\Phi}_{\text {red }}=\operatorname{diag}\left(\omega_{1}^{2}, \ldots, \omega_{M}^{2}\right)  \tag{2.56}\\
& \widehat{\boldsymbol{C}}_{\text {red }}=\boldsymbol{\Phi}_{\text {red }}^{\mathrm{T}} \boldsymbol{C} \boldsymbol{\Phi}_{\text {red }}=\operatorname{diag}\left(2 \zeta_{1} \omega_{1}, \ldots, 2 \zeta_{M} \omega_{M}\right)
\end{align*}
$$

It should be noted that equation (2.52) is an estimation of the true response. But the error will be small in many design situations as modes of higher order tend to give very low contribution to the dynamic response. Say that the response can be well estimated using twenty eigenmodes. Then a system containing perhaps several hundred degrees of freedom can be reduced to only twenty, which greatly reduces calculation time for the numerical integration procedures.

### 2.6 Numerical integration

To obtain the acceleration response over time there are in many design situations a need for numerical integration. Such procedures do not require any transformation of the equation of motion and are therefore commonly referred to as direct integration. The procedures are required in many design situations, for example when uncoupled modal coordinates cannot be established, when the system contains nonlinearities or when the load has a complicated time history.

There are several direct integration procedures used in the field of structural dynamics and these can be divided in two general groups. The first group directly treats the second order differential equation (the equation of motion) and the other group handles it through an equivalent set of first order differential equations. In the first group some of the most commonly used integration procedures are the central difference method, the Newmark- $\beta$ method and the Wilson- $\theta$ method. One example from the second group is the Sub-zero order hold method.

In this section focus have been put on the presentation of the Newmark- $\beta$ method which is the only numerical integration procedure used in this thesis. The reader is referred to Craig and Kurdila (2006) and Abrahamsson (2000) for insight in the other mentioned numerical integration procedures.

### 2.6.1.1 Newmark- $\beta$ method

The main assumption in the Newmark- $\beta$ method is that the systems displacement and velocity can be written as

$$
\begin{align*}
& \dot{\boldsymbol{u}}_{n+1}=\dot{\boldsymbol{u}}_{n}+\left[(1-\gamma) \ddot{\boldsymbol{u}}_{n}+\gamma \ddot{\boldsymbol{u}}_{n+1}\right] \Delta t  \tag{2.57}\\
& \boldsymbol{u}_{n+1}=\boldsymbol{u}_{n}+\dot{\boldsymbol{u}}_{n} h+\left[(1-2 \beta) \ddot{\mathbf{u}}_{\mathrm{n}}+2 \beta \ddot{\mathbf{u}}_{\mathrm{n}+1}\right] \frac{\Delta t^{2}}{2} \tag{2.58}
\end{align*}
$$

where the parameters $\beta$ and $\gamma$ are the integration coefficients and used to obtain integration stability and accuracy. Newmark introduced an unconditionally stable scheme with the relation $\beta=0.5$ and $\gamma=0.25$ and the method is therefore said to be an unconditionally stable implicit integration scheme. The Newmark- $\beta$ method using this relation for $\beta$ and $\gamma$ is called the Constant Average Acceleration method (or the trapezoidal rule). The Constant Average Acceleration method has no demand on the incremental time step to reach a stable solution, the acceleration variation is shown in Figure 2.6.


Figure $2.6 \quad$ Newmark's constant-average-acceleration scheme.

### 2.7 Inverse iteration

In section 0 it was mentioned that the eigenfrequencies and eigenvectors of MDOF systems in general are calculated numerically. In many situations more general procedures for solving eigenvalue problems are utilized. However when only a few of the lowest eigenvalues are required there are simpler and less time consuming alternatives. Inverse iteration is a procedure for calculating the first eigenfrequency and eigenvector. Consider again the eigenvalue problem

$$
\begin{equation*}
\left(\boldsymbol{K}-\lambda_{\mathrm{i}} \boldsymbol{M}\right) \boldsymbol{\phi}_{\mathrm{i}}=0 \tag{2.59}
\end{equation*}
$$

where $\lambda_{i}=\omega_{i}{ }^{2}$. The inverse iteration method is in the subgroup of vector iterations methods of solving eigenvalue problems and directly employs equation (2.59). It is possible to write the equation in the form

$$
\begin{equation*}
\phi=\lambda D \phi \tag{2.60}
\end{equation*}
$$

where

$$
\begin{equation*}
D=K^{-1} M \tag{2.61}
\end{equation*}
$$

is called the dynamical matrix. The iterative procedure of the inverse iteration method follows a few specific steps. A first guess of the eigenvector $\boldsymbol{u}_{s}$, is initially assumed and then a new vector $\boldsymbol{v}_{\boldsymbol{s}+\boldsymbol{1}}$ is calculated using the dynamical matrix as

$$
\begin{equation*}
v_{s+1}=D u_{s} \tag{2.62}
\end{equation*}
$$

The new guess for the next iteration is then calculated as

$$
\begin{equation*}
\boldsymbol{u}_{s+1}=\lambda_{\mathbf{s}+1} v_{\mathrm{s}+1} \tag{2.63}
\end{equation*}
$$

where $\lambda_{s+1}$ is a scaling factor for $\boldsymbol{v}_{\mathrm{s}+1}$ chosen in a consistent manner, for example to get the highest value equal to one. According to equation (2.60) $\boldsymbol{u}_{\mathrm{s}+1}$ becomes equal to $\boldsymbol{u}_{\mathrm{s}}$ and the scaling factor becomes the corresponding eigenvalue, if an iteration is made using the correct eigenvector. When the initial guess deviates from the true eigenvector this result is achieved after some iterations.

An alternative for the definition of the scaling factor is to use the Rayleigh quotient, i.e.

$$
\begin{equation*}
\lambda_{s+1}=\frac{\boldsymbol{v}_{s+1}^{\mathrm{T}} \boldsymbol{M} \boldsymbol{u}_{s}}{\boldsymbol{v}_{s+1}^{\mathrm{T}} \boldsymbol{K} \boldsymbol{v}_{s+1}} \tag{2.64}
\end{equation*}
$$

which gives a faster convergence.

### 2.8 Transverse vibration of Bernoulli-Euler Beams

In section 2.2 it was shown how the eigenfrequency and eigenvector can be calculated analytically for a SDOF system and in Section 2.4 it was said that numerical procedures often are required for solving the eigenfrequencies and eigenvectors of a

MDOF system. SDOF and MDOF systems are similar as they belong to the same group, as discrete parameter models. In this section an example of a continuous model will be presented, namely a simply supported beam. Continuous models can be used to derive the exact analytical solution for free vibration of structures, which presented as a discrete parameter model would require a MDOF system. The simply supported beam is of special interest as it can be seen as representation of a single-span railway bridge.

It can be shown that beams that are relatively long and thin, i.e. are applicable of Bernoulli-Euler beam theory, have the differential equation of motion for governing vertical vibration as

$$
\begin{equation*}
\frac{\partial^{2}}{\partial x^{2}}\left(E I \frac{\partial^{2} v}{\partial x^{2}}\right)+\rho A \frac{\partial^{2} v}{\partial t^{2}}=p_{y}(x, t), \quad 0<x<L \tag{2.65}
\end{equation*}
$$

where $E I$ is the bending stiffness, $\rho A$ is the mass per meter length, $x$ is the length coordinate and $v$ is the vertical deflection. If it is assumed that $E I$ is constant along the beam the differential equation of motion becomes

$$
\begin{equation*}
E I v^{\prime \prime \prime \prime}+\rho A \ddot{v}=0 \tag{2.66}
\end{equation*}
$$

As for the discrete parameter models a harmonic solution is assumed as

$$
\begin{equation*}
v(x, t)=V(x) \cdot \cos (\omega t-\alpha) \tag{2.67}
\end{equation*}
$$

and with this solution inserted in equation (2.66) it is obtained that

$$
\begin{equation*}
\frac{d^{4} V}{d x^{4}}-\lambda^{4} V=0 \tag{2.68}
\end{equation*}
$$

where

$$
\begin{equation*}
\lambda^{4}=\omega^{2} \frac{\rho A}{E I} \tag{2.69}
\end{equation*}
$$

The general solution to equation (2.68) can be written as

$$
\begin{equation*}
V(x)=A_{1} \sinh (\lambda x)+A_{2} \cosh (\lambda x)+A_{3} \sin (\lambda x)+A_{4} \cos (\lambda x) \tag{2.70}
\end{equation*}
$$

where $\mathrm{A}_{1,2,3,4}$ are constants determined by the boundary conditions. The boundary conditions for a simply supported beam are

$$
\begin{array}{ll}
v(x=0, t)=0 & v(x=L, t)=0 \\
\left.\frac{\partial^{2} v}{\partial x^{2}}\right|_{x=0}=0 & \left.\frac{\partial^{2} v}{\partial x^{2}}\right|_{x=L}=0 \tag{2.71}
\end{array}
$$

By combining equations (2.70) and (2.71) it can be shown that the eigenfrequencies and eigenvectors of a simply supported beam has the analytical solution

$$
\begin{align*}
& V_{(n)}(x)=\sqrt{\frac{2}{m l}} \sin \frac{n \pi x}{l}  \tag{2.72}\\
& \omega_{(n)}=(n \pi)^{2} \sqrt{\frac{E I}{m l^{4}}} \tag{2.73}
\end{align*}
$$

where

$$
n=1,2,3 \ldots
$$

The shapes of the first three eigenmodes are shown in Figure 2.7.


Figure 2.7
First three eigenmodes for a simply supported beam

## 3 Railway bridge code

In Sweden railway bridges subjected to high speed trains are design according to BV Bro (Banverket, 2006), with regulations set by the Swedish railway administration. This code is similar to the regulations set by Eurocode but differs in some aspects. In this Chapter the demands of dynamic design of railway bridges set by both BV Bro (Banverket, 2006) and Eurocode will be presented.

### 3.1 High-speed load models (HSLM)

Eurocode 1, CEN (2003), states that in a dynamic analysis of a bridge, characteristic values shall be used for the specified trains. The choice of train shall regard every permitted train configuration for each high speed train configuration that is expected to operate the bridge with velocities over $200 \mathrm{~km} / \mathrm{h}$. The dynamic analysis shall on international lines also include high-speed load models (HSLM). Load model HSLM consist of two separate universal train models, HSLM-A and HSLM-B, where the length of the coaches varies.

In BV Bro (Banverket, 2006) there are no demands on additional permitted train configurations to the HSLM loads suggested by Eurocode, and the HSLM loads are hence the only loads used in the dynamic design of railway bridges in Sweden. The HSLM load model consists of two separate universal train configurations, HSLM-A and HSLM-B, where the coach length varies. From Table 3.1 it is possible determine which train configuration that should be applied depending on the total length of the bridge.

Table 3.1 Load models for different structural configurations.

| Structural | Span | $L \geq 7 \mathrm{~m}$ |
| :--- | :--- | :--- |
|  | $L<7 \mathrm{~m}$ | HSLM-A |
| Simply supported <br> one span | HSLM-B | HSLM-A: trains A1 to A10 <br> inclusive |
| Continuous bridge | HSLM-A: trains A1 to A10 <br> inclusive | He |

### 3.1.1 HSLM-A

The HSLM-A load model consist of a load geometry determined by the coach length $D$, the bogie axle spacing $d$ and the numbers of coaches $N$ between the power cars, se Figure 3.1. There is also an intermediate coach after the power car in each train configuration.


Figure 3.1 The load geometry for the train model HSLM-A.
The HSLM-A train load defined in Eurocode 1, CEN (2003), include ten different train configurations, see Table 3.2. The designing dynamic response should be the most critical considering all ten load configurations.

Table 3.2 Train configurations for train HSLM-A.

| Train load | Intermediate <br> coaches <br> N [-] | Coach length | Boogie axel <br> spacing <br> d [m] | Point force |
| :--- | :---: | :---: | :---: | :---: |
| P [kN] |  |  |  |  |$|$| A1 |
| :--- |

### 3.1.2 HSLM-B

The HSLM-B train load consists of equally spaced point loads. Depending on the total span length, Eurocode defines the number of loads $N$ and the spacing $d$, as in Figure 3.2 and Figure 3.3.


Figure 3.2 The load distribution for train model HSLM-B.


Figure 3.3 The spacing and number of loads for HSLM-B.

### 3.1.2.1 Speeds to consider

According to Eurocode 1, CEN (2003), dynamic calculations of railway bridges should be considered from $40 \mathrm{~m} / \mathrm{s}(144 \mathrm{~km} / \mathrm{h})$, with a series of speed steps, up to the $1.2 \times$ maximum design speed. It is mentioned that speed steps should be smaller at speeds where resonance is likely to occur, but no specific value is given.

BV Bro (Banverket, 2006) requires that the dynamic response is examined for a speed interval of $100 \mathrm{~km} / \mathrm{h}$ to the maximum design speed $+20 \%$. The speed steps should be $5 \mathrm{~km} / \mathrm{h}$ in general and $2.5 \mathrm{~km} / \mathrm{h}$ close to resonance effects.

### 3.2 Demands on vertical acceleration

There are several requirements of the dynamic design of railway bridges for the serviceability limit state in Eurocode: Basis of Structural Design Annex A2, CEN (2003). Checks are required for vertical deflection, vertical acceleration, horizontal
transverse deflection, twist of the bridge deck and uplifting of the bearings. The scope of this Master's thesis is to investigate the vertical acceleration in the deck and therefore only this demand is presented in this section.

Bridge deck accelerations shall be limited to $3.5 \mathrm{~m} / \mathrm{s}^{2}$ for ballasted tracks and $5 \mathrm{~m} / \mathrm{s}^{2}$ for direct fastened decks according to both BV Bro (Banverket 2006) and Eurocode: Basis of Structural Design Annex A2, CEN (2003). The limit for vertical deflection along any track is L/600.

### 3.3 Limitation of considered eigenmodes

The criteria for traffic safety in Eurocode: Basis of structural Design Annex A2, CEN (2003), states that all members supporting the track shall consider frequencies and their associated mode shapes up to the greater of:

- 30 Hz
- 1.5 times the frequencies of the fundamental mode of vibration of the member being considered
- The frequencies of the third mode of vibration of the member

According to BV Bro (Banverket, 2006) it is enough to use the 30 Hz criteria

### 3.4 Structural damping

The acceleration response at resonance is highly dependent on the choice of structural damping. Therefore a lower bound value for the structural damping is suggested by Eurocode 1, CEN (2003) and the Swedish code BV Bro (Banverket, 2006). The suggested values for structural damping are shown in Table 3.3.

Table $3.3 \quad$ Values of damping for different spans and type of bridges.

| Bridge type | $\zeta$ Lower limit of damping [\%] |  |
| :--- | :--- | :--- |
|  | Span $L<20 \mathrm{~m}$ | Span $L \geq 20 \mathrm{~m}$ |
| Steel and composite | $\zeta=0.5+0.125(20-\mathrm{L})$ | $\zeta=0.5$ |
| Prestressed concrete | $\zeta=1.0+0.07(20-\mathrm{L})$ | $\zeta=1.0$ |
| Reinforced concrete | $\zeta=1.5+0.07(20-\mathrm{L})$ | $\zeta=1.5$ |

## 4 ANALYSIS IN 2D

Dynamic analysis of railway bridges is usually carried out using mode-superposition and therefore this thesis focuses solely on this method of calculation. The determination of the dynamic response in all structures when mode superposition is to be utilized is done in two steps. First the eigenfrequencies and the corresponding eigenvectors need to be determined. Secondly the response in time is determined. For simple structures this can be done analytically but for more complex structures (or complex loads) there is often a need to solve numerically.

One of the aims of this thesis is to find general guidelines for the design of railway bridges with regard to the dynamic aspects of the design process. To do that, both eigenfrequencies and the response in time have to be examined. The eigenfrequencies in relation to the load frequency applied on the bridge will affect the dynamic response behavior. The actual dynamic response will however also be affected by the configuration of geometric and material parameters of the bridge. The method applied in this thesis is to examine the effect from individual variation of material and geometric parameters to find how these affect the dynamic response. The dynamic response will throughout this thesis be referred to as the acceleration response. The reason is basically that the acceleration is the only examined outcome from the dynamic analysis performed in this thesis, since the demands on vertical acceleration govern the design process of railway bridges when considering the dynamic aspects.

A model that transforms a railway bridge seen in 2D into an SDOF model has been developed in the previous master thesis Ekström and Kieri (2007) and De Leon and Lasn (2008) carried out at Reinertsen Sverige AB. This model was shown to give accurate results for single-span bridges. The theory behind the SDOF model can be used to increase the understanding of the dynamic behavior of railway bridges. Therefore it is of interest to examine if the model is applicable on multi-span bridges. If the model can be proven to give sufficiently accurate results for multi-span bridges, it can also be a useful tool for calculating the dynamic response since it greatly reduces calculation time compared to a finite element model (FE model).

Section 4.1 will treat geometric definitions of span relation parameters and the crosssection geometry which has been kept constant for all calculations in 2D.

In Chapter 2 it was shown how the eigenfrequencies of a single-span bridge can be calculated analytically. The eigenfrequencies of multi-span bridges are examined in Section 4.2. The eigenfrequencies' complex dependence of geometric parameters is shown graphically, and also expressions for calculating the first eigenfrequencies are created empirically. These expressions allows for better examination of the acceleration response later on.

The load frequency and resonance phenomenon in railway bridges are treated in Section 4.3. It is shown that each HSLM load exerts one fundamental load frequency that dominates the dynamic response. It is also shown how the critical train speed (for which resonance in the bridge occurs) can be determined.

Section 4.4 treats the transformation of the bridge into an SDOF model. Theory behind the transformation is presented and accuracy of the model is tested.

After examining the eigenfrequencies for multi-span bridges and showing how the load frequency affect the dynamic response, a parameter study is presented in Section 4.4 treating the amplitude of the dynamic response. Geometric and material parameters' effect on the eigenfrequencies will influence the critical speed at which resonance occurs. As the examination of eigenfrequencies aim to show how the critical speeds can be determined in a simplified way, the parameter study made in Section 4.4 aims to show the acceleration amplitude at these resonance speeds depends on material and geometric parameters.

It will be shown that the geometric parameters affect the acceleration response in a rather complicated way, practically eliminating the possibility to determine single equations to express this behavior. Some simple dependencies are found though. All observations and conclusion from the previous sections are gathered in Section 4.6 and used to create a graphical calculation tool for some common railway bridge configurations. This tool can be used to determine the exact design acceleration for known material and geometric parameters, but also as a guideline of what parameter configuration that should be used to obtain a satisfying dynamic response in the bridge.

### 4.1 Geometric definitions

The calculations presented in this Chapter aim partly to show the effect from varying geometric and material parameters of a railway bridge. The cross-section however have been kept constant for simplicity since the effect from varying bending stiffness and bridge mass is better examined by variation of young's modulus and the density respectively.

The geometric parameters used in the 2D analysis will be presented in this section. Parameters that describe the span relations for two- and three-span bridges will also been defined.

### 4.1.1 Cross section

For all the examinations made in 2D analysis the same geometric parameters have been used, see Figure 4.1. The section chosen resembles the section of a portal frame bridge, with a total length of 9.5 m , designed by Reinertsen Sweden AB. The reason why a section from a real bridge design was chosen was to get an idea of how large accelerations and eigenfrequencies would get for different combinations of geometric and material parameters. For the calculations in 2D the section is treated as a rectangular cross-section, and hence the edge beams are disregarded.

(m)

Figure 4.1
Cross section used for all analysis in this Chapter.

Since the design code of the Swedish railway administration (BV Bro, Banverket 2006) has a limit of 30 Hz for eigenmodes that needs to be considered there is a special interest for the size of the eigenfrequencies. The eigenfrequencies also have a significant influence of how large the design acceleration will be, as will be shown later in this thesis.

### 4.1.2 Definitions of span relation parameters

The additional geometric parameters for a multi-span bridge in comparison to a single span bridge are the relations between the span lengths. For a two-span bridge one extra parameter is required and two extra parameters are required for a three-span bridge. The span relation for a two-span bridge has been defined as the length ratio between the second and the first span, i.e.

$$
\begin{equation*}
\mu=\frac{L_{2}}{L_{1}} \tag{4.1}
\end{equation*}
$$

where $L_{1}$ and $L_{2}$ are the lengths of the first and second span respectively. Figure 4.2 shows how the parameter $\mu$ can be used to describe the geometry of a two-span bridge.


Figure $4.2 \quad$ Explanatory figure of how the geometry of a two-span bridge can be described using the parameter $\mu$ defined in equation (4.1).

Two parameters are required to describe the geometry of a three-span bridge. These parameters are here defined as

$$
\begin{align*}
& \eta=L_{2} /\left(L_{1}+L_{3}\right)  \tag{4.2}\\
& \kappa=L_{3} / L_{1} \tag{4.3}
\end{align*}
$$

Figure 4.3 shows how these parameters can be used to describe the geometry of a three-span bridge.


Figure 4.3 Explanatory figure of how the geometry of a three-span bridge can be described using the parameters $\eta$ and $\kappa$ defined in equations (4.2) and (4.3).

### 4.2 Eigenfrequencies for multi-span bridges

It is possible to derive an expression that gives the exact value of the eigenfrequencies of any degree for a single span bridge, as was shown in Chapter 2. A bridge that has two or more spans contributes to a lot more complexity in this derivation and to the knowledge of the authors it is not possible to derive an explicit expression for the frequencies analytically.

In Appendix A there is a derivation similar to that of a single-span bridge for deriving the eigenfrequency of a two-span bridge made by the authors. In litterateur you can find similar derivations that lead to similar expressions to the one derived in Appendix A, see for example Gorman (1975). But the result will always be a matrix which determinant should equal zero, resulting in a complex trigonometric expression where an explicit formulation for the eigenfrequencies is impossible to derive.

This section will treat the determination of eigenfrequencies for multi-span bridges under the variation of geometric and material parameters. As eigenfrequencies cannot be determined analytically they have instead been examined graphically. The first eigenfrequency is of special interest since it often governs the dynamic response in railway bridges. Therefore an expression to easily determine the first eigenfrequency for a two- and three-span bridge has been derived.

It is important to remember the design rule set by the Swedish railway administration, stating that only frequencies below 30 Hz need to be considered when designing a railway bridge against the dynamic response criteria. Often only a few eigenmodes need to be considered in dynamic analysis of railway bridges because of this rule, and
for many bridges only the first. Therefore only the first eigenfrequencies have been examined in the following section.

The calculation of eigenfrequencies under the variation of material and geometric parameters has been made using a Matlab program presented in Appendix D made by the authors. Matlab have been used instead of commercial FE software since it better handles iterative calculations concerning the variation of geometric parameters. However ADINA has been used to verify the Matlab program, see Appendix D.

### 4.2.1 Two-span bridges

As for single-span bridges the frequency of a two span bridge will depend on the bending stiffness $E I$, the mass per meter length $m$ and the total length of the bridge $L_{\text {tot }}$. The addition to these parameters for a two span bridge is the ratio between the second and first span $\mu$, which was defined in Section 4.1.

In Appendix A, where an attempt to derive the eigenfrequencies analytically is made, it is shown that a complicated expression is obtained when determining an explicit expression for the eigenfrequencies' dependence on bridge parameters. This complicated expression however depends solely on the geometric parameters $L_{\text {tot }}$ and $\mu$. The eigenvalues' dependence on material parameters is defined in the same way as for single span bridges, see equation (2.69). This means that irrespective of the complicated dependence to geometric parameters the eigenfrequencies of a multi-span bridge depends in the same simple manner on material parameters as a single-span bridge.

It can also be shown that the same dependency, as for single-span bridges, exists for the total bridge length, see Figure 4.4 and Figure 4.5.


Figure 4.4 Eigenfrequencies for the first eigenmode for a two-span bridge with $\mu=2, E=30 G P a$, $\rho=3000 \mathrm{~kg} / \mathrm{m}^{3}$.


Figure 4.5 Eigenfrequencies for the second eigenmode for a two-span bridge with $\mu=2$, $E=30 G P a, \rho=3000 \mathrm{~kg} / \mathrm{m}^{3}$.

Figure 4.4 and Figure 4.5 show that both first and second eigenfrequency has a linear dependence against the inverse of the total bridge length in square. This means that the eigenfrequencies for a two-span bridge in a 2D analysis can be determined as

$$
\begin{equation*}
f_{i}=\frac{g_{i}(\mu)}{L_{t o t}{ }^{2}} \cdot \sqrt{\frac{E I}{m}} \tag{4.4}
\end{equation*}
$$

where $g_{i}(\mu)$ is a different function for each eigenfrequency that depends on the span ratio. These functions will in the remainder of this thesis be referred to as eigenfrequency functions. Figure 4.6 shows how these functions vary for the first five eigenfrequencies. Note that the function is dimensionless and independent of all geometric and material parameters with the exception of $\mu$.


Figure 4.6
Eigenfrequency functions for the first five frequencies for a two-span bridge.
Examining Figure 4.6 can help understand the difficulties in deriving an analytical expression for the eigenfrequencies. As can be seen they are all affected differently by
the change of $\mu$. Also the dependency for each individual frequency on $\mu$ is complex. This complexity can be explained by the eigenvectors corresponding to each frequency changing in shape. The shapes for the eigenvectors for the first, second and third eigenfrequency for different values of $\mu$ are shown in Figure 4.7- Figure 4.9. From Figure 4.7- Figure 4.9 it is shown that only the first eigenvector keeps a similar form when the ratio between the span lengths is altered, hence the more straight curve in Figure 4.6.


Figure 4.7
The first eigenmode for a two-span bridge with a total length of 15 m . The mode is normalized to get the modal mass equal to one.


Figure $4.8 \quad$ The second eigenmode for a two-span bridge with a total length of 15 m . The mode is normalized to get the modal mass equal to one.


Figure $4.9 \quad$ The third eigenmode for a two-span bridge with a total length of 15 m . The mode is normalized to get the modal mass equal to one.

It is of great value for the remaining examinations made in this thesis to have a simple expression for the first eigenfrequency. Therefore efforts have been made to create such an expression. The "curve fitting" toolbox in Matlab was used for this purpose. To get a more simple expression different equation types were fitted for different intervals of $\mu$. The equations shown in Table 4.1 are proposed to describe the first eigenfrequency for a two-span bridge.

Table 4.1 Proposed eigenfrequency function for determining the first eigenfrequency for a twospan bridge.

| $\mu$ | Eigenfrequency function, $g(\mu)$ | $\mathrm{R}^{2}$-value |
| :--- | :--- | :--- |
| $1-1.2$ | $2 \cdot \pi \cdot \sin (\pi / 2 \cdot \mu)$ | 0.991 |
| $1.2-3$ | $8.9 \cdot e^{-\mu}+3.34 \cdot e^{-0.00183 \cdot \mu}$ | 0.999 |

It should be noted that there is no derivation behind the expressions in Table 4.1, they have simply been chosen for their simplicity in comparison to other expressions. Figure 4.10 shows the curves fitted against the true acceleration.


Figure $4.10 \quad$ Chosen eigenfrequency functions $g$ for two-span bridges plotted against the true first frequency.

Three different sets of material parameters have been chosen to show the accuracy of the expressions. Table 4.2 and Table 4.3 show these parameter sets and a comparison of eigenfrequencies calculated using the simple method presented in this section and a FE analysis.

Table $4.2 \quad$ Parameter sets used for different verifications throughout this thesis.

| Parameter | Set 1 | Set 2 | Set 3 |
| :--- | :---: | :---: | :---: |
| $E[\mathrm{GPa}]$ | 30 | 40 | 50 |
| $\rho\left[\mathrm{~kg} / \mathrm{m}^{3}\right]$ | 2400 | 2000 | 2800 |
| $L_{\text {tot }}[\mathrm{m}]$ | 25 | 20 | 15 |
| $\mu[-]$ | 1.1 | 2.5 | 1.5 |

Table 4.3 Eigenfrequencies calculated using equation (4.4) and eigenfrequency functions in Table 4.1 compared to eigenfrequencies from a FE analysis.

|  | Set 1 | Set 2 | Set 3 |
| :--- | :---: | :---: | :---: |
| FE analysis | 6.07 | 7.85 | 17.31 |
| Equations | 6.08 | 7.85 | 17.29 |

It is concluded that the frequency functions in Table 4.1 together with equation (4.4) gives satisfactory accuracy for determining the first eigenfrequency for a two-span bridge.

### 4.2.2 Three-span bridges

A three-span bridge have the same dependence on the bending stiffness $E I$, the mass per meter $m$ and the total length of the bridge $L_{\text {tot }}$, with the same motivation as described in Section 4.2.1 for a two-span bridge. Two extra geometric parameters are however required to describe the relation of the three spans. In this thesis they are called $\eta$ and $\kappa$ and are defined in Section 4.1.

Since a two-span bridge only depends on one geometric parameter in a complex way it is possible to find expressions for describing the eigenfrequency functions for any combination of geometric and material parameters for this type of bridge. It gets more complex for a three-span bridge where two geometric parameters ( $\eta$ and $\kappa$ ) have a complex effect on the eigenfrequencies, see Figure 4.11.


Figure 4.11
Eigenfrequency function for the first eigenfrequency for the variation of $\kappa$ and $\eta$.
Figure 4.11 shows that the influence of the geometric parameter $\kappa$ changes depending on the value of $\eta$. This means that it is not possible to formulate two separate equations for describing the eigenfrequencies' dependence on the two variables. Instead one equation depending on both variables is required, making it much harder to find a suitable expression.

It is rather common for three-span bridges to have $\kappa$ equal to one, meaning the first and third span are equally sized while the middle span may vary in size. Since complexity in finding an expression for the frequencies depend on both $\eta$ and $\kappa$ a limitation has been made to bridges where $\kappa=1.0$. This type of bridge has a similar appearance of the eigenfrequencies for the variation of $\eta$ as for a two span bridge, see Figure 4.12.


Figure 4.12 Eigenfrequency functions for the first seven eigenfrequencies for a three-span bridge with $\kappa=1.0$.

As for a two-span bridge it is of great value for further examinations to have a simplified expression for calculating only the first eigenfrequency for variations in $\mu$. Again the Matlab curve fitting toolbox was used for this purpose, and the expressions in Table 4.4 are proposed for different intervals of $\mu$.

Table 4.4 Proposed expressions for the eigenfrequency function for the first eigenfrequency of a three-span bridge with $\kappa=1.0$.

| Mode | $\eta=L 2 /(L 1+L 3)[-]$ | Eigenfrequency function, $g(\eta)$ | $\mathrm{R}^{2}$-value |
| :--- | :---: | :---: | :---: |
| First mode | $0.5-0.8$ | $1.6 \cdot \sin (7.2 \cdot \eta-1.95)+12.55$ | 0.999 |
| First mode | $0.8-2.5$ | $2 \cdot p i \cdot \eta^{-1.137}+3.511$ | 0.999 |

Figure 4.13 shows the curves fitted against the true acceleration.


Figure 4.13
Chosen eigenfrequency functions $g$ for two-span bridges plotted against the true first frequency.

A verification of the accuracy of the expressions in Table 4.4 has been made using three different parameter sets which are defined in Table 4.5. Table 4.6 shows the results from this verification.

Table 4.5
Parameter sets used for different verifications throughout the thesis

| Parameter | Set 1 | Set 2 | Set 3 |
| :--- | :---: | :---: | :---: |
| $E[\mathrm{GPa}]$ | 30 | 40 | 50 |
| $M\left[\mathrm{~kg} / \mathrm{m}^{3}\right]$ | 2400 | 2000 | 2800 |
| $L_{\text {tot }}[\mathrm{m}]$ | 30 | 25 | 20 |
| $\mu[-]$ | 0.6 | 1.5 | 0.9 |

Table 4.6 Eigenfrequencies calculated using equation (4.4) and eigenfrequency functions in Table 4.4 compared to eigenfrequencies from a FE analysis.

|  | First eigenfrequency [Hz] |  |  |
| :--- | :---: | :---: | :---: |
|  | Set 1 | Set 2 | Set 3 |
| FE analysis | 9.301 | 9.254 | 19.402 |
| Equations | 9.298 | 9.262 | 19.385 |

### 4.2.3 Comparison to single-span bridges

Single-span railway bridges are very common and their first eigenfrequency can as mentioned be calculated analytically. Hence it may be of interest to know the increase in eigenfrequency that the additional spans of a multi-span bridge contribute with compared to a single-span bridge corresponding to the largest span length. The eigenfrequency of a multi-span bridge can alternatively be calculated as

$$
\begin{equation*}
f_{i}=U \cdot f_{i, \text { single }} \tag{4.5}
\end{equation*}
$$

where $f_{i, \text { single }}$ is the i:th eigenfrequency of a single-span bridge corresponding to the largest span length of the multi-span bridge and

$$
\begin{equation*}
U=\frac{2 \cdot g_{i}(\mu) \cdot(1+\mu)^{2}}{\pi \cdot \mu^{2}} \tag{4.6}
\end{equation*}
$$

for a two-span bridge and

$$
\begin{equation*}
U=\frac{2 \cdot g_{i}(\eta) \cdot(1+\eta)^{2}}{\pi \cdot \eta^{2}} \tag{4.7}
\end{equation*}
$$

for a three-span bridge. Figure 4.14 shows the appearance of $U$ for calculation of the first eigenfrequency of a two- and three-span bridge.


Figure $4.14 \quad$ Appearance of the factor $U$ for determining the first eigenfrequency for a two- and three span bridge through the eigenfrequency of a single-span bridge corresponding to the largest span length.

### 4.2.4 Remarks

By examining the derivation of eigenfrequencies for single- and multi-span bridges it was seen that the eigenfrequencies for a multi-span bridge depend on material parameters in the same way as a single-span bridge. By graphic examination it was seen that this also is true for the total length of the bridge. One important feature to note is that all frequencies have the same dependence on these parameters. This means the eigenfrequencies have the same relation to each other for any combination of material parameters and total bridge length. This is important in later calculations since it allows searching for simpler dependencies that the acceleration response has on these parameters. However the geometric span relation parameters affect the eigenfrequencies differently for all multi-span bridges. This means that the dynamic response behavior will be affected in a different way when span relations are altered depending on which eigenmode that is examined. It is therefore not possible to define exactly how the dynamic response depends on the span relation parameters since each eigenmode has its own effect on the dynamic response, and the number of eigenmodes in a continuous structure is infinitely large.

### 4.3 Load frequency and resonance effects

According to BV Bro (Banverket 2006) all railway bridges with a total length above 7 $m$ have to be designed to withstand the HSLM-A train loads in terms of maximum allowed vertical acceleration in the superstructure of the bridge, see Chapter 3. As shown in Chapter 3 all HSLM-A loads have different load parameters in the term of coach length $D$, bogie axel length $d$ and the number of coaches. In this section it is discussed how the train loads will affect a railway bridge is presented. The main load frequency is defined and it is shown how it in relation to the eigenfrequencies of a bridge affects the dynamic response due to resonance.

To understand the meaning of load frequency of a train load it is convenient to start with a single point load moving over a single-span bridge with a fixed velocity, see Figure 4.15.


Figure 4.15 A single-span bridge subjected to one point load. (a): constant load moving over the bridge with constant velocity. (b): point load that varies in time acting in the middle of the bridge.

Figure 4.15 implies that a point load moving over a single-span bridge can be transformed from a load varying in the space domain to a stationary load varying in the time domain. The transformed load in Figure 4.15 will hence vary in time and if
several loads with equal distance moves over the bridge there will be a load effect on the bridge varying with a certain frequency affected by the speed of the loads and the distance between them. As an example let us consider the static analysis of the bridge in Figure 4.15. The influence line for the mid-span deflection caused by one point load as it moves over the bridge is very close to a sinusoidal curve. If several influence lines are superpositioned in the time domain a resulting influence of the loads is received, see Figure 4.16.


Figure 4.16
The influence line for static displacement in a single-span bridge subjected to several equally distanced point loads travelling over the bridge with constant speed.

The static response shown by Figure 4.16 gives an idea of how a bridge will be affected by moving point loads, even though the full complexity in a dynamic analysis is not considered. It should be noted that the appearance of the total influence line differ as the load distance is altered or bridge spans are added. For more information about the reliability/possibility to describe load effect on the dynamic response using a static model, see Section 4.4.

The train load models HSLM-A are all a combination of point loads traveling with different distances in between, namely the coach length $D$ and the boogie axel length $d$. However, the coach length $D$ can be considered as the "main" distance between the loads, as it serves as the distance between the resultants of the boogie axle loads of the train, see Figure 3.1. In Figure 4.16 it can been seen that the time between peaks in the total load influence is equal to the time between the traveling loads. As the frequency is the inverse of this time it can then be concluded that the "main" frequency of a HSLM-A train load can be defined as

$$
\begin{equation*}
\Omega=\frac{v}{D} \tag{4.8}
\end{equation*}
$$

where

$$
\begin{aligned}
& \Omega=\text { Load frequency } \\
& D=\text { Coach length of considered HSLM-A load } \\
& v=\text { Train velocity }
\end{aligned}
$$

Resonance will occur when the load frequency coincides with one of the eigenfrequencies of any structure, see Chapter 2. However, for railway bridges resonance will also occur when an eigenfrequency coincides with any multiple of the load frequency, i.e.

$$
\begin{equation*}
\Omega=\frac{f_{i}}{k} \tag{4.9}
\end{equation*}
$$

where

$$
\begin{aligned}
f_{i} & =\text { Eigenfrequency of the i:th degree } \\
k & =1,2,3 \ldots
\end{aligned}
$$

The behavior behind equation (4.9) can be understood by looking at a spring-weight SDOF system without damping. Say the system is vibrating with its eigenfrequency. If we tap the weight (put a small impulse load on it) every time it is in its highest position we will increase the amplitude of the vibration every cycle and the amplitude goes to infinity with time, see Figure 4.17. If we tap the weight every second time it is in its highest position the amplitude also goes to infinity in time but it increases slower. These two examples correspond to adding a load with a frequency coinciding with the eigenfrequency and the eigenfrequency divided by two. In a real structure the first of these most often gives the largest dynamic responses since the structure has less time to damp out the vibrations.


Figure 4.17
Illustrative figure of how the displacement increases for a SDOF model when an impulse load is applied with a frequency corresponding to the systems eigenfrequency and half the eigenfrequency in (a) and (b) respectively

The ratio between the load frequency and the first eigenfrequency is defined as

$$
\begin{equation*}
\beta=\frac{\Omega}{f_{1}} \tag{4.10}
\end{equation*}
$$

This parameter will be frequently used in this thesis since it can be used to describe all bridge parameters and the train velocities' effect on the maximum acceleration response.

FE analyses have been performed on a two-span bridge to show how the dynamic response in a railway bridge is dominated by resonance phenomenon. Three different data sets have been used, see Table 4.2. In the analyses mode superposition is used and limited to the first two eigenmodes. The examined result is the maximum acceleration for the variation of train velocity, see Figure 4.18.


Figure 4.18
Maximum acceleration response for the variation of train velocity using the three parameter sets defined in Table 4.2, two eigenmodes and a modal damping of $2 \%$. The train load HSLM-A1 is used for all three parameter sets.

It can be seen in Figure 4.18 that resonance effects occur at $\beta=1$ and $1 / 2$ for all three parameter sets and thus confirming the correctness of equation (4.9).

An additional parameter denoted the critical speed is defined before proceeding to the following sections.

$$
\begin{equation*}
v_{c r}=D \cdot f_{1} \tag{4.11}
\end{equation*}
$$

Where:

$$
\begin{aligned}
v_{c r}= & \text { The critical speed for a specific bridge configuration and } \\
& \text { HSLM-A load }
\end{aligned}
$$

The critical speed is the speed resulting in $\beta=1.0$ for a given coach length $D$ and consideration to the first eigenfrequency. In reality this speed is never reached as it is unadvisable considering dynamic aspects of railway bridge design. Figure 4.18 gives an understanding of how much larger the accelerations become if the speed is
reached. It serves nevertheless as a good reference value and it will be used frequently in the remainder of the thesis.

By combining equation (4.8), (4.9) and (4.11) the resonance speed can be shown to be

$$
\begin{equation*}
v_{\text {resonance }}=\frac{v_{c r}}{k} \tag{4.12}
\end{equation*}
$$

where

$$
\begin{aligned}
v_{\text {resonance }}= & \text { train velocity at which resonance effects occur in a } \\
& \text { railway bridge }
\end{aligned}
$$

Finally by combining equations (4.4), (4.11) and (4.12) an analytical expression for calculating the resonance speeds by hand is obtained as

$$
\begin{equation*}
v_{\text {resonance }}=\frac{D}{k} \cdot \frac{g_{i}}{L_{\text {tot }}{ }^{2}} \cdot \sqrt{\frac{E I}{m}} \tag{4.13}
\end{equation*}
$$

### 4.4 The SDOF model

A model meant for describing the dynamic response of multi-span bridges consisting of only one degree of freedom has been established. The purpose of creating such a model is partly because it describes the complex dynamic response in a simple way which may increase the understanding of this complexity, but also because it greatly decreases calculation time and gives an analytical equation of the response that could be used for finding guidelines in the design process of railway bridges.

In this section the theory behind the creation of the SDOF model will be explained. This is followed by a verification of the SDOF model's accuracy.

### 4.4.1 Theory

The examination of a SDOF model that resembles a railway bridge is a continuation from previous master carried out at Reinertsen Sweden AB, namely Ekström and Kieri (2007) and De Leon and Lasn (2008).

In Ekström and Kieri (2007) the transformation of a single-span railway bridge was made using transformation factors which was first introduced by Nyström (2006), also a master thesis carried out at Reinertsen Sweden AB. The transformation factor approach creates a transformation factor for stiffness, mass and load respectively. These factors are defined so that kinetic energy stored in the beam and SDOF system becomes equal, the moved mass and the load exerts the same amount of work. The transformation factors are based on the assumption that the beam deforms as a sinusoidal curve.

De Leon and Lasn (2008) compared the transformation factor approach to an additional method of transformation, called the force scaling approach. The force scaling approach requires considerably less calculations and is also easier to apply on
more complex structures. The method was proven to give accurate results for a singlespan bridge.

One of the aims of this thesis is to examine if it is possible to transform also a multispan bridge into a SDOF model. Because of the complexity in multi-span bridges the force scaling approach has hence been chosen for this task, and its theory will be presented in this section. The interested reader is referred to Nyström (2006) for theory on the transformation factor approach.

A SDOF model has the equation of motion shown in equation (2.2). To establish this model the stiffness $K$, mass $M$, damping $C$ and load $p(t)$ needs to be determined. The force scaling approach has the following calculation procedure:

- Chose an arbitrary mass of the SDOF model
- Calculate the stiffness of the SDOF model so the first frequency of the SDOF model and the railway bridge becomes equal
- Determine the damping
- Scale the load acting on the bridge to use on the SDOF model so that the static deflection of the bridge and SDOF model becomes equal

Let us go through the four steps in more detail. First an arbitrary mass of the SDOF model is chosen as

$$
\begin{equation*}
M=T \cdot M_{\text {bridge }}, \text { where } \mathrm{T}>0 \tag{4.14}
\end{equation*}
$$

In Section 4.3 it was shown that the response in railway bridges is governed by the ratio between load frequency and first eigenfrequency. It is hence important that the SDOF model has a frequency that corresponds to the first eigenfrequency of the bridge. The fact that an SDOF model only has one eigenfrequency limits the model since only the first eigenfrequency of the bridge can be regarded. Resonance effects that occur with regard to higher eigenmodes are not given consideration in the model. By combining the expression for the eigenfrequency in equation (2.5) and equation (4.4) an expression for determining the stiffness is obtained as

$$
\begin{equation*}
K=M \cdot \frac{f_{i}(\mu)}{L_{t o t}{ }^{2}} \cdot \sqrt{\frac{E I}{m}} \tag{4.15}
\end{equation*}
$$

The SDOF model is made to resemble the dynamic response using the first eigenmode. The damping should therefore be chosen to resemble the modal damping of the first eigenmode. Equation (4.16) shows how the damping $C$ can be calculated from a known modal damping.

$$
\begin{equation*}
C=2 \cdot \zeta \cdot M \cdot \omega_{s d o f} \tag{4.16}
\end{equation*}
$$

where

$$
\omega_{s d o f}=\sqrt{\frac{K}{M}}
$$

Finally the load acting on the bridge is scaled so that the static deflection of the bridge and SDOF model becomes equal. The static displacement of the SDOF model is defined as

$$
\begin{equation*}
u(t)_{s d o f}=\frac{p(t)_{\text {sdof }}}{K} \tag{4.17}
\end{equation*}
$$

Setting the static displacement in equation (4.17) equal to the static displacement of the bridge gives an expression for determining the load acting on the SDOF model as

$$
\begin{equation*}
p(t)_{\text {sdof }}=K \cdot u(t)_{\text {bridge }} \tag{4.18}
\end{equation*}
$$

The force scaling approach has its name since equation (4.18) re-written forms a relation between the load acting on the bridge and on the SDOF model equal to $\mathrm{K} / \mathrm{K}_{\text {bridge }}$ which in De Leon and Lasn (2008) has been called the scaling factor.

A point in which the largest accelerations are assumed to occur in the bridge needs to be chosen before the deflection of the bridge can be calculated. It is assumed in this thesis that the largest acceleration will occur in the middle of the largest span. The static deflection in time is calculated by using influence lines for the deflection in the critical point with consideration to one point load. The influence lines are based on the coordinate of the point load which can be replaced with velocity of the point loads times the time to get the static deflection in time. By superpositioning the influence of point loads in a correct manner a HSLM-A train load can be resembled. The added complications with applying the force scaling approach on a multi-span bridge compared to on a single-span bridge, as has been made in De Leon and Lasn (2008), lies solely in the calculation of the static deflection of the bridge.

### 4.4.2 Verification

The whole concept of the SDOF model is based on the limit of required eigenmodes in BV Bro Banverket (2006). As described in Chapter 3 this limit is 30 Hz . It means that the eigenmodes that has an eigenfrequency above 30 Hz can be excluded when calculating the dynamic response.

The SDOF model created is meant for calculating the dynamic accelerations. These accelerations are strongly affected by higher eigenmodes, as shown by Figure 4.19. The only time the SDOF model is capable of describing a more accurate dynamic acceleration response is when only the first eigenmode has to be considered. However, this is often the case for multi-span bridges.


Figure $4.19 \quad$ Comparison of the maximum acceleration response in a two-span bridge subjected to a point load of 100 kN moving with $v=50 \mathrm{~m} / \mathrm{s}$ using different number of eigenmodes. The used parameters are $E=40 \mathrm{GPa}, \rho=2400 \mathrm{~kg} / \mathrm{m}^{3}, L=15 \mathrm{~m}, \mu=2, c=0.02$.

A SDOF model based on the theory presented in Section 4.4.1 has been established and its Matlab code can be found in Appendix D. The SDOF model's accuracy have been verified against a FE program that also can be found in Appendix D. In the verification only the first eigenmode was used in the FE analysis since this corresponds to what the SDOF model describes. Table 4.7 shows two sets of parameters that have been used in the verification. Figure 4.20 and Figure 4.21 shows the acceleration response in time for the SDOF model and the FE analysis using data set 1 in Table 4.7.

Table 4.7 Two parameter sets used in the verification of the SDOF models accuracy. Both parameter sets concern a two-span bridge.

| Bridge <br> type | $E[\mathrm{GPa}]$ | $\rho\left[\mathrm{kg} / \mathrm{m}^{3}\right]$ | $L[\mathrm{~m}]$ | $\mu[-]$ | $c[-]$ | HSLM-A | $v[\mathrm{~m} / \mathrm{s}]$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Set 1 | 40 | 2400 | 15 | 2.0 | 0.02 | A1 | 50 |
| Set 2 | 40 | 2400 | 15 | 2.0 | 0.02 | A4 | 75 |



Figure $4.20 \quad$ Acceleration response in the middle of the second span using data set 1 in Table 4.7.


Figure $4.21 \quad$ Acceleration response from Figure 4.20 during the first second.
Figure 4.22 and Figure 4.23 shows the acceleration response in time for the SDOF model and the FE analysis using data set 2 in Table 4.7.


Figure 4.22
Acceleration response in the middle of the second span using data set 2 in Table 4.7.


Figure 4.23
Acceleration response from Figure 4.22 during the 2: nd second.
For the first data set in Figure 4.20 and Figure 4.21there is a close resemblance between the SDOF model and the FE analysis. Both the appearance and the amplitude of the acceleration response are similar. For the second data set Figure 4.22 and Figure 4.23 there is however some lack in accuracy when it comes the amplitude of the accelerations, even though the appearance is accurate.

A more extensive comparison of the acceleration amplitude has been made by examining the variation of maximum acceleration in time for different train velocities, see Figure 4.24 and Figure 4.25. A comparison between the acceleration response in the middle of the largest span (in this case the second span) and the maximum acceleration response in the whole bridge is also made in Figure 4.24 and Figure 4.25.


Figure 4.24 Maximum acceleration response using data set 1 in Table 4.7 with variation in the train velocity.


Figure $4.25 \quad$ Maximum acceleration response using data set 2 in Table 4.7 with variation in the train velocity.

When higher modes are included the maximum acceleration is likely to occur somewhat off the center of the largest span. When only one mode is included however Figure 4.24 and Figure 4.25 show that the maximum acceleration occurs in the middle of the largest span with just a small error. This is good for the SDOF model since it is based on the maximum acceleration appearing in the middle of the largest span. This makes the acceleration behavior in the SDOF model more correct. Figure 4.24 and Figure 4.25 do however also show that the SDOF model lack accuracy and also that this lack of accuracy is not consist for the variation of train velocity.

In the following section calculations have been made to capture different material and geometric parameters dependence on the acceleration amplitude. For the possibility of finding such dependencies it is necessary to use a calculation tool that can show if the dependencies are consistent for the variation of train velocity. It was stated in the beginning of this Chapter that the SDOF model could be used to decrease the calculation time for the upcoming calculations if proven accurate enough. But since the accuracy varies depending on train velocity the SDOF model has not been used further in this thesis.

It can however be stated that the accuracy of the SDOF model is remarkably good considering the models simplicity even for multi-span bridges.

### 4.5 Amplitude of the acceleration response

By calculating the eigenfrequencies of a railway bridge, as shown in Section 4.2, it is possible to decide the critical speeds at which resonance occur for a certain train load. In Section 4.3 it was also shown that resonance effects occurs for multiples of the critical speed for railway bridges and how the resonance speeds can be calculated.

The equations in Sections 4.2 and 4.3 give together an exact formulation of how the resonance speed is calculated. They do not however give any guidance to how large the dynamic response will be at this resonance speed. As this thesis partly aims to give guidelines for how to estimate the dynamic response it is required to further examine the variation in acceleration amplitude of the dynamic response.
In this section the results from a parameter study made on two- and three span bridges is presented. In this parameter study the amplitude of the dynamic response in terms of acceleration has been examined at different resonance speeds for the variation of both material and geometric parameters. Section 4.5 .3 will present the results from the parameter study for the individual variation of each material and geometric parameter, respectively. A discussion and the conclusions made concerning the amplitude of the dynamic acceleration response are presented in Section 4.5.4.

The dynamic response in the parameter study has been calculated using a Matlab FE program developed within this thesis by the authors, see Appendix D. For details about meshing, element choice, load application and time integration the reader is referred to Appendix D. In comparison to commercial FE software the variation of geometric parameters is substantially easier to handle in Matlab. It is also easier to apply train loads and variation of train velocity in Matlab.
Equation (4.13) has been used for calculating the resonance speeds both in the Matlab program and when handling the data from it. In the FE calculations mode superposition was used utilizing only the first eigenmode to save calculation time.

Parameter behaviors that differ for different eigenmodes are beyond the time limit of this thesis and therefore the use of additional eigenmodes is unnecessary.

In this section a selection of figures are presented to give the reader a comprehensible presentation of the results from the parameter study. For the interested reader a more extensive gathering of figures from all the analysis made can be found in Appendix B.

### 4.5.1 Method

In this section the method behind the parameter study will be presented. The section is meant to clarify what calculations that have been performed and how the data for the figures that will be presented have been extracted.

The time response of the vertical acceleration when a train runs over a bridge is complex. However the shape of this response is not of interest in the process of designing a railway bridge. Only the maximum acceleration in time is of interest. Therefore there is no necessity in examining the actual response in time, but rather only the maximum acceleration. When designing a railway bridge a range of train velocities need to be considered however, see Chapter 3. Changing the train velocity may result in a completely different acceleration response in time and hence also a different maximum acceleration.

The process of creating the Matlab programs that have been used for the parameter study can be divided in steps. First a program that calculates the response in time for a fixed train velocity and parameter set was created and verified against ADINA. Then this program was modified to calculate the response in time for every train velocity of interest and from these extracting the maximum acceleration response, see Figure 4.26.


Figure $4.26 \quad$ Illustrative figure of how the acceleration response in time has been gathered in a $v$ a plot.

The method applied in the parameter study has been to first calculate the maximum acceleration response in time for the variation of train velocity with a certain train load and set of material and geometric parameters, as shown by Figure 4.26 procedure is then repeated with the independent variation of one variable. The variation in amplitude of the acceleration response for changing that variable can then be studied by examining the amplitude for fixed values of $\beta$. Since the acceleration response at resonance often governs the determination of design acceleration in real life projects the $\beta$ values $0.5,0.33,0.25$ and 0.2 have been examined which corresponds to resonance speeds. Figure 4.27 shows how the data at a constant value of $\beta$ have been extracted to show the accelerations dependence upon a variable of interest.


Figure $4.27 \quad$ Illustrative figure of how the acceleration at a $\boldsymbol{\beta}$ value have been extracted and gathered in a figure that shows the accelerations dependence on the examined variable.

### 4.5.2 Extent of the study

In this section an overview of all the performed calculations will be presented. The section is meant to show the extent of the parameter study together with the required performance to achieve this extent.

The examined parameters are the Young's modulus, E, the density, $\rho$, the total bridge length, $\mathrm{L}_{\text {tot }}$, the span relation, $\mu$, and the modal damping c . Calculations have been performed on both two- and three-span bridges. Table 4.8 shows the extent of the parameter study.

Table $4.8 \quad$ Extent of the parameter study

| $\begin{aligned} & \frac{0}{0} \\ & \frac{\pi}{5} \\ & \stackrel{y}{7} \end{aligned}$ | $\begin{aligned} & \stackrel{5}{50} \\ & \frac{0}{2} \\ & \frac{0}{2} \\ & \ddot{0} \end{aligned}$ | Two-span bridge |  |  | Three-span bridge |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | A1 | A4 | A8 | A1 | A4 | A8 |
| $E$ [GPa] | 5 | 15-60 | 15-60 | - | 15-60 | - | - |
| $\rho\left[\mathrm{kg} / \mathrm{m}^{3}\right]$ | 200 | 2800-4600 | - | - | 2200-4000 | - | - |
| $L_{\text {tot }}[\mathrm{m}]$ | 0.5 | 13-22.5 | 13-22.5 | 15-24.5 | - | - | - |
| $\mu[-], \eta[-]$ | 0.05 | 1-1.95 | 1-1.95 | 1-1.95 | 0.5-1.45 | 0.5-1.45 | 0.5-1.45 |
| $c$ [-] | 0.001 | 0.01-0-029 | - | - | - | - | - |

It is seen in Table 4.8 that the train loads HSLM-A1, -A4 and -A8 have been chosen for the parameter study. These loads have been chosen as they have different values of every parameter that defines the HSLM-A train loads.

The calculations behind the parameter study are very time consuming. As mentioned in Section 4.5.1 the response in time must be calculated for every train velocity of interest to receive the response from one single set of parameters. Table 4.8 show that this calculation procedure has been repeated 250 times using different parameter sets. This is despite limiting the parameter study to three HSLM-A train loads.

For the examination of every parameter the train velocity has been varied within the interval of $100-360 \mathrm{~km} / \mathrm{h}$. Both the interval of the train velocity and the size of the examined parameters shown by Table 4.8 have been chosen to resemble realistic values.

When examining one variable the remaining data set have been chosen to capture the resonance peaks of interest, namely at $\beta=0.5,0.33,0.25$ and 0.2 . This means that different data sets have been chosen for the examination of different variables. The data set used is shown in the caption of every figure presented from the parameter study later on.

### 4.5.3 Results

The effect on the acceleration amplitude in railway bridges from varying different material and geometric parameters have been examined. Each parameter has been examined individually and the results are presented in this section in a corresponding way. The cross-section defined in Section 4.1 have been used for all analysis presented in this section.

### 4.5.3.1 Variations of mass

The mass has been varied by changing the density. When the mass per meter of a bridge length is increased the eigenfrequencies, and consequently the resonance speeds, will decrease. As can be seen in equation (4.13) the critical speed depends on the inverse of the square root of the mass. Figure 4.28 shows an example of the maximum acceleration response in a bridge for different train velocities and it is seen how the resonance speeds decrease as the mass increase.


Figure 4.28
Acceleration response for a two-span bridge subjected to HSLM-A1 using $E=30 \mathrm{GPa}$, Ltot $=16 \mathrm{~m}, \mu=2$ and $\zeta=0.02$.

In Figure 4.28 it can be seen that the shape of the acceleration response curve is kept as the mass of the examined bridge varies. The accelerations at the resonance speeds have been collected and plotted against each respective density in Figure 4.29.


Figure 4.29 Acceleration response for a two-span bridge subjected to HSLM-A1 using $E=30 \mathrm{GPa}$, Ltot $=16 \mathrm{~m}, \mu=2$ and $\zeta=0.02$ at three resonance speeds.

Eurocode states that the amplitude of the acceleration response has an inverse proportionality against the mass. Figure 4.30 shows the normalized curves from Figure 4.29 , plotted against the inverse of each respective density. Here it can be seen that the statement in Eurocode is true. It is important to note that the acceleration amplitude has the same dependence for any speed, and hence only one line can be seen in Figure 4.30.


Figure $4.30 \quad$ Acceleration from Figure 4.29 plotted against the inverse of the density.

### 4.5.3.2 Variation of bending stiffness

From equation (4.13) it is known that the resonance speeds depend on the square root of the bending stiffness. An increase in the bending stiffness thus leads to the resonance peaks appearing at higher speeds. In this thesis the variation of bending stiffness has been examined by varying the module of elasticity. Figure 4.31 and Figure 4.32 shows examples of the acceleration response for the variation of train velocity for a two- and three-span bridge respectively.


Figure 4.31 Acceleration response for a two-span bridge subjected to HSLM-A1 using $=2800 \mathrm{~kg} / \mathrm{m}^{3}$, Ltot $=16 \mathrm{~m},=2$ and $=0.02$.


Figure 4.32 Acceleration response for a three-span bridge subjected to HSLM-A1 using $\rho=3000 \mathrm{~kg} / \mathrm{m}^{3}, \eta=1$, Ltot $=25 \mathrm{~m}, \kappa=1$ and $\zeta=0.02$.

It can be seen in Figure 4.31 and Figure 4.32 that both the acceleration amplitude and shape are completely independent of the choice of module of elasticity and hence also independent of the bending stiffness. For the purpose of giving more credibility to this shown phenomenon and give further insight in the response for the reader that is not too acquainted with the fields of dynamics additional calculations have been made in ADINA, see Figure 4.33 and Figure 4.34.


Figure 4.33 Acceleration response in time using HSLM-A1 and $E=30 \mathrm{GPa}, \rho=2400 \mathrm{~kg} / \mathrm{m}^{\wedge} 3$, Ltot $=15 \mathrm{~m}, \mu=2, \zeta=0.02$ and $v=399 \mathrm{~km} / \mathrm{h}$. The response is calculated in ADINA.


Figure 4.34 Acceleration response in time using HSLM-A1 and $E=50 \mathrm{GPa}, \rho=2400 \mathrm{~kg} / \mathrm{m}^{\wedge} 3$, Ltot $=15 \mathrm{~m}, \mu=2, \zeta=0.02$ and $v=516 \mathrm{~km} / \mathrm{h}$. The response is calculated in ADINA.

Figure 4.33 and Figure 4.34 shows the acceleration response in time using two different sets of module of elasticity and train velocity. The choice of parameter sets have been made to capture the response in time at $v_{c r} / 2$ or $\beta=0.5$. The figures show identical response curves but with different time axis implying that the maximum acceleration response are the same using the two different parameter sets. To show why the response in time becomes identical the analytical solution for a SDOF model has been examined, see Appendix E.

### 4.5.3.3 Variations of geometric parameters

The determination of dynamic acceleration response in railway bridges is a complex problem. The examinations of material parameters in Sections 4.5.3.1 and 4.5.3.2, though showed that the dynamic response has a simple relation to these parameters, which implies that the complexity solely lies in the variation of geometry.

How the acceleration response depends on geometric parameters cannot be expressed in separate expressions as for the material parameters. Instead they depend on each other. The geometric parameters for a bridge exist of the total bridge length and span relations. The reader should acknowledge however that the geometric parameters does not only include the total bridge length and span relation. There are also additional geometric parameters in the load in the form of bogie axle length and coach length. And it is a combination of these four geometric parameters together that affects the dynamic response.

With the goal of capturing the general variation trend of the acceleration response, examinations have still been performed on the geometric parameters. Variations in the load parameters have not been performed. Instead different HSLM-A loads have been used to examine the dynamic response during the variation of total bridge length and span relations independently.

### 4.5.3.4 Variation of total bridge length

The variation of total bridge length has been examined for two-span bridges only. The variation of total bridge length on two-span bridges will be shown to affect the dynamic response in such a complex way that no additional conclusions are presumed to be gained from analysis on three-span bridges. Figure 4.35 and Figure 4.36 shows examples of the acceleration response for a two-span bridge subjected to a HSLM-A load for different bridge lengths.


Figure 4.35 Acceleration response for a two-span bridge subjected to HSLM-A1 using $E=40 \mathrm{GPa}, \mu=2, \rho=3000 \mathrm{~kg} / \mathrm{m}^{3}$ and $\zeta=0.02$.


Figure 4.36 Acceleration response for a two-span bridge subjected to HSLM-A1 using $E=40 \mathrm{GPa}, \mu=2, \rho=3000 \mathrm{~kg} / \mathrm{m}^{3}$ and $\zeta=0.02$.

Just by examining Figure 4.35 and Figure 4.36 it can be seen that the acceleration response due to bridge length variation is not consistent. This means that the
acceleration has different dependencies on the bridge length for different values of $\beta$, i.e. the acceleration amplitude varies differently in different resonance peaks when the bridge length is changed, see Figure 4.35- Figure 4.38. More figures, similar to Figure 4.35- Figure 4.38 are presented in Appendix B.


Figure $4.37 \quad$ Acceleration response in resonance peak at $v_{c r} / 2$ for different train loads and a twospan bridge using $E=40 G P a, \mu=2, \rho=3000 \mathrm{~kg} / \mathrm{m}^{3}$ and $\zeta=0.02$.


Figure $4.38 \quad$ Acceleration response in resonance peak at $v_{c r} / 3$ for different train loads and a twospan bridge using $E=40 \mathrm{GPa}, \mu=2, \rho=3000 \mathrm{~kg} / \mathrm{m}^{3}$ and $\zeta=0.02$.

The following conclusions can be made concerning the acceleration amplitudes dependence of the total bridge length:

- The dependence differs for different $\beta$ values, meaning that the design velocity of the bridge considered will affect the dependence.
- The dependence will differ between eigenmodes for a constant train velocity since these correspond to different $\beta$ values.
- The dependence will differ between HSLM-A train loads for a constant train velocity since these correspond to different $\beta$ values.

These conclusions together make it difficult to derive an expression for how the acceleration amplitude will depend on the total bridge length, since the number of possible $\beta$ values and eigenmodes for a continuous structure are both infinitely large.

The general trend from varying the total bridge length that was found during the parameter study was that the acceleration amplitude decrease when the total bridge length is increased. However if this decrease depends on the combination of geometric parameters is unclear since increasing the total bridge length also increases the total bridge mass.

### 4.5.3.5 Variation of span relation for a two-span bridge

The eigenfrequencies for a two-span bridge was shown in Section 4.2 to be affected differently by the variation of span relation $\mu$ depending on the examined eigenmode. This knowledge alone makes it impossible to derive an expression for the acceleration amplitudes dependence on $\mu$. Each eigenfrequency will have a different dependence and the $\beta$ value will hence be affected differently for each eigenmode. Therefore the eigenmodes individual response is super-positioned differently for different values of $\mu$.

The variation of the span relation $\mu$ has still been examined with the purpose of capturing the general effect on the acceleration amplitude. Three different HSLM-A loads have been used in this examination, namely A1, A4 and A8. These three HSLM-A loads have been chosen as they differ in every load parameter that defines their appearance. Figure 4.39 and Figure 4.40 shows examples of the acceleration response for a two-span bridge for different values of the span relation using A1 and A8 respectively.


Figure 4.39 Acceleration response for a two-span bridge subjected to HSLM-A1 using $E=20 \mathrm{GPa}, L_{\text {tot }}=16 \mathrm{~m}, \rho=3000 \mathrm{~kg} / \mathrm{m}^{3}$ and $\zeta=0.02$.


Figure 4.40 Acceleration response for a two-span bridge subjected to HSLM-A8 using $E=20 \mathrm{GPa}, L_{\text {tot }}=16 \mathrm{~m}, \rho=6000 \mathrm{~kg} / \mathrm{m}^{3}$ and $\zeta=0.02$.

The acceleration amplitude has been gathered for constant $\beta$ values corresponding to resonance speeds, namely $0.5,0.33,0.25$ and 0.2 . Figure 4.41 and Figure 4.42 shows the variation of acceleration amplitude for $\beta$ equal to 0.5 and 0.33 . More detailed results are presented in Appendix B.


Figure $4.41 \quad$ Acceleration response in resonance peak at vcr/2 ( $\beta=0.5$ ) for different train loads using $E=20 \mathrm{GPa}, L_{\text {tot }}=16 \mathrm{~m}, \zeta=0.02$ and $\rho=3000,4000$ and $6000 \mathrm{~kg} / \mathrm{m}^{3}$ for A1, A4 and A8 respectively.


Figure $4.42 \quad$ Acceleration response in resonance peak at $v c r / 3(\beta=0.33)$ for different train loads using $E=20 \mathrm{GPa}, L_{\text {tot }}=16 \mathrm{~m}, \zeta=0.02$ and $\rho=3000,4000$ and $6000 \mathrm{~kg} / \mathrm{m}^{3}$ for Al , A4 and A8 respectively.

The acceleration amplitude seems to have a similar dependence on the span relation $\mu$ for all examined values of $\beta$ and all examined train loads. An interesting observation is that the acceleration amplitude increases when small the asymmetry is created in the bridge. For a two-span bridge it was shown in Section 4.2.1 that when asymmetry was increased the first eigenfrequency decreased while the second increased. The criterion set by the Swedish railway administration stating that only eigenfrequencies up to 30 Hz need to be considered then raises an interesting discussion. Because of this criterion there is often only need for considering the first, or first and second eigenmode. If creating small asymmetry always leads to increased accelerations this means that an increase of $\mu$ that puts the second eigenfrequency above 30 Hz may lead to lower design acceleration while the true acceleration increase. There might not be any conclusions to make of this observation, except that the 30 Hz criterion set by the Swedish railway administration might be considered questionable.

### 4.5.3.6 Variation of span relation for a three-span bridge

For the same reason as for a two-span bridge it is not possible to derive an expression for the acceleration amplitudes dependence on the span relation $\eta$ for three-span bridges. Examinations have still been performed with the purpose of capturing the general dependence. Resulting figures from this examination can be found in Appendix B. Figure 4.43 and Figure 4.44 show the variation of acceleration amplitude at $\beta=0.5$ and $\beta=0.33$ from the examination.


Figure $4.43 \quad$ Acceleration response in resonance peak at vcr/2 $(\beta=0.5)$ for different train loads using $\rho=2500 \mathrm{~kg} / \mathrm{m}^{3}, L_{\text {tot }}=20 \mathrm{~m}, \zeta=0.02, \kappa=1$ and $E=20,14$ and 10 GPa for A1, $A 4$ and $A 8$ respectively.


Figure 4.44
Acceleration response in resonance peak at vcr/3 ( $\beta=0.33$ ) for different train loads using $\rho=2500 \mathrm{~kg} / \mathrm{m}^{3}, L_{\text {tot }}=20 \mathrm{~m}, \zeta=0.02, \kappa=1$ and $E=20,14$ and 10 GPa for A1, A4 and A8 respectively.

The same conclusion as for a two-span bridge can be made for three-span bridges concerning the span relation $\eta$. The acceleration amplitude seems to increase when small asymmetry is created. However one deviating result was obtained in the analysis which increases the uncertainties in this conclusion, see Figure 4.43.

### 4.5.3.7 Variations of damping

The damping is the only examined variable that does not affect the eigenfrequencies and hence the resonance speeds. It does however affect the acceleration amplitude. The damping effect is not consistent but instead varies for different values of $\beta$. The largest effect occurs for $\beta$ values corresponding to resonance speeds while $\beta$ values in between these values can be unaffected, see Figure 4.45. The damping always lowers the acceleration response.


Figure 4.45 Acceleration response for a two-span bridge subjected to HSLM-A1 using $E=40 \mathrm{GPa}, \mu=2, \rho=4000 \mathrm{~kg} / \mathrm{m}^{3}$ and $L_{\text {tot }}=15 \mathrm{~m}$.

### 4.5.4 Conclusions

In Section 4.5.3 the results from the individual examinations of material and geometric parameters where presented. Several conclusions where made already in this section but will be gathered here for a more comprehensible presentation. Also some complementary observations will be presented concerning the general dynamic response behavior in railway bridges.

The following observations were made during the parameter study:

- The acceleration amplitude is independent of the module of elasticity.
- The acceleration amplitude is proportional to the inverse of mass.
- The total bridge length's effect on the acceleration amplitude depends on the value of $\beta$. This makes it hard to find an expression for this dependence since different train loads and eigenmodes will have a different effect for one constant train velocity.
- When comparing the acceleration amplitude in the resonance peaks for $v_{c r} / k$, with $k=2,3,4,5$, the highest amplitude is not always ordered from lowest to highest resonance peak, i.e. the following expression

$$
a\left(\frac{v_{c r}}{k_{i}}\right)>a\left(\frac{v_{c r}}{k_{i+1}}\right)
$$

does not hold for all combinations of geometric parameters. This depends partly on the additional observation of cancellation effects that can completely remove one of the resonance peaks for some sets of geometric parameters.

- The highest acceleration amplitude does not necessarily occur in a resonance peak if a maximum limit for the examined train speed is set (This mostly occurred when cancellation effects removed the resonance peak at $v_{c r} / 2$ ).
- The resonance effects does not always occur exactly at the critical speeds. This is especially true when cancellation effects occur but the phenomenon appears also for other speeds with small differences of around $5 \mathrm{~km} / \mathrm{h}$.


### 4.6 Case studies

The discussion presented in Section 4.5 .4 can be seen as the result from the examination of acceleration amplitudes. In this section three special case studies will be presented where all accumulated knowledge from the observations made will be applied to create guidelines and simplifications for three common railway bridge types. Two multi-span bridges have been examined: a two-span bridge with equal span lengths and a three-span bridge with equal span. Also a single-span bridge has been examined since it is a good application for the developed calculation tools in this section and serves as a basis for the 2D/3D comparison later in Chapter 5 .

The results presented in this section are in the form of diagrams, and referred to as design curves throughout this thesis. The name comes from the main concept of plotting the train velocities against the design acceleration that this velocity corresponds to instead of the actual acceleration. It is shown that the design curves can be used to decide which train loads that dominate the response for each specific bridge type. It is also possible to easily calculate the exact design acceleration for a certain set of material parameters. And as a third sector of application the curves can be used as guidance to which combination of bridge parameters that is preferable in a dynamic aspect.

### 4.6.1 Development of the design curves

In Section 4.5.3 and 4.5.4 it was shown that the geometric parameters of a railway bridge have a complicated effect on the acceleration response. Developing guidelines that treats an arbitrary combination of bridge length and span relations was shown to be a difficult task as the acceleration has a different dependence upon these variables for different eigenmodes and HSLM-A train loads. If a simplified model treating any combination of geometric parameters would be developed it would have to express different dependencies for every geometric parameter combination that is of interest.

In Section 4.5.4 some conclusions about the general response behavior of the accelerations were made concerning railway bridges. It was shown during the parameter study performed in Section 4.4 that it is not possible to predict which train velocity that will produce the design acceleration. Even if higher resonance peaks tend to give a higher acceleration response it is not consistently so. Also if a higher limit for the considered train velocities is set (which it always is in real life projects) it is
not necessarily the resonance peaks that contribute with the highest acceleration response.

However a discovery that contributes to increased possibility in finding guidelines for the dynamic response of railway bridges was found. The material parameters of a railway bridge have simple effect on both eigenfrequencies and the resulting acceleration amplitude. This conclusion forms the basis of the development of the design curves as the response behavior of the acceleration is kept under the variation of material parameters. Let us consider the difference in acceleration response from varying the module of elasticity (or bending stiffness), see Figure 4.31. An increase in the module of elasticity leads to increased eigenfrequencies and hence increased critical speed of the train loads, but the amplitude of the acceleration is unaffected. In Figure 4.31 the HSLM-A1 train load is used. If the different curves in Figure 4.31 are plotted against their corresponding $\beta$ value for the A1 load the curves will become identical.

The same approach can be used for the variation of density (or mass per meter of the bridge) with the difference that only the shape of the curves will be identical as the amplitude will vary. The difference in acceleration amplitude is however consistent for the variation of train velocity. For this reason the parameter $\tau$ is defined as

$$
\begin{equation*}
\tau=m \cdot a \tag{4.19}
\end{equation*}
$$

where

$$
\begin{array}{ll}
m= & \text { The mass per meter length of the bridge } \\
a= & \text { The acceleration response }
\end{array}
$$

Figure 4.46 is an illustration of the transformations from a $v-a$ plot to a $\beta-\tau$ plot.


Figure 4.46
Illustration of how different acceleration responses under the variation of material parameters can be expressed by one identical curve.

As mentioned above it is not possible to decide which train velocity (or $\beta$ value) will contribute to the highest acceleration response. Therefore it is here examined graphically. Since the material parameters' effect on the acceleration response, as shown in Figure 4.36, can be expressed by one single curve it is possible to develop one curve for each specific set of geometric parameters, from which it is possible to extract the exact acceleration response.

In the design of railway bridges the only acceleration response of interest is the highest one, the design acceleration. Therefore it is here suggested that a curve showing the corresponding design acceleration for different values of $\beta$ is developed.

This type of curve will be referred to as a design curve throughout this thesis. Figure 4.47 shows the continued transformation of the illustrative curve in Figure 4.46 into a design curve.


Figure $4.47 \quad$ Transformation of a regular response curve into a design curve.

In Section 4.5 .3 it was shown that the acceleration response has a complicated dependence on geometric parameters. The main conclusion stated was that the different dependencies for different geometric parameter sets have to be modeled separately. This is however not a possibility when the variation of several geometric parameters is of interest, since the number of different dependencies (or figures) to be developed becomes too extensive. If, on the other hand, special case studies are chosen it is possible to eliminate some of the geometric parameters. In this thesis three such case studies is made, which will be discussed later. What they all have in common is that fixed span relations have been chosen leaving only one geometric parameter for variation, namely the total bridge length. One design curve has to be developed for each length of interest for the different case studies. But if the design curves for one specific case study are gathered in one figure it would then serve as a useful comprehensive graphical calculation tool for the dynamic response of this bridge configuration.

There is however two additional complications whose solutions need to be clarified. First off the damping needs to be considered. It was shown in Section 4.5.3.7 that the damping's effect on the acceleration response varies depending on the value of $\beta$. The damping showed large change in effect at the resonance speeds but none in between. However since it was decided to create different design curves for each bridge length this is not a problem in practice. The code specifies the damping to be used for each bridge length, see Chapter 3. The design curve will hence include the effect of the considered bridge length and its corresponding damping.

Finally there are for each bridge length ten different HSLM-A train loads to account for. In real life design the only train load of interest is the one contributing to the highest acceleration response. However, it is hard to estimate which of the ten loads that will contribute to the highest acceleration, often resulting in a time consuming process where the response from all ten loads are examined. By plotting the design curves for each HSLM-A train load for a specific bridge length in one figure it is possible to graphically produce a design curve that includes the effect of all ten train
loads. One problem is that each train load corresponds to different $\beta-\tau$ plots since they all have different definitions of the corresponding $\beta$ value because of different coach lengths $D$. In this thesis this has been solved by simply plotting all the curves against the $\beta$ value of the HSLM-A1 train load. Figure 4.48 shows the final transformation of the illustrative design curve in Figure 4.47 into a design curve that includes the effect of all HSLM-A train loads.


Figure $4.48 \quad$ Illustration of how the response from different HSLM-A train loads can be combined in a single design curve

### 4.6.2 The design curves area of use

In real life design of railway bridges the dynamic aspect plays an important role in the design process. However, static analyses have been used in a larger extent in a historic perspective, since the requirement of dynamic analysis of railway bridges was introduced relatively recently in 2006. Static compared to dynamic analyses have in general simpler calculation procedures and engineers today possess larger experience with static aspects that govern railway bridge design. Therefore the early stage of the design process is today in a large extent still governed by calculations and guidelines based on experience from static analysis, even though the demands on dynamic criteria may prove to govern the final design. If a railway bridge is designed using guidelines based on static analysis and then shown to give unsatisfactory results in a dynamic aspect the design process becomes cumbersome and expensive.

Using the design curves it is possible to easily determine the design acceleration for a certain set of material parameters and total bridge length. It is important to note that the design acceleration determined from a design curve is not an approximation but in fact the exact design acceleration from a 2D analysis. So in the iterative process of statically designing the cross-section of a railway bridge the designer can easily check how the bridge performs in a dynamic aspect to avoid complications later in the project. The calculation is made in two steps. First the $\beta$ value corresponding to the set of bridge parameters and design train velocity is calculated as

$$
\begin{equation*}
\beta=\frac{v}{D} \cdot \frac{L^{2}}{g_{1}} \cdot \sqrt{\frac{m}{E I}} \tag{4.20}
\end{equation*}
$$

where
$E I=\quad$ Bending stiffness of the bridge
$m=\quad$ Mass per meter length of the bridge
$L=$ Total bridge length
$v=\quad$ Maximum allowed train velocity on the bridge
$D=\quad$ Coach length of the HSLM-A load chosen as reference for the $\beta$ value.
$g_{1}=\quad$ Value of the first eigenfrequency function for the given choice of span relations.

The calculated value of $\beta$ corresponds to a value of $\tau$ which is found by using the corresponding design curve. When the value of $\tau$ is known the design acceleration can be calculated using equation (4.19).

The design curves can also serve as guidance in the early design project life of a railway bridge since they basically give a complete overview of how the dynamic response change with modifications of the cross-section and changes in material parameters. Take the final design curve from the illustrative figures in the previous section as an example, see Figure 4.49.


Figure $4.49 \quad$ Illustration of the value $\beta^{*}$ which represents a value of $\beta$ which is preferable to fall below in a dynamic aspect if relevant values of $\beta$ are in the same range.

It will be shown in the upcoming presentation of real developed design curves that there, in the intervals of relevant $\beta$ values, often exist drastic changes in the design curves. Choosing the dimensions and materials of the cross-section with consideration to these changes greatly increases the possibility of satisfying the demands on maximum acceleration response. By combining equation (4.4), (4.8) and (4.10) an equation that gives guidance to the choice of cross-section and material parameters is obtained as

$$
\begin{equation*}
\frac{E I}{m}=\left(\frac{L^{2} \cdot v}{D \cdot \beta^{*} \cdot g_{1}}\right)^{2} \tag{4.21}
\end{equation*}
$$

where
$E I=$ Bending stiffness of the bridge
$m=\quad$ Mass per meter length of the bridge
$L=$ Total bridge length
$v=\quad$ Maximum allowed train velocity on the bridge
$D=\quad$ Coach length of the HSLM-A load chosen as reference for the $\beta$ value.
$g_{1}=\quad$ Value of the first eigenfrequency function for the given choice of span relations.
$\beta^{*}=$ Appropriate maximum value of $\beta$ according to Figure 4.49.

As a third sector of application the design curves can be used to verify more complex FE models created in commercial software. If a 3D model has been created it can be verified by using mode superposition with just the initial eigenmodes that does not cause rotation. This is further dealt with in Chapter 5.

### 4.6.3 Limitations of the design curves

The damping specified in Eurocode differs between reinforced concrete bridges and pre-stressed bridges. This means that different design curves need to be developed depending on the type of bridge that is designed.

Eurocode specifies that it is not allowed to interpolate between responses using different bridge lengths. If an interpolation is made between two design curves this result can hence only be treated as a rough estimation of the true response.

The number of eigenmodes used will affect the appearance of the design curves and must hence be chosen before design curves can be developed. The chosen number of eigenmodes for the special case studies in this thesis has been estimated based on the previous examination of eigenfrequencies for multi-span bridges. These choices will be presented in the upcoming section.

The design curves developed in this Chapter are all based on 2D analysis. If a 3D geometry is considered additional accelerations can be expected, since it basically mean that more eigenmodes have to be included. More about this limitation will be presented in Chapter 5 where dynamic analysis on railway bridges in 3D are treated.

### 4.6.4 Design curves for the three case studies

Three different bridge configurations have been chosen for the development of real design curves. These bridge configurations are:

- Single-span bridge
- Two-span bridge with equal span lengths
- Three-span bridge with equal span lengths

As mentioned in Section 4.6.2 one design curve for each bridge length of interest has to be developed. Also the number of eigenmodes for the bridge configuration needs to be chosen. Table 4.9 shows the chosen bridge lengths taken in consideration and the choice of eigenmodes used in the calculations.

Table $4.9 \quad$ Number of eigenmodes and considered bridge lengths for the four case studies

| Bridge configuration | Considered lengths | Number of eigenmodes used |
| :---: | :---: | :---: |
| Single-span bridge | $7-20 \mathrm{~m}$ | 2 |
| Two-span bridge | $16-30 \mathrm{~m}$ | 2 |
| Three-span bridge with <br> equal span length | $30-43.5 \mathrm{~m}$ | 5 |

The developed design curve are all for reinforced concrete. As mentioned pre-stress concrete corresponds to different damping according to BV Bro Banverket (2006) and their design curves must hence be developed separately.

The choices of lengths are meant to resemble common lengths for each respective bridge configuration. The creation of design curves has been limited to these length intervals since it is beyond the scope of this thesis to develop a complete calculation tool for the bridge configurations. The performed work has been restricted to the time limit of the thesis.

The chosen numbers of eigenmodes for each respective bridge configuration have been chosen based on the examination of eigenfrequencies in Section 4.2. For a single-span bridge it is known from experience at Reinertsen Sweden AB that a common value of the first eigenfrequency is above 8 Hz . In many design cases only the first eigenfrequency falls below 30 Hz since the second is four times larger, but it is not always so. For the purpose of getting design curves giving results on the safe side both the first and second eigenmode have been included in the development of design curves for single-span bridges. For the two- and three-span bridge all frequencies that fall below the 30 Hz criterion when the first frequency equals 8 Hz have been considered.

The calculation procedure presented in Section 4.5 .1 has been performed for the four case studies. A comparison of design curves from each respective HSLM-A train load for each considered bridge length has been made. These design curves have then been gathered in the final design curve for each bridge length and finally the design curves for every length of each bridge configuration have been gathered in one figure, see Figure 4.50 to Figure 4.52. For the purpose of giving a comprehensive presentation of the developed design curves all figures cannot be presented in the report. The interested reader is however referred to Appendix C where an extensive presentation of all figures from the development is made.


Figure 4.50
Design curves for a single-span bridge in the bridge length range of 7-20 $m$


Figure 4.51
Design curves for a two-span bridge with equal span lengths in the bridge length range of $16-30 \mathrm{~m}$


Figure $4.52 \quad$ Design curves for a three-span bridge with equal span lengths in the bridge length range of $30-43.5 \mathrm{~m}$

The developed design curves showed distinguished changes in the $\tau$ value for basically every bridge configuration and examined bridge length. This implies that there is great use in equation (4.21). This equation can be simplified for each examined bridge configuration for which the value of the eigenfrequency function is known. Also utilizing that the value of $D$ in equation (4.21) equals 18 m since the HSLM-A1 train load have been used as reference for the $\beta$ value in the design curves gives the following case specific expressions for each respective case study:

For a single-span bridge:

$$
\begin{equation*}
\frac{E I}{m}=\left(\frac{L^{2} \cdot v}{9 \pi \cdot \beta^{*}}\right)^{2} \tag{4.22}
\end{equation*}
$$

For a two-span bridge with equal span lengths:

$$
\begin{equation*}
\frac{E I}{m}=\left(\frac{L^{2} \cdot v}{36 \pi \cdot \beta^{*}}\right)^{2} \tag{4.23}
\end{equation*}
$$

For a three-span bridge with equal span lengths:

$$
\begin{equation*}
\frac{E I}{m}=\left(\frac{L^{2} \cdot v}{81 \pi \cdot \beta^{*}}\right)^{2} \tag{4.24}
\end{equation*}
$$

Every length specific design curve has been created as the maximum curve of all the corresponding load specific design curves. As mentioned above all the figures showing the load specific design curves cannot be included in the report, but as an example consider Figure 4.53.


Figure $4.53 \quad$ Design curves for each HSLM-A load for a three-span bridge with equal span lengths and a total length of 33 m

Similar curves as Figure 4.53 have been produced for every considered bridge length and bridge configuration. The curves show what HSLM-A load that governs the acceleration response. In Figure 4.53 it can be seen that A2, A7 and A9 governs the response for a three-span bridge with equal span lengths and a total length of 33 m .

It was found during the development of the design curves that there is a large variation in what HSLM-A load that governs the acceleration response. There was however one train load, namely HSLM-A8 that practically never governed the acceleration response for any length or bridge configuration. For the rare occasions when it did govern the response it was for high and short intervals of $\beta$. Another load that did not govern the response often was the A4, even if it did have influence for shorter lengths of the two-span bridge.

The HSLM-A loads that govern the acceleration response most frequently where A2 and A9. Both these loads have distinguishing load parameters compared to the other load combinations. A9 have together with A10 the largest value of the load $P$, namely 210 kN . A9 have however the smallest coach length of the two making resonance occur earlier. A2 has the largest bogie axle length at 3.5 m compared to the most common length of 2 m . It is interesting to note that only three HSLM-A loads have a bogie axle length larger than 2 m , and these are A2, A4 and A8. As mentioned A4 and A8 are the two HSLM-A loads that governed the acceleration response the least.

It was found that A2 governs the response for a lower interval of $\beta$ for every considered length and bridge configuration, with one single exception for a singlespan bridge of 17 m length.

### 4.6.5 Two calculation examples

In this section two examples of how to use the developed design curves is presented. The first example treats how the exact design acceleration can be calculated when all geometric and material parameters are known. The second example treats the use of a design curve as guideline in an early design stage when the cross-section is yet to be defined.

Consider first the calculation of design acceleration when all material and geometric parameters are known. Say we have a single-span bridge that is 10 m long with a rectangular cross-section that is 10 m wide and 0.6 m high. The material parameters are $E=40 \mathrm{GPa}$ and $\rho=2400 \mathrm{~kg} / \mathrm{m}^{3}$. The bridge should withstand train velocities of 250 $\mathrm{km} / \mathrm{h}$. According to BV Bro Banverket (2006) the bridge should then be designed for the train velocity $1.2 \cdot 250=300 \mathrm{~km} / \mathrm{h}$ or $83.33 \mathrm{~m} / \mathrm{s}$, see Chapter 3 . Using equation (4.20) we get that

$$
\beta=\frac{v}{D} \cdot \frac{L^{2}}{g_{1}(\mu)} \cdot \sqrt{\frac{M}{E I}}=\frac{83.33}{18} \cdot \frac{10^{2}}{\pi / 2} \cdot \sqrt{\frac{A \cdot 2400}{I \cdot 40 \cdot 10^{9}}} \approx 0.417
$$

From Figure 4.54 we get the corresponding design curve for the studied bridge configuration.


Figure $4.54 \quad$ Design curve for a single-span bridge with a total length of 10 m .
The corresponding $\tau$ value for our bridge configuration and $\beta=0.417$ is approximate $\tau=70000 \mathrm{~N} / \mathrm{m}$. Using equation (4.19) we can calculate the design acceleration to be

$$
a=\frac{\tau}{m}=\frac{70000}{A \cdot 2400} \approx 4.34 \mathrm{~m} / \mathrm{s}^{2}
$$

Consider now the following example. Say we are designing a two-span railway bridge with equal span-lengths which is supposed to be 20 m in total bridge length and have a width of 10 m . Again the bridge is designed for the train velocity $300 \mathrm{~km} / \mathrm{h}$. From Figure 4.51 we get the design curve corresponding to the bridge configuration, see Figure 4.55.


Figure 4.55
Design curve for a two-span bridge with equal span lengths and a total length of 20 m .

From Figure 4.55 it is seen that a drastic change in the acceleration occurs at $\beta$ values higher than 0.38 . Let us call this appropriate choice of $\beta$ for $\beta^{*}$. Falling below $\beta^{*}$ should be beneficial in a dynamic aspect with consideration to the demands on vertical acceleration. Using equation (4.23) we can calculate that

$$
\frac{E I}{m}=\left(\frac{L^{2} \cdot v}{36 \pi \cdot \beta^{*}}\right)^{2}=\left(\frac{20^{2} \cdot 83.33}{36 \pi \cdot 0.38}\right)^{2} \approx 6 \cdot 10^{5}
$$

correspond to an appropriate ratio between the bending stiffness and the mass per meter of the bridge to achieve this. This limit for the ratio between bending stiffness and mass per meter length is a lower limit. Increasing the ratio leads to an increased first eigenfrequency and a lower $\beta$ value. Say we are interested in a rectangular crosssection. The expression can then be simplified even further since

$$
\frac{E I}{m}=\frac{E \frac{b h^{3}}{12}}{b h \rho}=\frac{E}{\rho} \cdot \frac{h^{2}}{12}
$$

The expression becomes independent of the width and a lower limit for allowed height is received. If we use a density of $2400 \mathrm{~kg} / \mathrm{m}^{3}$ and module of elasticity of 40 GPa and assume that there is no ballast we get that

$$
h \approx \sqrt{6 \cdot 10^{5} \cdot 12 \frac{\rho}{E}}=0.657 \mathrm{~m}
$$

Basically we have now calculated that if we use a rectangular cross-section and a height above 0.66 m we will avoid the drastic peak in the acceleration design curve. To be sure this corresponds to the acceleration demand set by BV Bro Banverket (2006) we can calculate an upper limit for $\tau$ using equation (4.19). Say the limit for maximum allowed vertical acceleration is $5 \mathrm{~m} / \mathrm{s}^{2}$, which corresponds to a bridge without ballast. Then we get that

$$
\tau=b \cdot h \cdot \rho \cdot a_{d}=10 \cdot 0.657 \cdot 2400 \cdot 5=78840
$$

The calculated higher limit for $\tau$ is higher than the $\tau$ value corresponding to the chosen $\beta^{*}$ value of 0.38 . It can thus be concluded that a rectangular cross-section with a height larger than 0.66 m would be satisfactory in a dynamic aspect for the considered bridge.

## 5 Analysis in 3D

In Chapter 4 it was shown that the design acceleration for one specific set of geometric parameters can be presented by one single curve, independent of the choice of material parameters and designing train velocity. This curve is throughout this thesis called a design curve. The design curves were shown to give great support in the design process of railway bridges with regard to the demands on vertical acceleration. There were, however, some limitations for the use of the curves. This Chapter treats the main limitation, namely that the design curves were developed through 2D analysis.

All real life structures are three-dimensional. However the designs of many structures and details are made in 2D since these analyses in many cases can be made to resemble the real behavior. 3D analyses are significantly more complicated and time consuming than 2D analyses and hence it is preferable to use a 2 D analysis if possible.

This chapter treats the difference between the acceleration response in 2D and in 3D for railway bridges. The major difference comes from the fact that a 3D analysis model has eigenmodes that twist around the bridge length axis. This means that the response is be governed by eigenmodes that are affected differently by geometric changes and eccentricity of the load. The aim with the 2D/3D comparison is to get an understanding of complications that arise when three-dimensional geometry is considered. Also it is of interest to see how large difference in acceleration response that may exist between a 2D and a 3D model.

Section 5.1 presents the calculations that have been made in 3D and the corresponding assumptions for these calculations. Because of the 30 Hz criterion set by the Swedish railway administration all performed calculations and comparisons are based on FE analysis utilizing mode-superposition.

The studies of the acceleration response in 3D can be divided in three main areas. Eigenmodes and eigenfrequencies, variation in bridge width, and variation in load eccentricity. Each area is treated separately in Section 5.2 to 5.4. The difference in design acceleration is made in Section 5.5 by comparing the design curves created in 3D and 2D. Finally a discussion about the accumulated results is presented in Section 5.6.

### 5.1 Performed calculations and limitations

In this section all performed calculations on three-dimensional bridges is presented. Due to limit in time of this thesis several limitations have been made and these are also presented.

### 5.1.1 Geometry

Two different length configurations for a single-span bridge have been examined in the 3D analyses. Single-span bridges have been chosen since its geometry is simpler than multi-span bridges. The geometry is preferred since it simplifies the examination
of the added complexity that three-dimensional geometry contributes with. Two different lengths were chosen to increase the range of the examinations.

There are a number of different cross-sections used for railway bridges today. The design curves developed from 2D analysis in Chapter 4 are applicable for any form of cross-section geometry since the curves solely depend on area and moment of inertia of the cross-section, which independently affect the mass per meter length and bending stiffness respectively. However in 3D analysis when rotation around the longitudinal axis is allowed there is also an influence from bending stiffness perpendicular to the longitudinal axis and the rotational stiffness to consider. Different shapes of cross-section geometry will therefore have different effect on the acceleration response behavior and amplitude. The 3D calculations performed in this thesis have been limited to a rectangular cross-section. Cross-sections whose appearances require additional geometric parameters than cross-section height and width are presumed to add complexity in the calculations and have hence been avoided.

### 5.1.2 Method

Calculations have been limited to single-span bridges for variations in bridge width and load eccentricity as mentioned above. Table 5.1 shows the considered bridge lengths, widths and eccentricities in the performed calculations.

Table 5.1
Considered bridge configurations for the calculations in 3D

|  | Bridge length [m] | Bridge width | Load eccentricity for <br> width of 10 m |
| :---: | :---: | :---: | :---: |
| Single-span <br> bridge | 10 | $5 \mathrm{~m} / 1 \mathrm{~m} / 10 \mathrm{~m}$ | - |
| Single-span <br> bridge | 15 | $7.5 \mathrm{~m} / 1.5 \mathrm{~m} / 15 \mathrm{~m}$ | $0 \mathrm{~m} / 0.5 \mathrm{~m} / 4.5 \mathrm{~m}$ |

Table 5.1 shows that 12 different geometric configurations have been used for the variation of bridge width. For each configuration the total acceleration response and each eigenmodes separate response has been calculated in the whole bridge and in two specific nodes as shown in Figure 5.1. The calculated acceleration responses include the total and individual effect of all HSLM-A train loads.


Figure 5.1 Definition of bridge width, bridge length and load eccentricity for the calculations in 3D. Also the two nodes from where separate data has been extracted are shown.

The effect on the acceleration response by varying the load eccentricity has been made on a single-span bridge with a length of 15 m and width of 10 m . The train load HSLM-A1 was used for the examinations on eccentricity variation.

All the 3D calculations have been made using Matlab programs created by the authors. The same main program as for the 2D analysis has been used with the difference that new function files treating the mesh and load creation have been used. All the programs used for the analysis can be found in Appendix D together with explanations of the program structures and verifications against ADINA. The reasons for using Matlab are the same as for the calculations in 2D, i.e. that the variation of geometric parameters and train velocity and the application of load are substantially easier to handle in Matlab compared to commercial FE software.

### 5.2 Eigenmodes and eigenfrequencies

The differences in eigenfrequencies and eigenmode shapes between 2D and 3D analyses are what make the dynamic responses differ. The change in frequencies will cause resonance to occur for different train speeds and rotation around the longitudinal axis will increase the influence of bridge width and load eccentricity.

Two different forms of eigenmodes can be distinguished in a 3D analysis, in this thesis they are referred to as bending modes and torsion modes. Bending modes are basically the eigenmodes from 2D analyses where bending in the longitudinal direction governs the eigenmode shape. Torsion modes are governed by torsion around the length coordination of the bridge and bending in the transverse direction. These eigenmodes cannot be described by 2D analyses. So basically bending and torsion modes represent eigenmodes from 2D and the additional modes in 3D respectively. It should be noted that the definitions of bending and torsion modes have been created by the authors, and are not generally accepted. The reason why they are used here is to easier explain the performed study, which is focused on differences between 2D and 3D analysis. Figure 5.2 shows the two bending and torsion modes with lowest eigenfrequency for a single-span bridge with a length of 15 m and a width of 10 m .


Figure 5.2 The two bending modes $(a$ and $b)$ and torsion modes ( $c$ and $d$ ) with lowest corresponding eigenfrequency for a single-span bridge with a length of 15 m , width of 10 m and height of 0.6 m .

As shown in Figure 5.2 there is small change in appearance in the bending modes compared to a 2D analysis where the section perpendicular to the length coordination remains straight. It will be shown that this bending in the cross-section plane will make the bending modes contribute to changes in acceleration with variations in bridge width and load eccentricity. The eigenfrequency is also affected but remain very close to that of a 2 D analysis.

The torsion modes can have rather low eigenfrequencies and hence easily fall below the 30 Hz limit in BV Bro Banverket (2006). In most design cases the two eigenmodes with lowest corresponding eigenfrequencies are the first bending and torsion mode, respectively. If the 30 Hz criterion is considered this basically means that a larger number of eigenmodes need to be considered in a 3D analysis since the bending modes from 2D analysis in principal are kept while an addition of torsion modes is made. Consequently an increased acceleration response can be expected.

In 2D analyses an important feature of the eigenfrequencies was that the relation between eigenfrequencies of different degree was kept constant for any change in material parameters, section geometry or bridge length. A single-span bridge as an example has in a 2D analysis a second and third eigenfrequency that correspond to four and nine times the first eigenfrequency, respectively. This made the creation of design curves possible since the eigenmodes corresponding $\beta$ values also had a constant relation to each other. Take Figure 4.33 and Figure 4.34 as an example. These figures show the response in time for the first eigenmode using different modules of elasticity with train velocities chosen to get $\beta=0.5$ in both cases. If similar curves for the second eigenmode where created they would both correspond to $\beta=0.125$ for the same bridge parameters, since the used eigenfrequency in the calculation of $\beta$ would be four times larger in both cases. This means that the shape of the time response in the two figures would be identical also for the second eigenmode, and hence also for the total response considering both eigenmodes. If a parameter that
changes the relation between first and second eigenfrequency (the span relation $\mu$ in 2D analysis as an example) is altered instead, the nice feature of unaffected shape of the total response, considering both eigenmodes, would be lost. Choosing two train velocities that correspond to the same $\beta$ value for the first eigenmode would then correspond to two different $\beta$ values for the second eigenmode. Consequently it can be stated that a parameter that changes the relation between eigenfrequencies when altered affect the acceleration response in a complex way, and hence, its variation cannot be described by a design curve.

In a 2 D analysis the cross-section geometry has no effect on the relation between eigenfrequencies since it only affects the moment of inertia and area of the bridge which has the same effect on all eigenfrequencies. In 3D analyses though the shape and size of the cross-section will affect the relation between eigenfrequencies. Figure 5.3 to Figure 5.5 shows the eigenfrequencies and the ratio between the eigenfrequencies for the eigenmodes shown in Figure 5.2 for three different lengthwidth configurations and the variation of cross-section height. The frequencies have been calculated in ADINA.


Figure 5.3 Left: Eigenfrequencies for the eigenmodes in Figure 5.2 with $L=5 \mathrm{~m}, b=5 \mathrm{~m}$, $E=30 G P a, v=0.02, \rho=2400 \mathrm{~kg} / \mathrm{m}^{3}$. Straight lines correspond to the two lowest analytical calculated frequencies in 2D. Right: Corresponding rations of the frequencies to that of the first bending mode.


Figure $5.4 \quad$ Left: Eigenfrequencies for the eigenmodes in Figure 5.2 with $L=10 \mathrm{~m}, b=10 \mathrm{~m}$, $E=30 \mathrm{GPa}, v=0.02, \rho=2400 \mathrm{~kg} / \mathrm{m}^{3}$. Straight lines correspond to the two lowest analytical calculated frequencies in 2D. Right: Corresponding rations of the frequencies to that of the first bending mode.


Figure $5.5 \quad$ Left: Eigenfrequencies for the eigenmodes in Figure 5.2 with $L=10 \mathrm{~m}, \mathrm{~b}=5 \mathrm{~m}$, $E=30 \mathrm{GPa}, v=0.02, \rho=2400 \mathrm{~kg} / \mathrm{m}^{3}$. Straight lines correspond to the two lowest analytical calculated frequencies in 2D. Right: Corresponding rations of the frequencies to that of the first bending mode.

Several observations can be made from Figure 5.3 to Figure 5.5.

- Variations in height affect the relation between eigenfrequencies and the dependence changes for different length-width configurations.
- Frequencies in 3D for bending modes deviate from the analytically calculated frequencies in 2D for larger heights.
- Increasing the width of the bridge has basically no effect on the bending modes.
- Increasing the width of the bridge decreases the frequency of the torsion modes.

For real life designs the height of rectangular cross-sections does not vary extensively but commonly lies in the range $0.5-0.8 \mathrm{~m}$. Even though it is shown that the relations between eigenfrequencies are affected by changes in cross-section height they are small within the commonly used range. Therefore no further investigations on the variation of height have been made in this thesis. The height 0.6 m has been used for all remaining calculations concerning variation of bridge width and load eccentricity.

In 2D analysis the fact that all eigenmodes have the same relation to each other means that the different eigenmode shapes will be ordered in the same way for any choice of parameter set. For a single-span bridge the first eigenmode will approximately have the shape of a half sinus curve, the second a full sinus curve and so on. In 3D analysis the shape of the bending modes will come in the same order as in 2D analysis as they are very similar to these eigenmodes. The shape of the additional torsion modes may however change order. Consider Figure 5.2 where the two torsion modes with lowest corresponding eigenfrequency for a single-span bridge of 15 m length and 10 m width are shown. The first torsion mode will always be the one shown in Figure 5.2c for any choice of length-width ratio. The shape of the second eigenmode however depends on the length-width ratio of the bridge, where a smaller length-width ratio causes the appearance of the eigenmode to change. Figure 5.6 shows the second and third torsion mode of a single-span bridge with lower length-width ratio.


The second and third torsion mode for a single-span bridge with a length and width of 10 m and height 0.6 m .

As shown by Figure 5.2 and Figure 5.6 the second torsion mode becomes the third when the bridge length-width ratio is decreased. For the following examinations on bridge width and load eccentricity a choice of included eigenmodes must be made. Since the relation between torsion modes do not remain for the variation of crosssection geometry examination on width variation becomes more difficult when more torsion modes are included. Therefore a choice has been made to only include the first torsion mode in the upcoming calculations on bridge width and load eccentricity variations. However for the purpose of comparing the change in design curves in 2D and 3D for the two examined single-span bridges it is a requirement that two bending modes are used in order to be consequent with the design curves created in Chapter 4 for single-span bridges. This adds complications in the calculations since the second bending mode will change to higher eigenmode degree when width is increased because of lowered eigenfrequency of torsion modes. As a solution to this problem a Matlab program that can find the three specific eigenmodes of interest utilizing inverse iteration has been made by the authors. The Matlab program can be found in Appendix D.

### 5.3 Variation of width

It was mentioned in Section 5.2 that variation of cross-section geometry does not affect the relation between eigenfrequencies in 2D analysis and therefore it is possible to develop one single design curve for each length that corresponds to an arbitrary choice of cross-section. It was also shown that this is not possible for the variation of bridge width in a 3D analysis as the bridge width affect the bending and torsion modes in a more complex manner. This will be explained in more detail in this section and a graphical presentation of how the response differs for the variation of bridge width will be made.

In Section 5.1 it was shown that the calculations on bridge width variation in 3D have been made on two single-span bridges with the length of 10 and 15 m . Every figure presented in this section is a result from these calculations. As mentioned the calculations have included all the HSLM-A train loads and gathered the total acceleration response in different nodes and also the individual response of different
eigenmodes in these nodes. The results presented in this section are all for the HSLMA1 train load and the node at the edge, in the middle of the bridge, see Figure 5.1.

### 5.3.1 Influence on bending modes

For a rectangular cross-section the increase of width does not affect the design curve from a 2D analysis. This is because the shape of the $v-a$ plot is unaffected. Figure 5.7 shows a comparison of three $v-a$ plots from 3D analysis where the only difference is the width of the cross-section.


Figure 5.7 Maximum acceleration in time for variation of train velocity for a single-span bridge with length of 10 m , height of 0.6 m and $E=30 \mathrm{GPa}, \rho=4000 \mathrm{~kg} / \mathrm{m}^{3}, v=0.02$ and $\zeta=0.022$

As can be seen in Figure 5.7 the shape of the $v-a$ plot changes with changed width, and hence also the corresponding design curve. The individual responses from the bending modes and the torsion mode have been studied for the purpose of explaining why and how the $v-a$ curves change. Consider first the response from the bending modes. Figure 5.8 shows different $v-a$ plots for the single-span bridge of 10 m length considering only the first bending mode.


Figure $5.8 \quad$ Maximum acceleration in time for variation of train velocity considering the individual response of the first bending mode for a single-span bridge with length of 10 m , height of 0.6 m and $E=30 \mathrm{GPa}, \rho=4000 \mathrm{~kg} / \mathrm{m}^{3}, v=0.02$ and $\zeta=0.022$.

As Figure 5.8 shows the shape of the $v-a$ curve from the first bending mode remains as the width is varied. The change in amplitude can be explained by the change of bridge mass, where the amplitude was shown in Chapter 4 to have an inverse proportionality against the mass. Figure 5.9 and Figure 5.10 shows the $v-a$ plots in Figure 5.8 transformed to $\beta$ - $\tau$ plots.


Figure 5.9 Corresponding $\beta$ - $\tau$ curves for the $v$-a curves in Figure 5.8.


Figure 5.10
Magnification of Figure 5.9 for the interval $0.48<\beta<0.52$.
It can be seen in Figure 5.10 that the $\beta-\tau$ curves for the bending mode do not have identical shapes for the variation of width as they would in a 2D analysis. The curves have small differences in position and amplitude. Both these observations are believed to be explained by small changes in eigenmode shape. As shown in Figure 5.2 the bending modes do include rotation around the bridge length axis, and this rotation causes the width to influence both eigenfrequency and acceleration amplitude. The widths effect on the acceleration behavior and amplitude are however small and it can be concluded that the major changes in acceleration response that were seen in Figure 5.7 are caused by the addition of a torsion mode.

### 5.3.2 Influence on torsion modes

Consider now the individual response of the first torsion mode. The maximum acceleration responses for different widths under the variation of train velocity are shown in Figure 5.11. A constant eccentricity of 2 m has been used in the calculations.


Figure 5.11 Maximum acceleration in time for variation of train velocity considering the individual response of the first torsion mode for a single-span bridge with length of 10 m , height of 0.6 m and $E=30 \mathrm{GPa}, \rho=4000 \mathrm{~kg} / \mathrm{m}^{3}, v=0.02, \zeta=0.022$, $e=2 \mathrm{~m}$.

As mentioned in Section 5.2 the eigenfrequencies of the torsion modes are affected by the change in width of the cross-section. Figure 5.11 clearly shows how the placement of the response curve moves to the left because of the lowered eigenfrequency with increased width. The $v-a$ curves in Figure 5.11 have been transformed to $\beta$ - $\tau$ curves with the purpose of examining the acceleration response behavior for the torsion mode, see Figure 5.12.


Figure 5.12 Corresponding $\beta$ - $\tau$ curves for the $v$-a curves in Figure 5.11.

Figure 5.12 shows that the shape of the $\beta$ - $\tau$ curve for the first torsion mode is kept as the width is increased. Also the amplitude of the acceleration is decreased with an increase in width. The change in amplitude of the accelerations seem however to have a different dependence on the increase of bridge width. The parameter $\tau$ can therefore not be used to create a curve that is independent of the bridge width as it can for bending modes. Figure 5.13 shows the acceleration for a fixed $\beta$ value in Figure 5.12 plotted against the inverse of corresponding bridge width.


Figure 5.13 The acceleration amplitudes relation to the inverse of the bridge width in Figure 5.12 for $\beta=0.19$.

The curve in Figure 5.12 would have had a constant derivative of one if the torsion mode had the same dependence on the increase of bridge width as bending modes.

Basically Figure 5.12 shows that the response from the torsion mode has the same type of behavior as a bending eigenmode when plotted against its own corresponding $\beta$ and an observation like resonance peaks at $\beta \approx 0.33,0.25$ and 0.2 can be made. It would be possible to create a $\beta$ curve with constant amplitude if a parameter (like $\tau$ for bending modes) were established that described the material and cross-section parameters effect on the acceleration amplitude. Additional calculations are however required to achieve this since the performed width variation affects bending stiffness, rotational stiffness and mass and hence does not give sufficient understanding of the problem.

### 5.3.3 Combined response

If the response from the bending modes and the torsion mode were considered separately two different design curves could be developed that gave a simple comprehensive view of the dynamic response, corresponding to an arbitrary choice of bridge width. Unfortunately the direct summation from the two curves would not correspond to the true response, considering both eigenmodes. The reason is that the two design curves would correspond to the maximum acceleration in time considering the individual response from the bending and torsion modes. Hence the $v-a$ plots for the different eigenmodes cannot be superpositioned. In mode-superposition it is the
response in time from different eigenmodes that are superpositioned to get the total response in time. The maximum accelerations in time for different eigenmodes are likely to occur at different times. Figure 5.14 shows the superposition of the individual $v-a$ plots for the different eigenmodes and the $v$ - $a$ plot considering both eigenmodes.


Figure 5.14 Comparison of the total response considering both the first bending and torsion mode with the superposition of the individual response of the eigenmodes.

The true contribution from the torsion mode is at all times equal to or smaller than the response received from superpositioning the $v-a$ curves of the eigenmodes individual response. This contribution from the torsion mode is considerably more difficult to map compared to the maximum acceleration in time. The true contribution for a certain set of geometric, material and train parameters will be the contribution at the time that corresponds to the maximum acceleration for the combined response of eigenmodes. This time may occur when the torsion mode has its individual maximum but also for a time when the contribution is zero or even negative, causing a decrease of the maximum acceleration in time. Figure 5.15 shows an example of the true contribution from the torsion mode compared to the contribution if the eigenmodes individual $v$ - $a$ curve is added.


Figure 5.15
The true contribution from the torsion mode to the total response compared to the contribution if the eigenmodes individual $v$-a curve is added for a single-span bridge with length 10 m , width 6 m , height $0.6 \mathrm{~m}, E=30 \mathrm{GPa}, \rho=4000 \mathrm{~kg} / \mathrm{m}^{3}, v=0.02$, $\zeta=0.0185, e=2 \mathrm{~m}$.

Finally Figure 5.16 shows the true contribution from the torsion mode for the variation in bridge width.


Figure 5.16
The true contribution from the torsion mode to the response considering only the first bending mode for a single-span bridge with length 10 m , width 6 m , height 0.6 m , $E=30 \mathrm{GPa}, \rho=4000 \mathrm{~kg} / \mathrm{m}^{3}, v=0.02, \zeta=0.0185, e=2 \mathrm{~m}$.

Figure 5.16 show that there are no general similarities between the true contributions from the torsion mode for different widths. Interesting to note is that the torsion mode in some cases does decrease the maximum acceleration in time. It should also be noted that the contributions shown by Figure 5.16 are for the HSLM-A1 train load in
the edge node shown by Figure 5.1. Only added complexity is presumed to be added when all HSLM-A trainloads and the maximum response for arbitrary node is considered.

### 5.4 Variation of load eccentricity

In a 2 D analysis no consideration is given to the eccentricity of the train load. Eigenmodes that include large deformation parallel the length coordination are highly influenced by the eccentricity as it by definition decides the coordinate of load application parallel the length axis. As was shown in Section 5.2 it is mainly the torsion modes that hence is affected by eccentricity but also the bending modes to some extent.

The eccentricities' effect on each individual eigenmode will be presented initially in this section and then followed by the effect on the total response. The examinations on eccentricity variation have been in the middle line of the bridge in the two nodes shown in Figure 5.1. In these nodes there is no contribution from the second bending eigenmode and hence only responses from the first bending and torsion mode will be presented. All results presented in this section are the individual response of the HSLM-A1 train load.

### 5.4.1 Influence on bending modes

The bending eigenmode will be affected by the eccentricity as it includes rotation perpendicular to the length coordination. The effect is however small, see Figure 5.17.


Figure 5.17 Maximum acceleration in time for the first bending mode in the middle of the bridge for a single-span bridge with length 15 m , width 10 m , height $0.6 \mathrm{~m}, E=100 \mathrm{GPa}$, $\rho=4000 \mathrm{~kg} / \mathrm{m}^{3}, v=0.2$. The eccentricity varies from 0-4.5 m.

The accelerations at $v=196.25 \mathrm{~m} / \mathrm{s}(\beta=0.5)$ have been gathered in Figure 5.18 with the purpose of examining the eccentricities' effect on the acceleration response through the first bending mode.


Figure $5.18 \quad$ Acceleration in the middle and edge node in Figure 5.1 from the first bending mode for a single-span bridge with length 15 m , width 10 m , height $0.6 \mathrm{~m}, E=100 \mathrm{GPa}$, $\rho=4000 \mathrm{~kg} / \mathrm{m}^{3}, v=0.2$ and HSLM-A1 with velocity of $196.25 \mathrm{~km} / \mathrm{h}$.

An increase in acceleration with increased eccentricity is expected. The bending eigenmode can be said to be a combination of two different motions. One corresponds to the motion in 2D and is unaffected by the eccentricity and the other consist of the rotation perpendicular to the length coordination and is more influenced as the eccentricity increase. In general it can be said that the influence a load has on an eigenmode is largest when the load is placed where the eigenmode deforms the most. Larger influence on the eigenmode creates larger acceleration through it. It can be seen that the accelerations seem to increase with the same shape as the cross-section deforms for the considered bending mode, see Figure 5.2a.

### 5.4.2 Influence on torsion modes

The largest impact from a change in eccentricity will be through torsion modes which are governed by rotation perpendicular the length coordination. Figure 5.19 shows the individual response from the first torsion mode.


Figure 5.19 Maximum acceleration in time for the variation of train velocity and load eccentricity for a single-span bridge with length 15 m , width 10 m , height $0.6 \mathrm{~m}, \mathrm{E}=100 \mathrm{GPa}$, $\rho=4000 \mathrm{~kg} / \mathrm{m}^{3}, v=0.2, \zeta=0.0185$

The accelerations at $v=152.5 \mathrm{~m} / \mathrm{s}$ has been gathered with the purpose of examining the eccentricities' effect on the acceleration response through the bending eigenmode, see Figure 5.20.


Figure 5.20 Acceleration in the edge node in Figure 5.1 from the first torsion mode for a singlespan bridge with length 15 m , width 10 m , height $0.6 \mathrm{~m}, E=100 \mathrm{GPa}$, $\rho=4000 \mathrm{~kg} / \mathrm{m}^{3}, v=0.2$ and HSLM-A1 with velocity of $196.25 \mathrm{~km} / \mathrm{h}$.

The same arguments for why the accelerations increase with the eccentricity for the bending mode can be used for the torsion mode. It can be seen in Figure 5.20 that the accelerations seem to increase with the same shape as the cross-section deforms also for the torsion mode.

### 5.4.3 Combined response

Finally we consider the total response of the single-span bridge for which the individual response of the first bending and torsion mode have been presented. For the edge node, the $v-a$ plots for the single-span bridge with variation in load eccentricity are shown in Figure 5.21.


Figure 5.21 Maximum acceleration in time for different train velocities considering the first bending and torsion mode for a single-span bridge with length 15 m , width 10 m , height $0.6 \mathrm{~m}, E=100 \mathrm{GPa}, \rho=4000 \mathrm{~kg} / \mathrm{m}^{3}, v=0.2, \zeta=0.185$ and HSLM-Al with velocity of $196.25 \mathrm{~km} / \mathrm{h}$.

The change in acceleration due to load eccentricity from Figure 5.21 for constant train velocities are shown by Figure 5.22.


Figure 5.22
Accelerations from Figure 5.21 for constant train velocities with increased load eccentricity.

As seen in Figure 5.22 there is no exact influence that can be derived with regards to the eccentricity. In general the eccentricity will increase the total acceleration response in a linear manor, the same way as it affects the torsion mode. It can be seen that the increase in acceleration is linear for higher values of eccentricity for all three considered values of train velocity. But for lower values of eccentricity an increase in acceleration is not guaranteed. The reason for this is that the time for which the maximum acceleration occur shifts for lower eccentricities. If the torsion mode has an acceleration response with opposite sign at the bending modes maximum there might even be a decrease in maximum acceleration, as is the case for the train velocity 152.5 $\mathrm{km} / \mathrm{h}$ in Figure 5.22.

### 5.5 Design curves

Last in the investigation of variation of bridge width and load eccentricity for a singlespan bridge is to show the difference between design curves in 2D and 3D analysis. The design curves in 3D for each width have been compared with the 2D design curve obtained in Chapter 4, see Figure 5.23 and Figure 5.24. In both figures the 2D design curve is for all velocities the lowest design curve. It is the first torsion mode that causes different total response for the variation of bridge width, as mentioned in Section 5.3. The complexity in the response due to varying width comes from superposition of the bending and torsion modes, whose frequencies are affected differently by the variation. Figure 5.23 and Figure 5.24 show that it is not possible to predict which width that will have the largest acceleration.


Figure 5.23
Comparison of design curves for the variation of bridge width in 3D and the $2 D$ design curve, with $L=10 \mathrm{~m}$.


Figure 5.24
Comparison of design curves for the variation of bridge width in $3 D$ and the $2 D$ design curve, with $L=15 \mathrm{~m}$.

The percental contribution from the torsion mode is shown in Figure 5.25 and Figure 5.26, for a bridge length of 10 and 15 m respectively. It can be concluded that the percental contribution varies for different combination of width, length and velocity. In the case of the 10 meter bridge the contribution is as high as $50 \%$ for $\beta=0.4$, which is a $\beta$ value likely to occur in a real life project. Hence the use of the design curve from 2D without consideration to 3D effects is for $\beta=0.4$ a large underestimation. In Figure 5.25 and Figure 5.26 it can be seen that the difference between a 2D and 3D analysis generally is smaller in percent for the longer bridge, even though the relation between length and width are the same. It is however not possible to draw any general conclusions if a longer bridge with the same length and width relation implies lower contributions from the torsion mode. For this a larger range of calculations are required.


Figure 5.25 Differences in design acceleration between $2 D$ and $3 D$ analysis for $L=10 \mathrm{~m}$


Figure 5.26 Differences in design acceleration between 2D and $3 D$ analysis for $L=15 \mathrm{~m}$.
The design curves for different eccentricities in 3D and the 2D design curve are compared in Figure 5.27, for a bridge with length 15 meter. It can be seen that the 2D design curve form the lower limit of the design curves and that the resulting acceleration response will increase with increasing eccentricity. This is an expected behaviour because of the increase in acceleration response from the torsion mode due to increased eccentricity that could be observed in Section 5.4. The percental increase in size of the design curves in 3D for variation of load eccentricity compared to the 2D design curve is shown in Figure 5.28.


Figure 5.27
Comparison of $3 D$ design curves for different eccentricities in the edge node and the design curve from the $2 D$ analysis, with $L=15 \mathrm{~m}$ and $b=10 \mathrm{~m}$.


Figure 5.28
Differences in design acceleration between 2D and 3D analysis in the edge node for $L=15 \mathrm{~m}$ and $b=10 \mathrm{~m}$.

### 5.6 Conclusions from the analysis in 3D

The main difference between 2D and 3D analysis is that torsion modes are included. In the previous sections of this Chapter it has been shown that a 3D analysis contributes with significantly more complexity compared to a 2D analysis. This complexity is mainly because of the additional torsion modes but also changes in the shape of the bending modes. Below follows the conclusions made from the 3D analysis in this thesis, with regard to single-span bridges and rectangular crosssection.

- Bending modes change in shape when going from 2D to 3D. In 3D the bending modes include torsion around the lengths coordination and the amount included depends on geometric parameters like width, height and length of the bridge.
- Torsion modes are included in 3D analysis. The eigenfrequencies of the lowest torsion modes are in range of the lowest bending modes eigenfrequencies and the two lowest eigenfrequencies seem to always be that of the first bending and torsion mode for the studied configurations.
- Cross-section height affects the relation between all eigenmodes, including bending modes. There may be a distinct deviation of the second bending modes eigenfrequency in 3D compared to the analytical from 2D for larger cross-section heights. However the height has a small influence on the ratio between the first bending and torsion mode.
- Eigenfrequencies of bending modes are in principal unaffected by variation in cross-section width but torsion modes are highly affected.
- The variation of acceleration response from the first bending mode due to changed width can be described by a design curve. The accelerations change in amplitude is in principal proportional to the cross-section width and hence also the bridge mass.
- The variation of acceleration response from the first torsion mode from changing the width cannot be described by a design curve using the parameter $\tau$. The change in acceleration amplitude cannot be explained solely by the change in bridge mass.
- The true contribution when considering the first bending and torsion mode compared to only the first bending mode for variation in bridge width is very complex. The complexity depends mainly on the change in eigenfrequency of the torsion mode, which changes the total acceleration response in time for a fixed train velocity (or value of $\beta$ for the bending mode).
- The load eccentricity's effect on bending modes is small but still present. The load's influence on the bending modes increase with increased eccentricity.
- The eccentricity has a large influence on the effect from the torsion modes. For the first torsion mode the increased acceleration is almost proportional to the eccentricity.
- The true change in design acceleration, for variation in eccentricity, is unpredictable when considering the first bending and torsion mode. A small added eccentricity may change the time for which the maximum acceleration occurs; it can even result in a decrease of the maximum acceleration.
- The increase in design acceleration in a 3D analysis compared to a 2D analysis can be as high as $100 \%$ for large load eccentricities.


## 6 Discussion

As mentioned in the background to this thesis there is a need today to find guidelines and simplifications for the dynamic design of railway bridges subjected to high speed trains. Basically guidelines and simplifications cover two different areas of use, where guidelines are meant to be used parallel static analysis in an early design process and simplifications are meant to reduce calculation time for the advanced dynamic analysis.

In this thesis the main focus has been put on developing guidelines for the early design process and the main result has been the idea behind the design curves. The design curves developed in 2D do however not give an estimation of accelerations but the exact design accelerations and are hence also a first step for simplifying the time consuming advanced dynamic analysis.

The discussion presented in this chapter will attend the advances made in this thesis concerning both guidelines and simplifications in dynamic design of railway bridges. Is the design curves developed in 2D applicable in real life projects? Is it possible to develop guidelines considering 3D geometry? Is it possible to develop guidelines that completely exclude the advanced and time consuming dynamic analysis?

Discussion will also be presented on how advanced dynamic analysis could be used for a more time efficient design process.

### 6.1 Analysis in 2D

As mentioned in the introduction of this thesis several earlier theses have been performed regarding dynamic analysis of railway bridges within the last years. Three master theses have been carried out in cooperation with Reinertsen Sweden AB, namely Ekström and Kieri (2007), De Leon and Lasn (2008) and Gustavsson (2008). These theses are all based on 2D analysis.

One of the advances made in this thesis concerning 2D analysis is the application of the SDOF model on multi-span bridges. It was shown that the force scaling approach gives a good comparison of results to a FE analysis considering the first eigenmode. There is however some lack in accuracy which varies for different HSLM-A loads and train velocities. In De Leon and Lasn (2008) it was concluded that the SDOF model can be used for approximating the dynamic response. That thesis were however limited to single-span bridges for which the second eigenfrequency (in 2D analysis) commonly falls above 30 Hz .

So to what extent is the SDOF model applicable as a guideline or simplification in a real life design? The answer depends on the benefits of having a model that quickly can calculate an approximate response. According to the authors the use of the SDOF model as a guideline is ineffective compared to graphical alternatives like the design curves developed in this thesis. Since the model also is approximate it is difficult to implement it as a simplification of the advanced analysis, even though it could be used to verify complex models by limiting them to the first bending mode.

According to the authors the main use of the SDOF model is that the analytical expressions behind it gives an increased understanding of the influence from material, geometric and load parameters on the dynamic response. This is also the main purpose which it has served in previous master thesis. It might be possible that the added complexities when considering 3D analysis could be explained by the SDOF model. All the complications in the model are based on calculating influence lines of the travelling loads. If a SDOF model is be able to describe the dynamic behaviour in 3D depends hence on the possibility of calculating the deflection in 3D analytically.

The examinations of acceleration amplitude aimed to map the influence of geometric and material parameters. De Leon and Lasn (2008) encountered the accelerations inverse proportionality to the mass. They could however only state that there was an approximate inverse proportionality since they used their SDOF model for all calculations throughout the thesis, which as mentioned include accuracy errors compared to the response from a FE analysis. The choice in this thesis to use a FE program for the examinations of individual parameters influence on the dynamic response was made to avoid these errors. With the use of FE analysis it was stated in this thesis that both young's modulus and the density and hence also cross-section geometry has a simple influence on both eigenfrequencies and acceleration amplitude in 2D analysis.

The first idea of using a graphical tool as guideline in the design process was made by De Leon and Lasn (2008). Their idea was to create transformation graphs based on their observation of the accelerations approximate inverse proportionality to the bridge mass with the use of the SDOF model. The idea is according to the authors complicated and time consuming. Four different graphs were required for calculating an approximate result and 300 graphs had to be developed for the possibility of describing just a single-span bridge. In this thesis it has been shown that the dynamic response for single-span bridges can be illustrated by one single figure through so called design curves, and that the approach is applicable for any bridge configuration of multiple spans.

Gustavsson (2008) had a different approach compared to De Leon and Lasn (2008) as his thesis focused on developing a FE program in Matlab that performed the dynamic calculations of railway bridges subjected to HSLM-A train loads. Gustavsson meant that, using his program, it was possible to quickly determine the acceleration response for a certain set of bridge parameters, and that this approach was easier and better than the earlier ideas of simplifications through SDOF models and graphical presentation. A FE program for calculating the dynamic response of railway bridges has been created by the authors in this thesis as well. However it is not the authors' opinion that a 2D FE program alone can be used as a good guideline or simplification in the design process of railway bridges, but rather the possibilities of creating graphical calculation tools from it. As the reader is well aware of now the design curves are the way of representation created and preferred by the authors. A graphical presentation gives more than just a value of the acceleration outcome, it gives an overview and understanding of the response. As been shown the design curves can be used to easily calculate the acceleration for a certain set of bridge parameters and give guidance to a suitable choice of these parameters.

Calculating the design acceleration and giving guidance to a suitable choice of bridge parameters are the main areas of use for the design curves. However in the creation of
design curves the responses from the ten HSLM-A train loads is calculated individually, and the comparison between these individual responses serves an additional purpose as the critical load can be determined from it. In this thesis design curves were calculated for a single- two- and three-span bridge with equal span lengths. It was found that there is a large variation in what HSLM-A load that governs the acceleration response for these bridge configurations. Based on this observation it can be stated that no general exclusion rule can be established. However for the situation when it is preferable to only include a few loads for a fast approximate calculation, it was found that A2 and A9 most frequently governed the design acceleration.

### 6.2 Analysis in 3D

In 3D analysis eigenmodes will include torsion. As was shown eigenmodes that are governed by bending compared to those governed by torsion are affected differently by changes in cross-section geometry, and especially cross-section width. It is hence impractical to apply the design curves made for 2D analysis on 3D geometry. Too many curves have to be developed if it is not possible to create one single design curve, describing an arbitrary choice of cross-section geometry.

So what possibilities of developing guidelines and simplification for dynamic 3D analysis of railway bridges are there? The first torsion mode was shown to have the same behaviour as the bending modes when plotted against its own $\beta$-value. It showed resonance at the same values and for the width variation the shape of the acceleration response was unaffected. It was mentioned that it might be possible to develop a design curve with consideration to the first torsion mode if the effect from variation of cross-section geometry on the acceleration could be decided (as it was clear that increase in width's effect on the acceleration came from more than the increase in bridge mass). One large disadvantage with such a curve is that it does not include an arbitrary choice of torsion modes, meaning that separate curves would have to be developed for every torsion mode of interest.

The bending modes in 3D analysis can be described by one design curve, as the relation between their corresponding eigenfrequencies in principal is unaffected by cross-section geometry (the height was shown to affect the relation but it was small within the commonly used range of heights). A torsion mode cannot be included in this design curve because the relation between bending and torsion modes' eigenfrequencies changes with the cross-section geometry. For the same reason different torsion modes require individual design curves. However for the case when only one torsion mode is considered it should be possible to use a separate design curve that considers width and load eccentricity to give an estimation of the acceleration response. This estimation based on summarizing the bending modes and the first torsion mode's maximum response in time will always be on the safe side.

Another possibility could be to use the design curves developed in 2D with a safety factor considering the additional 3D effects. The design curves developed in 3D in this thesis showed large similarities with those from 2D. It was seen that the response for the examined single-span bridges was governed by the first bending mode and that distinct changes in the design curves occurred at the same values of $\beta$. The use of the design curves from 2D as a guideline could hence work even if 3D geometry is considered. If a safety factor is used for calculating the design acceleration
corresponding to a certain $\beta$ value the 3D effects could be fully represented by the 2D curves. It is however complicated to establish such a safety factor. More extensive calculations in 3D are required to determine how high the addition in acceleration can be. This addition varies within a large range as already shown by the examination of the two single-span bridges in this thesis. There is a large possibility that the use of a safety factor would hence give results that is too much on the safe side for some sets of bridge parameters. If the factor is made dependent on parameters like load eccentricity this could be avoided, but further studies are required to determine if this is at all possible.

### 6.3 Practical use of design curves

It is the authors' opinion that the design curves are of great guidance in the 2D dynamic design of railway bridges. However the question of when 2D analyses are applicable remains. In this thesis 3D analyses were performed on two single-span bridges and the first torsion mode was included, which is highly governed by load eccentricity. The results basically showed that design curves from 2D analysis are not sufficient for describing the acceleration response for a rectangular cross-section subjected to high eccentric loading.

The case when the load has no eccentricity was not studied in this thesis. It was however still shown that the bending modes from 2D and 3D are very similar (for single-span bridges and rectangular cross-section). If only bending modes are included in the 3D analysis it might hence be possible to use 2D design curves with a small magnification factor to get accurate results for any parameter set. There are however torsion modes that are affected when the load has no eccentricity which has to be studied further.

Irrespective of when 3D effects need to be considered, it is the authors' opinion that the design curves from 2D should be used as a basis and 3D effects considered using correlation rules. This because the 2D design curves are a much more comprehensible presentation of the acceleration response, as they allow arbitrary choice of crosssection geometry. As mentioned the 2D design curves could be used for both calculating the design acceleration and as guidance when only bending modes need to be considered. The impact of 3D effects then depends on the individual torsion modes contribution, but also which torsion modes that needs to be considered.

The question of what eigenmodes that should be included could be both beneficial and unbeneficial for consideration to 3D effects. The 3D response studied in this thesis for example would become even more complicated if the code demands that an additional torsion mode should be considered. But in the situation where no torsion mode needs to be considered, which might be the case when there is no load eccentricity, the 2D design curves could be used directly.

### 6.4 Advanced dynamic analysis

The development of guidelines and simplifications that completely exclude the time consuming advanced dynamic analysis requires more examinations and experience of the dynamic response of railway bridges. As the analysis in 3D showed several additional complications compared to 2D analysis, such simplifications are still not in
reach, and might never be. However the advanced analysis can be gradually simplified. One such simplification is to use adequate software. As mentioned in this thesis, commercial software (like ADINA) might be unfitted for dynamic analysis of railway bridges as this analysis includes a combination of travelling point loads. Problems might for example arise when choosing the mesh size and time step for the numerical integration with consideration to a range of train velocities. The Matlab code presented in Appendix D is a good alternative to commercial software and it is preferred by the authors for iterative calculations on railway bridges. The use of a FE program that can produce a design curve also has the benefit of giving in overview of the dynamic response, which at least shows the effect of varying material parameters and design train velocity.

## 7 Concluding remarks

The conclusions from this master thesis and thoughts about further studies are presented in this chapter.

### 7.1 Conclusions

Previous master theses regarding dynamic design of railway bridges, carried out at Reinertsen Sweden AB during recent years, have all focused on single-span railway bridges. This thesis has had a more general approach, as the response of railway bridges with multiple spans also has been included. Examinations have therefore been made on the eigenmodes of multi-span bridges. It was shown to be complicated to derive the eigenfrequencies for multi-span bridges analytically. But through numerical examinations it was shown that the eigenfrequencies can be calculated through an equation that is very similar to that of single-span bridges, with the addition of so called eigenfrequency functions. These functions depend on the span relation and differ between eigenmodes.
A transformation procedure called the force scaling approach was applied on multispan bridges as a continuation of previous master theses. The method transforms a railway bridge of interest into a SDOF model. The method have earlier been shown to give accurate results for single-span bridges and been said to be compatible with more complex structures. However, it was shown in this thesis that it lacks some accuracy for multi-span bridges. The method can still be used to approximate the dynamic response, but it will be on the unsafe side. The model was not used further in the thesis because of the lack in accuracy.

A parameter study has been performed in this thesis, with the purpose of examining the individual effect on the acceleration response due to different bridge parameters. It was found that material parameters affect the acceleration in a predictable way. The amplitude of the acceleration, for a specific frequency ratio (the ratio between load frequency and eigenfrequency), is inverse proportional against the density and independent of Young's modulus. The total bridge length and span relations were found to affect the acceleration in a way that is hard to predict, though it was seen that the accelerations tend to decrease with increased bridge length and increase for small asymmetry between spans.
During the parameter study several observations were made concerning resonance in railway bridges. It was seen that the design acceleration does not always occur in the resonance peek corresponding to the highest train velocity. It was also seen that resonance effects may disappear for some combinations of geometry, through so called cancellation effects. The summarising conclusion from these observation is that it is difficult, in advance, to determine which train speed that will govern the response, and hence the whole speed range suggested by BV Bro (Banverket, 2006) need to be considered.

The observations regarding resonance and individual parameters effect on the accelerations were gathered to create a graphical guideline, called design curves. A design curve captures the dynamic response of a railway bridge with a certain bridge configuration. It allows an arbitrary choice of train speed, material parameters and cross-section geometry and includes the effect of all ten HSLM-A train loads. It was shown that such design curves can be used to quickly calculate the design acceleration for a specific set of parameters. Also it serves as guidance in the early stage of a
project, as the graphical form gives an overview of how the response is affected by changes in bridge parameters.

The design curves main limitation is that they are based on 2D analysis. In this thesis a comparison between 2D analysis and 3D analysis was performed to show the difference in acceleration response. Calculations were made on two single-span bridges under the variation of bridge width and load eccentricity. In 2D analysis two bending modes were included and in 3D analysis one additional torsion mode was also included. The results showed that both eigenfrequencies and the acceleration response are affected differently for bending and torsion modes. The eigenfrequency of bending modes are in principal unaffected by width variation while width strongly affects the torsion modes' eigenfrequencies. The torsion modes acceleration response is also highly governed by load eccentricity, while the bending modes' response is not. The differences between eigenmodes in a 3D analysis make it more difficult to find guidelines and simplifications compared to 2D analysis. However it can be concluded that 3D effects need to be considered somehow, as it was shown that they may give a very high contribution when large load eccentricity is present. In this thesis the possibilities of producing separate design curves or magnification factors have been discussed, though additional studies are required to determine these possibilities.

### 7.2 Suggestions of continued work

The question of how many eigenmodes that should be included in the dynamic design of railway bridges (based on eigenfrequencies) definitely needs more attention. Developing guidelines for railway bridges, of any kind, whether it is a simplified system or a design curve requires that the number of included eigenmodes is known. In this thesis, choices were made for how many eigenmodes to include for the three bridge configurations of which design curves where created in 2D. However for the development of design curves, to be used in real life design, these choices need to be made on a larger basis.

The knowledge of which torsion modes that needs to be considered for a specific set of bridge parameters is essential when considering 3D effects. This knowledge itself might serve as guidance in the design process if it can be shown that all torsion modes can be excluded for some sets of bridge parameters. Knowing how many eigenmodes that should be included is hence an important step for the possibility of creating guidelines for the dynamic design of railway bridges considering both 2D and 3D.

To determine the required number of eigenmodes for different bridge configurations, it is suggested by the authors that a study which gather used cross-section geometry from railway bridges designed for HSLM train loads is performed. It might also be possible to use the guidance of the design curves to create a range of suitable crosssections. However, this creates a circular reasoning as the result from the design curves depend on the included eigenmodes and vice versa.

Even if the torsion modes to include in a railway bridge design are known, there is still a question of determining these modes' contribution to the acceleration response in a simplified way. The torsion modes individual contribution could be presented by individual design curves or magnification factors as mentioned in Section 6.3. It might also be possible to develop simplified models like a rod subjected to torsion for the first torsion mode or a simply supported plate for the second torsion mode. As
different torsion modes behave in different ways it is likely that different factors/curves/models need to be developed for the respective modes.

There is also a need to increase the understanding of how both bending and torsion modes are affected by changes in cross-section shape. It was shown in this thesis that bending modes in 3D compared to 2D had small differences, but the 3D calculations in this thesis were also limited to rectangular cross-sections. For other cross-section shapes there might be larger differences between the bending modes, which then would cause larger differences between 2D and 3D analysis irrespective of the inclusion of torsion modes.

To summarize the authors would like to see further studies regarding which eigenmodes to include, individual effect of torsion modes and changes in eigenmodes regarding cross-section shape in 3D. An idea could be to examine all three areas, but for one specific bridge configuration at a time. This way useful guidance for real life projects could be continuously produced.

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## A. Appendix A

## Attempt of analytically deriving an explicit expression for the eigenfrequencies of a two-span bridge

The procedure of deriving an analytical expression for the eigenfrequencies of a single span bridge is rather simple and results in an explicit expression for eigenfrequencies of any degree. For two span bridges it gets more complicated however and an explicit expression is impossible to derive. This appendix is meant to give the reader an idea of why.

The differential equation of motion for transverse vibration is shown by equation (A.1).
$\frac{\partial^{2}}{\partial x^{2}} *\left(E I \frac{\partial^{2} v}{\partial x^{2}}\right)+\rho A \frac{\partial^{2} v}{\partial t^{2}}=p_{y}(x, t)$
For free vibrations with constant bending stiffness equation (A.1) reduces to:
$E I v^{\prime \prime \prime \prime}+\rho A \ddot{v}=0$
If we assume the following solution to equation (A.2):
$v(x, t)=V(x) * \cos (\omega t-\alpha)$
We get:
$\frac{d^{4} V}{d x^{4}}-\lambda^{4} V=0$
Where:
$\lambda^{4}=\omega^{2} \frac{\rho A}{E I}$
The general solution to (A.4) can be written as:

$$
\begin{equation*}
V(x)=A_{1} \sinh (\lambda x)+A_{2} \cosh (\lambda x)+A_{3} \sin (\lambda x)+A_{4} \cos (\lambda x) \tag{A.6}
\end{equation*}
$$

This equation has four variables which must be determined. For a single span bridge we could simply use the boundary condition $V=d^{2} V / d x^{2}=0$ at both ends. This gives us that $\mathrm{A}_{2}=\mathrm{A}_{4}=0$ and the two remaining linear equations gives us that:

$$
\left|\begin{array}{cc}
\sin (\lambda L) & \sin (\lambda L)  \tag{A.7}\\
\lambda^{2} \sin (\lambda L) & -\lambda^{2} \sin (\lambda L)
\end{array}\right|=0
$$

Which because of the simplicity of the matrix gives us an explicit expression for $\lambda=$ $r * \pi / L$ for $r=1,2,3$ and so on.

For a two span bridge we have two beams connected together. This means we instead have eight variables and eight boundary conditions. The boundary conditions at the middle support will connect the two beams together.

Figure (A-1) defines two different parameters along the bridge. The definitions of these parameters are made to give as simple result as possible.


Figure A-1 Definition of the variables $x$ and $y$
We have the following general solutions for the displacement of beam 1 and 2 :
$V_{1}(x)=A_{1} \sinh (\lambda x)+A_{2} \cosh (\lambda x)+A_{3} \sin (\lambda x)+A_{4} \cos (\lambda x)$
$V_{2}(x)=B_{1} \sinh (\lambda y)+B_{2} \cosh (\lambda y)+B_{3} \sin (\lambda y)+B_{4} \cos (\lambda y)$
With the definition of $x$ and $y$ we have the following boundary conditions:
$V_{1}(0)=V^{\prime \prime}{ }_{1}(0)=0 \quad V_{2}(0)=V^{\prime \prime}{ }_{2}(0)=0$
$V_{1}(L 1)=0 \quad V_{1}(L 2)=0$
$V^{\prime}{ }_{1}(L 1)=-V^{\prime}{ }_{2}(L 2) V^{\prime \prime}{ }_{1}(L 1)=V^{\prime \prime}{ }_{2}(L 2)$
This gives us that:
$A_{2}=A_{4}=B_{2}=B_{4}$
And the following linear equations give us that:
$\left|\begin{array}{cccc}\sinh (\lambda L 1) & 0 & \sin (\lambda L 1) & 0 \\ 0 & \sinh (\lambda L 2) & 0 & \sin (\lambda L 2) \\ \cosh (\lambda L 1) & \cos (\lambda L 1) & \cosh (\lambda L 2) & \cos (\lambda L 2) \\ \sinh (\lambda L 1) & -\sin (\lambda L 1) & -\sinh (\lambda L 2) & \sin (\lambda L 2)\end{array}\right|=0$

And this results in a complicated expression where there is no possibility of deriving an explicit analytical expression for $\lambda$.

The last steps can be better performed. The first two rows in (12) could be used to eliminate two variables. But even if we get that the determinant should be zero for a two times two matrix we still get a very complicated expression. The interested reader is referred to Gorman (1975).

## B. Appendix B

## Results from a parameter study on multi-span railway bridges

The following document contains figures showing the results from a parameter study made on continuous railway bridges. The parameter study covers two- and three-span bridges.

The parameter study aims to show the effect from different parameters on the dynamic response. The studied parameters are:

- $\mathrm{E}=$ Module of elasticity
- Ltot $=$ Total bridge length
- $\mu=\quad$ L2/L1 for two-span bridges
- $\eta=\quad \mathrm{L} 2 /(\mathrm{L} 1+\mathrm{L} 3)$ for three-span bridges
- $\rho=$ density
- $\mathrm{c}=$ damping

Two different types of figures are presented. The first type shows different acceleration responses with the variation of train speed. The second type shows the variation of acceleration amplitude around resonance peeks.
The corresponding parameter set for the variation of a certain variable are presented in the figure explanation. Different parameter sets have been chosen for different comparisons.

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## B. 1 Variation of module of elasticity, E, for a two-span bridge



Figure B-1 Acceleration response for a two-span bridge subjected to HSLMA1 using $\rho=2800 \mathrm{~kg} / \mathrm{m}^{3}$, Ltot $=16 \mathrm{~m}, \mu=2$ and $c=0.02$


Figure B-2
Acceleration response for a two-span bridge subjected to HSLM-A4 using $\rho=2800 \mathrm{~kg} / \mathrm{m3}$, Ltot $=16 \mathrm{~m}, \mu=2$ and $c=0.02$

## B. 2 Variation of density, $\rho$, for a two-span bridge



Figure B-3 Acceleration response for a two-span bridge subjected to HSLMA1 using $E=30$ GPa, Ltot $=16 \mathrm{~m}, \mu=2$ and $c=0.02$

Figure B-4


Figure B-5 Acceleration from figure B-4 plotted againsst the inverse of the density

## B. 3 Variation of span relation, $\mu$, for a two-span bridge



Figure B-6
Acceleration response for a two-span bridge subjected to HSLM-A1 using $E=20 \mathrm{GPa}$, Ltot $=16 \mathrm{~m}, \rho=3000 \mathrm{~kg} / \mathrm{m}^{\wedge} 3$ and $c=0.02$


Figure B-8
Acceleration response for a two-span bridge subjected to HSLM-A4 using $E=20 \mathrm{GPa}$, Ltot $=16 \mathrm{~m}, \rho=4000 \mathrm{~kg} / \mathrm{m}^{\wedge} 3$ and $c=0.02$


Figure B-7
Acceleration response for a two-span bridge subjected to HSLM-A1 using $E=20 \mathrm{GPa}$, Ltot $=16 \mathrm{~m}, \rho=3000 \mathrm{~kg} / \mathrm{m}^{\wedge} 3$ and $c=0.02$


Figure B-9
Acceleration response for a two-span bridge subjected to HSLM-A4 using $E=20 \mathrm{GPa}$, Ltot $=16 \mathrm{~m}, \rho=4000 \mathrm{~kg} / \mathrm{m}^{\wedge} 3$ and $c=0.02$


Figure B-10 Acceleration response for a two-span bridge subjected to HSLM-A8 using $E=20 \mathrm{GPa}$, Ltot $=16 \mathrm{~m}, \rho=6000 \mathrm{~kg} / \mathrm{m}^{\wedge} 3$ and $c=0.02$


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Figure B-11 Acceleration response for a two-span bridge subjected to HSLM-A8 using $E=20 \mathrm{GPa}$, Ltot $=16 \mathrm{~m}, \rho=6000 \mathrm{~kg} / \mathrm{m}^{\wedge} 3$ and $c=0.02$


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Acceleration response for a two-span bridge subjected to HSLM-A1 using $E=40 \mathrm{GPa}, \mu=2, \rho=4000 \mathrm{~kg} / \mathrm{m}^{\wedge} 3$ and Ltot $=15 \mathrm{~m}$

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Figure B-35
Acceleration response for a three-span bridge subjected to HSLM-A1 using $E=40 \mathrm{GPa}, \mu=1$, Ltot $=25 \mathrm{~m}, ~, \kappa=1$ and $c=0.02$

## B. 7 Variation of span relation, $\eta$, for a three-span bridge



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Figure B-38
Acceleration response for a three-span bridge subjected to HSLM-A8 using $E=10 \mathrm{GPa}, L=20 \mathrm{~m}, \rho=2500 \mathrm{~kg} / \mathrm{m}^{\wedge} 3, \kappa=1$ and $c=0.02$


Figure B-37 Acceleration response for a three-span bridge subjected to HSLM-A4 using $E=14 \mathrm{GPa}, L=20 \mathrm{~m}, \rho=2500 \mathrm{~kg} / \mathrm{m}^{\wedge} 3, \kappa=1$ and $c=0.02$


Figure B-39
Acceleration response for a three-span bridge subjected to HSLM-A8 using $E=10 \mathrm{GPa}, \mathrm{L}=20 \mathrm{~m}, \rho=2500 \mathrm{~kg} / \mathrm{m}^{\wedge} 3, \kappa=1$ and $c=0.02$


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Acceleration response in resonance peek at vcr/4 from the response shown by figures B-36-39


Figure B-45
Normalized acceleration response in resonance peek at vcr/4 from the response shown by figures B-36-39

## C. Appendix C

The design curves from the three different case studies are gathered in this appendix. The different bridge configurations are:

- Single-span bridge - Section C. 1
- Two-span bridge with equal span lengths - Section C. 2
- Three-span bridge with equal span lengths - Section C. 3

Further description and explanation of the three case studies can be found in section 3.6.4. The layout of each section in this appendix is first a summation of the design curves for all bridge lengths considered in the each case study, followed by the design curves for each specific bridge length with the possibility to see which load that governs the response at a specific choice of $\beta$.
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## C. 1 Bridge with one span



Figure C-1 Design curves of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=7 \mathrm{~m}$ to $L_{\text {tot }}=20 \mathrm{~m}$.


Figure C-2 Design curves of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=7 \mathrm{~m}$ to $L_{\text {tot }}=20 \mathrm{~m}$.

## C.1.1 Total length 7 m



Figure C-3 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with one span and $L_{\text {tot }}=7 \mathrm{~m}$.


Figure C-4 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=7 \mathrm{~m}$.


Figure C-5 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=7 \mathrm{~m}$.

## C.1.2 Total length 8 m



Figure C-6 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with one span and $L_{\text {tot }}=8 \mathrm{~m}$.


Figure C-7 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=8 \mathrm{~m}$.


Figure C-8 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=8 \mathrm{~m}$.

## Total length $9 \mathbf{m}$



Figure C-9 Relation between $\beta$ and $\tau$ for load Al to A10 for a bridge with one span and $L_{\text {tot }}=9 \mathrm{~m}$.


Figure C-10 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=9 \mathrm{~m}$.


Figure C-11 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=9 \mathrm{~m}$.

## C.1.3 Total length 10 m



Figure C-12 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with one span and $L_{\text {tot }}=10 \mathrm{~m}$.


Figure C-13 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=10 \mathrm{~m}$.


Figure C-14 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=10 \mathrm{~m}$.

## C.1.4 Total length 11 m



Figure C-15 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with one span and $L_{\text {tot }}=11 \mathrm{~m}$.


Figure C-16 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=11 \mathrm{~m}$.


Figure C-17 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=11 \mathrm{~m}$.

## C.1.5 Total length 12 m



Figure C-18 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with one span and $L_{\text {tot }}=12 \mathrm{~m}$.


Figure C-19 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=12 \mathrm{~m}$.


Figure C-20 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=12 \mathrm{~m}$.

## C.1.6 Total length 13 m



Figure C-21 Relation between $\beta$ and $\tau$ for load Al to A10 for a bridge with one span and $L_{\text {tot }}=13 \mathrm{~m}$.


Figure C-22 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=13 \mathrm{~m}$.


Figure C-23 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=13 \mathrm{~m}$.

## C.1.7 Total length 14 m



Figure C-24 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with one span and $L_{\text {tot }}=14 \mathrm{~m}$.


Figure C-25 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=14 \mathrm{~m}$.


Figure C-26 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=14 \mathrm{~m}$.

## C.1.8 Total length 15 m



Figure C-27 Relation between $\beta$ and $\tau$ for load Al to A10 for a bridge with one span and $L_{\text {tot }}=15 \mathrm{~m}$.


Figure C-28 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=15 \mathrm{~m}$.


Figure C-29 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=15 \mathrm{~m}$.

## C.1.9 Total length 16 m



Figure C-30
Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with one span and $L_{t o t}=16 \mathrm{~m}$.


Figure C-31 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=16 \mathrm{~m}$.


Figure C-32 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=16 \mathrm{~m}$.

## C.1.10 Total length 17 m



Figure C-33 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with one span and $L_{\text {tot }}=17 \mathrm{~m}$.


Figure C-34 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=17 \mathrm{~m}$.


Figure C-35 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=17 \mathrm{~m}$.

## C.1.11 Total length 18 m



Figure C-36 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with one span and $L_{\text {tot }}=18 \mathrm{~m}$.


Figure C-37 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=18 \mathrm{~m}$.


Figure C-38 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=18 \mathrm{~m}$.

## C.1.12 Total length 19 m



Figure C-39
Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with one span and $L_{\text {tot }}=19 \mathrm{~m}$.


Figure C-40 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=19 \mathrm{~m}$.


Figure C-41 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=19 \mathrm{~m}$.

## C.1.13 Total length 20 m



Figure C-42
Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with one span and $L_{\text {tot }}=20 \mathrm{~m}$.


Figure C-43 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with one span and $L_{\text {tot }}=20 \mathrm{~m}$.


Figure C-44 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with one span and $L_{\text {tot }}=20 \mathrm{~m}$.

## C. 2 Bridge with two equal spans



Figure C-45 Design curves of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with two equal spans and $L_{\text {tot }}=16 \mathrm{~m}$ to $L_{\text {tot }}=30 \mathrm{~m}$.


Figure C-46 Design curves of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with two equal spans and $L_{\text {tot }}=16 \mathrm{~m}$ to $L_{\text {tot }}=30 \mathrm{~m}$.

## C.2.1 Total length 16 m



Figure C-47 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with two equal spans and $L_{\text {tot }}=16 \mathrm{~m}$.

$\beta$ [-]
Figure C-48 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with two equal spans and $L_{\text {tot }}=16 \mathrm{~m}$.


Figure C-49 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with two equal spans and $L_{\text {tot }}=16 \mathrm{~m}$.

## C.2.2 Total length 18 m



Figure C-50 Relation between $\beta$ and $\tau$ for load A1 to Al0 for a bridge with two equal spans and $L_{\text {tot }}=18 \mathrm{~m}$.


Figure C-51 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with two equal spans and $L_{\text {tot }}=18 \mathrm{~m}$.


Figure C-52 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with two equal spans and $L_{\text {tot }}=18 \mathrm{~m}$.

## C.2.3 Total length 20 m



Figure C-53 Relation between $\beta$ and $\tau$ for load A1 to Al0 for a bridge with two equal spans and $L_{\text {tot }}=20 \mathrm{~m}$.


Figure C-54 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with two equal spans and $L_{\text {tot }}=20 \mathrm{~m}$.


Figure C-55 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with two equal spans and $L_{\text {tot }}=20 \mathrm{~m}$.

## C.2.4 Total length 22 m



Figure C-56 Relation between $\beta$ and $\tau$ for load A1 to Al0 for a bridge with two equal spans and $L_{\text {tot }}=22 \mathrm{~m}$.


Figure C-57 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with two equal spans and $L_{\text {tot }}=22 \mathrm{~m}$.


Figure C-58 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with two equal spans and $L_{\text {tot }}=22 \mathrm{~m}$.

## C.2.5 Total length 24 m



Figure C-59 Relation between $\beta$ and $\tau$ for load A1 to Al0 for a bridge with two equal spans and $L_{\text {tot }}=24 \mathrm{~m}$.


Figure C-60 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with two equal spans and $L_{\text {tot }}=24 \mathrm{~m}$.


Figure C-61 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with two equal spans and $L_{\text {tot }}=24 \mathrm{~m}$.

## C.2.6 Total length 26 m



Figure C-62 Relation between $\beta$ and $\tau$ for load A1 to Al0 for a bridge with two equal spans and $L_{\text {tot }}=26 \mathrm{~m}$.


Figure C-63 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with two equal spans and $L_{\text {tot }}=26 \mathrm{~m}$.


Figure C-64 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with two equal spans and $L_{\text {tot }}=26 \mathrm{~m}$.

## C.2.7 Total length 28 m



Figure C-65 Relation between $\beta$ and $\tau$ for load A1 to Al0 for a bridge with two equal spans and $L_{\text {tot }}=28 \mathrm{~m}$.


Figure C-66 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with two equal spans and $L_{\text {tot }}=28 \mathrm{~m}$.


Figure C-67 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with two equal spans and $L_{\text {tot }}=28 \mathrm{~m}$.

## C.2.8 Total length 30 m



Figure C-68 Relation between $\beta$ and $\tau$ for load A1 to Al0 for a bridge with two equal spans and $L_{\text {tot }}=30 \mathrm{~m}$.

$\beta$ [-]
Figure C-69 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.5$, for a bridge with two equal spans and $L_{\text {tot }}=30 \mathrm{~m}$.


Figure C-70 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with two equal spans and $L_{\text {tot }}=30 \mathrm{~m}$.

## C. 3 Bridge with three equal spans



Figure C-71 Design curves of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with three equal spans and $L_{\text {tot }}=30 \mathrm{~m}$ to $L_{\text {tot }}=43.5 \mathrm{~m}$.


Figure C-72 Design curves of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with three equal spans and $L_{\text {tot }}=30 \mathrm{~m}$ to $L_{\text {tot }}=43.5 \mathrm{~m}$.

## C.3.1 Total length 30 m



Figure C-73 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with three equal spans and $L_{\text {tot }}=30 \mathrm{~m}$.

$\beta$ [-]
Figure C-74 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.4$, for a bridge with three equal spans and $L_{\text {tot }}=30 \mathrm{~m}$.


Figure C-75 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with three equal spans and $L_{\text {tot }}=30 \mathrm{~m}$.

## C.3.2 Total length 31.5 m



Figure C-76 Relation between $\beta$ and $\tau$ for load A1 to Al0 for a bridge with three equal spans and $L_{\text {tot }}=31.5 \mathrm{~m}$.

$\beta$ [-]
Figure C-77 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.4$, for a bridge with three equal spans and $L_{\text {tot }}=31.5 \mathrm{~m}$.


Figure C-78 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with three equal spans and $L_{\text {tot }}=31.5 \mathrm{~m}$.

## C.3.3 Total length 33 m



Figure C-79 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with three equal spans and $L_{\text {tot }}=33 \mathrm{~m}$.


$\beta$ [-]
Figure C-80 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.4$, for a bridge with three equal spans and $L_{\text {tot }}=33 \mathrm{~m}$.


Figure C-81 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with three equal spans and $L_{\text {tot }}=33 \mathrm{~m}$.

## C.3.4 Total length 34.5 m



Figure C-82 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with three equal spans and $L_{\text {tot }}=34.5 \mathrm{~m}$.

$\beta$ [-]
Figure C-83 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.4$, for a bridge with three equal spans and $L_{\text {tot }}=34.5 \mathrm{~m}$.


Figure C-84 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with three equal spans and $L_{\text {tot }}=34.5 \mathrm{~m}$.

## C.3.5 Total length 36 m



Figure C-85 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with three equal spans and $L_{\text {tot }}=36 \mathrm{~m}$.

$\beta$ [-]
Figure C-86 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.4$, for a bridge with three equal spans and $L_{\text {tot }}=36 \mathrm{~m}$.


Figure C-87 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with three equal spans and $L_{\text {tot }}=36 \mathrm{~m}$.

## C.3.6 Total length 37.5 m



Figure C-88 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with three equal spans and $L_{\text {tot }}=37.5 \mathrm{~m}$.

$\beta$ [-]
Figure C-89 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.4$, for a bridge with three equal spans and $L_{\text {tot }}=37.5 \mathrm{~m}$.


Figure C-90 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with three equal spans and $L_{\text {tot }}=37.5 \mathrm{~m}$.

## C.3.7 Total length 39 m



Figure C-91 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with three equal spans and $L_{\text {tot }}=39 \mathrm{~m}$.


Figure C-92 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.4$, for a bridge with three equal spans and $L_{\text {tot }}=39 \mathrm{~m}$.


Figure C-93 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with three equal spans and $L_{\text {tot }}=39 \mathrm{~m}$.

## C.3.8 Total length 40.5 m



Figure C-94 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with three equal spans and $L_{\text {tot }}=40.5 \mathrm{~m}$.

$\beta[-]$
Figure C-95 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.4$, for a bridge with three equal spans and $L_{\text {tot }}=40.5 \mathrm{~m}$.


Figure C-96 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with three equal spans and $L_{\text {tot }}=40.5 \mathrm{~m}$.

## C.3.9 Total length 42 m



Figure C-97 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with three equal spans and $L_{\text {tot }}=42 \mathrm{~m}$.

$\beta$ [-]
Figure C-98 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.4$, for a bridge with three equal spans and $L_{\text {tot }}=42 \mathrm{~m}$.


Figure C-99 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with three equal spans and $L_{\text {tot }}=42 \mathrm{~m}$.

## C.3.10 Total length 43.5 m



Figure C-100 Relation between $\beta$ and $\tau$ for load A1 to A10 for a bridge with three equal spans and $L_{\text {tot }}=43.5 \mathrm{~m}$.


Figure C-101 Relation between $\beta$ and $\tau$ for load A1 to A10, with a zoom at $\beta=0.2$ to $\beta=0.4$, for a bridge with three equal spans and $L_{\text {tot }}=43.5 \mathrm{~m}$.


Figure C-102 Design curve of the relation between $\beta$ and $\tau$ for the HSLM-A loads of a bridge with three equal spans and $L_{\text {tot }}=43.5 \mathrm{~m}$.

## D. Appendix D Matlab Programs

In this appendix the layout of several Matlab programs used in the dynamic analysis of railway bridges will be presented. In the first three sections "main programs" are presented. With "main program" the authors refer to a program which purpose mainly is to gather the calculations from different function files to achieve a sought result. Used function files are presented in section D.4.

Each section starts with an overview of the created program and a description of the area of use for the program. In each section the code for the specific program will be printed and available to be copied into Matlab. Verifications for some of the programs are presented in the end of the sections.

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## D. 1 SDOF

This section will present the "main program" for the transformation of a railway bridge into a SDOF model using the force scaling approach.

## D.1.1 About the SDOF program

The force scaling approach is made in three steps. First the mass of the SDOF system is chosen arbitrary. Then the stiffness is calculated to get the first eigenfrequency equal to the first eigenfrequency of the railway bridge. Finally the load is scaled so that the static displacement of the SDOF system and the critical point in the railway bridge become equal.
The scaling of the load requires that the influence line for the displacement in the critical point of the bridge is calculated for a point load moving across the bridge. This is done in a function file called SdofLoad.m, see section D.4.1.
The Newmark- $\beta$ method is used as the numerical time integration method, see section D.4.6.
D.1.2 Code

```
%-----------------------------------------------------------------------------------
% MAIN PROGRAM
% SDOF MODELING OF A RAILWAY BRIDGE
%--------------------------------------------------------------------------------------
%% STARTING DEFINITIONS
clear all
dof = 3;
neli = 20;
%% DEFINITION OF BRIDGE PARAMETERS
L_tot = 15;
kap = 1;
my = 2;
Li = L_tot*[1/(1+my) (1-1/(1+my))];
w=11.2;
h = 0.6;
E = 40*10^9;
Mass = 2400;
A = w*h;
I = w*h^3/12;
ep =[E A I Mass*W*h [0 0]];
%% EIGENFREQUENCIES
[Coord, Edof, Ndof, Elnod, Ex, Ey, nel, non, ndof, b] =...
    MeshTwoD(dof, neli, Li);
K = zeros(ndof);
M = zeros(ndof);
for i=1:nel
    [Ke,Me] = beam2d(Ex(i,:),Ey(i,:),ep);
    K = assem(Edof(i,:),K,Ke);
    M = assem(Edof(i,:),M,Me);
end
[L,X] = eigen(K,M,b);
%% MASS, STIFFNESS AND DAMPING OF THE SDOF SYSTEM
M_sdof = 1;
```

```
K_sdof = M_sdof * L(1);
damp = 0.02;
C_sdof = 2*damp*M_sdof*sqrt(L(1));
speed_step = 2.5/3.6;
max_speed = 100;
q = waitbar(0, 'Jag är en muffin..');
for i=1:((max_speed-28)/speed_step)
    speed = 28+i*speed_step; %Speed in m/s
    ntimes = 10000;
    [F_sdof,K_bridge,h]= SdofLoad(speed,ntimes,Li,E,I, 4);
    %[F_sdof,K_bridge,h]= SdofLoadSimple(speed,ntimes,Li,E,I,1);
    S = K_sdof/K_bridge;
    F = S*F_sdof;
    u0 = 0;
    v0 = 0;
    bet = 0.25;
    gam = 0.5;
    [u,v,a,t] = NEWMARK(K_sdof,C_sdof,M_sdof,F,h,u0,v0,bet,gam);
    A_max(i) = max(a);
    V(i) = speed;
    waitbar(i/((max_speed-28)/speed_step))
end
close(q)
%% PLOT RESULTS
hold on
%plot(t,F)
plot(V.*3.6,A_max,'r')
%plot(t,a)
%plot(t,F)
```


## D. 2 Acceleration response in 2D

This section will present a "main program" for Matlab for calculating the acceleration response of a continuous railway bridge in 2D. The program has several areas of use. The main use is to calculate the maximum vertical acceleration in time in the railway bridge for the variation of train velocity and bridge parameters. But as several calculation procedures are required to achieve this it can also be used to produce the result of each separate calculation. With small modifications of the program code presented later it is possible to extract eigenfrequencies, eigenmodes and the time response concerning displacement, velocity and accelerations.

## D.2.1 About the FE program

The program calculates the acceleration response using FE analysis. The analysis is based on mode-superposition, but numerical time integration is still used because of complexity in the train load. The calculation procedure can be divided in the following steps:

- Meshing
- Assembly of stiffness and mass matrix
- Calculation of eigenfrequencies, eigenmodes, modal mass matrix, modal stiffness matrix and modal damping
- Load creation
- Numerical time integration
- Data saving and plotting

The meshing is made using a function file called MeshTwoD.m. In short the mesh file creates a 2D beam with the possibility of choosing the number of spans, span lengths, degrees of freedom in each node and the number of elements in each span. The reader is referred to section D.4.4 for more information on the mesh function file.

Creation of each elements local stiffness and mass matrix and assembly into the global stiffness matrices is made using two Calfem files called beam2d.m and assem.m. The reader is referred to CALFEM (Austell P.-E., et. Al., 2004) for more information on Calfem and the used function files.

Calculation of eigenfrequencies and eigenmodes is also made using a Calfem file, called eigen.m. This file uses the built in eigensolver in Matlab but orders the output data in a correct way for the use of mode-superposition. It also normalizes the eigenmodes to get a modal mass matrix equal to a diagonal of ones.

Calculation of the train load is made using a function file called TrainTwoD.m. The function file calculates the load matrix representing the load in each node for every time step for an arbitrary HSLM-A load. The file has the possibilities of easily switching between different HSLM-A train loads and adjusting these loads to any choice of number of elements and time step length. The reader is referred to section D.4.2 for more information on the load function file.

The time integration function file uses the Newmark- $\beta$ method with the parameters $\beta$ and $\gamma$ equal to 0.25 and 0.5 respectively to make it unconditionally stable. The time integration has been made using the function file Newmark.m, see D.4.6.

## D.2.2 Matlab code



```
[Coord, Edof, Ndof, Elnod, Ex, Ey, nel, non, ndof, b] =...
    MeshTwoD(dof, neli, Li);
%ASSEMBLY OF BEAM ELEMENTS
%The Calfem function file beam2d is used for creating the local
%stiffness matrices and these are assembled using the Calfem file
%assem
K = zeros(ndof);
M = zeros(ndof);
for i=1:nel
    [Ke,Me] = beam2d(Ex(i,:),Ey(i,:),ep);
    K = assem(Edof(i,:),K,Ke);
    M = assem(Edof(i,:),M,Me);
end
%EIGENVALUES AND EIGENMODES
[L,X] = eigen(K,M,b);
%Cropping of the stiffness, mass and eigenvector matrices with
%consideration to boundary conditions.
K (b, : ) = [];
K (:,b)=[];
M(b,:)=[];
M(:,b)=[];
X(b,:) = [];
%MODE SUPERPOSITION
%number of eigenmodes used
nmodes = 2;
%modal mass and stiffness matrix
M_diag = X(:,1:nmodes)'*M*X(:,1:nmodes);
K_diag = X(:,1:nmodes)'*K*X(:,1:nmodes);
%modal damping
damp = 0.02;
C_diag = diag(2*damp*sqrt(L(1:nmodes)));
%CONSIDERED RANGE OF TRAIN VELOCITY
speed_step = 1.25/3.6;
max_speed = 315;
%Loop that calculates the response in time using mode superposition
%for a chosen range of train velocities
q = waitbar(0, ['Iteration ' num2str(j)]);
for i=1:((max_speed)/speed_step)
    %Train velocity
    speed = i*speed_step;
    %Number of time steps
    ntimes = 10000;
    %Starting displacement and velocity
    u0 = zeros(nmodes,1);
    v0 = zeros(nmodes,1);
    %TRAIN LOAD
    [F_train] = TrainTwoD(speed,ntimes,ndof,Coord,Elnod,Edof,Li,1);
    F_train = [zeros(length(F_train(:,1)),1) F_train];
    F_train(b,:) = [];
    F = X(:,1:nmodes)'*F_train;
    %TIME INTEGRATION
    bet = 0.25;
    gam = 0.5;
    [u,v,a,t] = NEWMARK(K_diag,C_diag,M_diag,F,h,u0,v0,bet,gam);
    %Calculates the true displacements and accelerations and adds the
    %removed degrees of freedom
    u = X(:,1:nmodes)*u;
```

```
        a = X(:,1:nmodes)*a;
        e=1:1:ndof;
        e(b) = [];
        u_true = zeros(ndof,ntimes+1);
        a_true = zeros(ndof,ntimes+1);
        u_true(e,:) = u;
        a_true(e,:) = a;
        %MAXIMUM ACCELERATION
        Max = max(max(a_true));
        Min = min(min(a_true));
        A_maxmax(i,j) = max([Max abs(Min)]);
        V(i,j) = speed;
        waitbar(i/((max_speed)/speed_step))
    end
    close (q)
    %SAVES USED PARAMETERS FOR EACH ITERATION
    EE(j) = E;
    MASS(j) = Mass;
    L_TOT(j) = L_tot;
    MY(j) = my;
    DAMP(j) = damp;
end
%% PLOT RESULTS
% A huge variation of plot options can be used. The code below is just an
% example of possible plot options
hold on
colors = ['y', 'g', 'b', 'r', 'k'];
for k=1:5
plot(V(:,k)*3.6, A_maxmax(:,k), colors(k))
end
xlabel('Speed [km/h]')
ylabel('Acceleration [m/s^2]')
title('HSLM-A1')
vari = 'E = ';
legend([vari num2str(EE(1))], [vari num2str(EE(2))],...
    [vari num2str(EE(3))], [vari num2str(EE(4))],...
    [vari num2str(EE(5))])
```


## D.2.3 Verification

Verification of the program has been made using the commercial software ADINA. A 2D model of a two-span bridge has been created in both ADINA and the Matlab FE program with the purpose of comparing the responses from the two programs. The geometry and boundary conditions in the model are shown by Error! Reference source not found.


Figure D-1 Geometry and boundary conditions in the model used for the verification.
The material parameters in the form of module of elasticity and concrete density have been chosen to 30 GPa and $2400 \mathrm{~kg} / \mathrm{m}^{3}$ respectively.

The time integration is made using the Newmark- $\beta$ method with the integration parameters $\beta$ and $\gamma$ equal to 0.25 and 0.5 respectively. 10000 time steps are used in the integration.

Both models use mode-superposition limited to the three eigenmodes with lowest corresponding eigenfrequency. Damping has been included through a modal damping of $2 \%$ for all considered eigenmodes.

Beam elements are used in both programs. One hundred elements in each span have been used in the Matlab program, which corresponds to a length of 0.05 and 0.1 m . The element length in Adina has been chosen to 0.25 in both spans. The reason why the element lengths have been chosen differently between the models is because the Matlab model is built to choose the number of elements equally in every span, while the Adina model gets substantially easier to create with regard to the train load using equal span lengths.

As a first verification the frequencies from the two models are compared, see Table D.1.

Table D. $1 \quad$ Eigenfrequencies for the two models

|  | Frequencies |  |
| :---: | :--- | :--- |
| Mode | ADINA | Matlab |
| 1 | 14.21 Hz | 14.2 Hz |
| 2 | 44.17 Hz | 44.17 Hz |
| 3 | 61.71 Hz | 61.71 Hz |

The HSLM-A load case A1 and A4 have been used to verify the dynamic response of the program. The loads have been simulated with the velocities $180 \mathrm{~km} / \mathrm{h}$ and 270 $\mathrm{km} / \mathrm{h}$ respectively. The displacements and accelerations have been gathered in the middle of the second span. Figure D-2 to Figure D-5 show the comparison of responses between the two programs for the two first seconds.


Figure D-2 Comparison between displacements for the HSLM-A1 train load with velocity of 180 km/h


Figure D-3 Comparison between accelerations for the HSLM-A1 train load with velocity of 180 km/h


Figure D-4 Comparison between displacements for the HSLM-A4 train load with velocity of 270 km/h


Figure D-5
Comparison between accelerations for the HSLM-A4 train load with velocity of 270 km/h

## D. 3 Acceleration response in 3D

The purpose with this "main program" is to calculate the maximum vertical acceleration in time considering a 3D model of railway bridges for the variation of load and bridge parameters. It is also possible to modify the program to calculate eigenfrequencies, eigenmodes and the time response concerning displacement, velocity and accelerations.

## D.3.1 About the FE program

The program calculates the acceleration response for a railway bridge modelled as a simply supported plate with the possibility to use eccentric loading. The analysis uses mode-superposition and numerical time integration to get the vertical acceleration response in time. The program can be divided in several steps:

- Meshing
- Assembly of the stiffness matrix
- Calculation of eigenfrequencies, eigenmodes, modal mass matrix, modal stiffness matrix and modal damping
- Creation of HSLM train loads
- Numerical time integration
- Plot figures and saves data

The function file MeshThreeD.m is used for creating the mesh. The mesh file creates a 3D plate consisting of shell elements. It is possible to choose span length and the number of elements in both length of the span and the width. More information about the mesh function file can be found in section D.4.5.

The Calfem files platre.m, planre and assem.m have been used in the creation of each elements local stiffness and mass matrix and assembly into the global stiffness matrices. Further information about the Calfem files is found in CALFEM - A finite element toolbox version 3.4 (Austell P.-E., et. Al., 2004)

The 3D program uses a different method for calculating the eigenfrequencies and eigenmodes compared to the 2D "main program". Instead of using a Calfem file and the built in eigensolver in Matlab a method called inverse iteration and inverse iteration with shift is used. The method provides the possibility of calculating only a few of the lowest eigenmodes to save calculation time and can be used to choose specific eigenmodes. The gained eigenmodes from the function file Inverse_iteration.m are all normalized so it is possible to get a modal mass matrix equal to a diagonal of ones. The function file is described in detail in section D.4.7.

The train load in this program is defined by the function file TrainThreeD.m. It calculates the load representing the load in each node for every time step for an arbitrary HSLM-A load. It is possible to switch between different HSLM-A train loads and adjust to arbitrary choice of time steps and mesh size. The function file is described further in D.4.3.

As in the 2D program Newmark integration are used for the time integration, for more information see D.4.6.

## D.3.2 Matlab code



```
%% INTRODUCTION
%---------------------------------------------------------------------------
% Program that calculates the response in time using Newmark %
integration.
% Mode superposition is used for fastening the calculation process. %
Meshing, and load creation is made in function files.
%% MAIN CALCULATIONS
% Clear all previous data
clear all
clc
clf
%Indata
qz = 0;
ty = 0;
```

\%Number of elements used in the length of the bridge
n11 = 16;
\%Number of elements used in the width of the bridge
n2 = 16;
\%\% Mesh the geometry in separate file.
\%A loop that allows for the variation of any bridge parameter. Each \%loop
defines a set of material, geometric and load parameters and \%calculates the
response in time for a defined range of train \%velocities. The program can
therefore be used for the response in otime for one specific set of
parameters but also extensive \%calculations like the maximum acceleration in
time under the \%variation of train velocities for a chosen number of
parameter sets.
for $j=1: 10$
\%GEOMETRIC PARAMETERS
\%Total length
$11=15$;
Li = [l1 12];
\%Cross-section width
b=11;
\%Cross-section height
h=0.7;
\%Area
$\mathrm{A}=\mathrm{b}$ *h;
\%Moment of inertia
$I=b * h \wedge 3 / 12$;
\%MATERIAL PARAMETERS
\%Module of elasticity
$\mathrm{E}=40 \mathrm{e} 9$;
\%Density
d=2800;
\%poisson's ratio
$\mathrm{v}=0.2$;
\%MESHING
\%Meshing is made using a separate function file
[B1, B2, B3, B4, P1, P2, P3, P4, Ex, Ey, ndof, nel, non, Edof] =MeshThreeD (11, b, n11, n2)
;

```
%ASSEMBLY OF SHELL ELEMENTS
%The Calfem function files platre and planre are used for %creating the
local stiffness matrices and these are assembled %using the Calfem
function file assem
K = zeros(ndof);
f = zeros(ndof,1);
D = hooke(1,E,v);
ep = [1 h];
eq = [0 0 qz]';
for i=1:nel
    ex = Ex(i,:);
    ey = Ey(i,:);
    [Ke1,fe1]=platre(ex,ey,ep(2),D,eq(3));
    [Ke2,fe2]=planre(ex([1 2]),ey([[2 3]),ep,D,eq([1 2]));
    Ke = zeros(20);
    fe = zeros(20,1);
    Ke([11 2 6 6 7 11 12 16 17],[1 [ 2 6 % 7 11 12 16 17])=Ke2;
    Ke([3 4 5 8 8 9 10 13 14 15 18 19 20],[3 [3 4 5 8 9 10 13 14 15 18 19
        20])=Ke1;
    fe([1 1 2 6 7 7 11 12 16 17])=fe2;
    fe([[3 4 5 5 8 9 10 13 14 15 18 19 20])=fe1;
    [K,f] = assem(Edof(i,:),K,Ke,f,fe);
end
%% MASS MATRICE
%Mass for each element
M_el_1 = l1*b*d*h/nel;
%Rotational lumped mass (see ADINA manual)
M_rot_1 = (M_el_1/4)*(1/12)*(h*h);
k1 = [M_el_1/4 M_el_1/4 M_el_1/4 M_rot_1 M_rot_1...
        M_el_1/4 M_el_1/4 M_el_1/4 M_rot_1 M_rot_1...
        M_el_1/4 M_el_1/4 M_el_1/4 M_rot_1 M_rot_1...
        M_el_1/4 M_el_1/4 M_el_1/4 M_rot_1 M_rot_1];
Me_1 = diag(k1);
M = zeros(ndof);
for i=1:nel
    M = assem(Edof(i,:),M,Me_1);
end
%% BOUNDARY CONDITIONS
for i=1:length(B4(:,1))
    bc1((3*i-2):(3*i)) = B4(i,1:3)';
    bc2((2*i-1):(2*i)) = B2(i, 2:3)';
    bc5((2*i-1):(2*i)) = B4(i,2:3)'+5*(n2+1)*n11/2;
end
bc3 = zeros(1,2);
for k=2:(length(B3(:,1))-1)
    b3 = B1 (k,[[1 2 3}])+5*n2/2
    bc3e = [b3; 0 0 0]';
    bc3 = [bc3;bc3e];
end
bc3(1,:) = [];
bc4 = zeros(1,2);
for k=2:(length(B3 (:,1))-1)
    b4 = B1 (k, 4) +5*n2/2;
    bc4e = [b4; 0]';
    bc4 = [bc4;bc4e];
end
bc4(1,:)=[];
bc = [bc1';bc2'];
BC2 = [bc1';bc3(:,1);bc2'];
BC3 = [bc1';bc4(:,1);bc5';bc2'];
%% Solve the eigen-value problem
Ltot = l1;
[X,L]=Inverse_iteration(K,M,bc,BC2,BC3, ndof,Ltot, E,b*h^3/12,b*h,d);
```

```
    %SAVE USED PARAMETERS FOR EACH ITERATION
    EE(z,j) = E;
    L_TOT(z,j) = l1;
    MASS(z,j) = d;
    BB}(z,j)=b
    H(z,j) = h;
    %Cropping of the stiffness, mass and eigenvector matrices with
    %consideration to boundary conditions.
    K(bc, :) = [];
    K(:,bc)=[];
    M(bc, :) = [];
    M(:,bc)=[];
    X(bc,:) = [];
    %%MODE SUPERPOSITION
    %Number of eigenmodes used
    nmodes = 3;
    %% Creation of diagonal K,M and C for mode superposition
    M_diag = X(:,1:nmodes)'*M*X(:,1:nmodes);
    K_diag = X(:,1:nmodes)'*K*X(:,1:nmodes);
    %modal damping (According to BV Bro, Banverket 2006)
    if Ltot<=20
    damp = 0.015 + 0.0007*(20-Ltot);
    else
    damp = 0.015;
    end
    C_diag = diag(2*damp*2*pi*L(1:nmodes));
    %CONSIDERED RANGE OF TRAIN VELOCITY
    speed_step = 1.25/3.6;
    max_speed = 0.8*vcr;
    vcr = pi/2/Ltot^2*sqrt(E*I/(d*A))*18;
    %% Newmark integration
    bet = 0.25;
    gam = 0.5;
    %Number of time steps
    ntimes = 10000;
    %Starting displacements
    u0 = zeros(nmodes,1);
    v0 = zeros(nmodes,1);
    q = waitbar(0, ['Iteration' num2str(j)]);
    %Loop that calculate the response in time using mode
    %superposition for a chosen range of train velocities
        for i=1:(max_speed/speed_step)
            e = 3.5;
            %Train velocity
            speed = i*speed_step;
            %% Train load
            [F_train h] =
TrainLoadThreeD(Li,speed,ntimes,ndof,Ex,Ey,Edof,b,e,n2,n11,n12,j);
            F_train(bc,:) = [];
            F = X(:,1:nmodes)'*F_train;
            %TIME INTEGRATION
            [a,t] = NEWMARK(K_diag,C_diag,M_diag,F,h,u0,v0,bet,gam);
                %Calculate the true displacements and acceleration and %adds
            the removed degrees of freedom.
                e = 1:1:ndof;
                e(bc) = [];
                a = X(:,1:nmodes)*a;
                a_true = zeros(ndof,ntimes);
```

```
    a_true(e,:) = a;
    a = [];
    % MAX IN THE WHOLE BRIDGE
    % MaxMax = max(max(a_true));
    % MinMin = min(min(a_true));
    % A_maxmax(i,j) = max([MaxMax abs(MinMin)]);
    V(i,j) = speed;
    e2 = 3:5:(ndof-2);
    u_vert = u_true(e2,:);
    a_vert = a_true(e2,:);
    a_true = [];
    MaxVert = max(max(a_vert));
    MinVert = min(min(a_vert));
    A_maxvert(i,j) = max([MaxVert abs(MinVert)]);
        end
end
%% PLOT RESULTS
% A huge variation of plot options can be used. The code below is just an
% example of possible plot options
hold on
colors = ['y', 'g', 'b', 'r', 'k'];
for k=1:5
plot(V(:,k)*3.6, A_maxmax(:,k), colors(k))
end
xlabel('Speed [km/h]')
ylabel('Acceleration [m/s^2]')
title('HSLM-A1')
vari = 'E = ';
legend([vari num2str(EE(1))], [vari num2str(EE(2))],...
    [vari num2str(EE(3))], [vari num2str(EE(4))],...
    [vari num2str(EE(5))])
```


## D.3.3 Verification

ADINA and the Matlab 3D FE program has been used to build a similar model of a one span bridge with the purpose of comparing the responses from the two programs. The boundary conditions and geometry used in the verification model can been seen in Figure D-6.


Figure D-6 Geometry and boundary conditions.
The material parameters in the form of module of elasticity and concrete density have been chosen to 30 GPa and $2400 \mathrm{~kg} / \mathrm{m}^{3}$ respectively.

The time integration is made using the Newmark- $\beta$ method with the integration parameters $\beta$ and $\gamma$ equal to 0.25 and 0.5 respectively. 10000 time steps are used in the integration.

Both models use mode-superposition limited to the three eigenmodes with lowest corresponding eigenfrequency. Damping has been included through a modal damping of $2 \%$ for all considered eigenmodes.

Plate elements are used in the Matlab program to create the model above and in ADINA shell elements are used. The element size in Adina has been chosen to $0.25 \cdot 0.25 \mathrm{~m}^{2}$. In Matlab the element size were chosen to $16 \cdot 16$ elements which correspond to $0.625 \cdot 0.9375 \mathrm{~m}^{2}$. The reason why a much more coarse mesh was chosen for the Matlab program is because of the memory limit that the program has.

The result from comparing the frequencies from the two models can be seen in Table D-1.

Table D-1 Comparison between the eigenfrequencies for the two programs.

|  | Frequencies |  |
| :---: | :--- | :--- |
| Mode | ADINA | Matlab |
| 1 | 4.3 Hz | 4.3 Hz |
| 2 | 9.93 Hz | 10.03 Hz |
| 3 | 17.25 Hz | 17.26 Hz |

The further verification of the FE Matlab 3D plate program is divided in two parts. First the total response is monitored for one specific train configuration, certain bridge parameters and a velocity of $180 \mathrm{~km} / \mathrm{h}$. The response gained from the Matlab FE program is compared to the response obtained from an analysis in ADINA with the same input parameters. After the total response is verified, for one single case, max accelerations for different bridge configurations and speeds for the trainload will be compared. It is mainly the width of the bridge that will be changed in the comparison but also two different velocities of the traveling train will be tested.

The load case HSLM-A1 has been used to represent the high-speed train. It has been simulated with the velocity $180 \mathrm{~km} / \mathrm{h}$ and accelerations has been monitored in the middle of the span. Results from the comparison between total acceleration response in ADINA and Matlab are shown in Figure D-7 and Figure D-8. The response from each program shows very similar behavior and has only small deviations from one other.


Figure D-7 Comparison between displacements for the HSLM-A1 train load with velocity of 180 km/h


Figure D-8 Comparison between displacements for the HSLM-A1 train load with velocity of 180 km/h.

Since the response is nearly identical for this single case next step is to compare the maximum acceleration with some different input parameters. Results from the analysis in ADINA and Matlab are shown in Table D-2 and Table D-3 respectively. The comparison between the results in each program is shown in Table D-4.
Table D-2 Max accelerations from the different input parameters in ADINA.

| ADINA |  |  |
| :--- | :---: | :---: |
| HSLM-A1 | $\mathrm{L}=10 \mathrm{~m}, \mathrm{~b}=6 \mathrm{~m}$ | $\mathrm{~L}=10 \mathrm{~m}, \mathrm{~b}=10 \mathrm{~m}$ |
| $180.0 \mathrm{~km} / \mathrm{h}$ | $2.925 \mathrm{~m} / \mathrm{s}^{2}$ | $1.239 \mathrm{~m} / \mathrm{s}^{2}$ |
| $242.5 \mathrm{~km} / \mathrm{h}$ | $15.446 \mathrm{~m} / \mathrm{s}^{2}$ | $9.177 \mathrm{~m} / \mathrm{s}^{2}$ |

Table D-3 Max accelerations from the different input parameters in Matlab.

| Matlab |  |  |
| :--- | :---: | :---: |
| HSLM-A1 | $\mathrm{L}=10 \mathrm{~m}, \mathrm{~b}=6 \mathrm{~m}$ | $\mathrm{~L}=10 \mathrm{~m}, \mathrm{~b}=10 \mathrm{~m}$ |
| $180.0 \mathrm{~km} / \mathrm{h}$ | $2.816 \mathrm{~m} / \mathrm{s}^{2}$ | $1.262 \mathrm{~m} / \mathrm{s}^{2}$ |
| $242.5 \mathrm{~km} / \mathrm{h}$ | $15.439 \mathrm{~m} / \mathrm{s}^{2}$ | $8.897 \mathrm{~m} / \mathrm{s}^{2}$ |

Table D-4 Comparison of max accelerations from the result in ADINA and Matlab.

| ADINA vs Matlab |  |  |
| :--- | :---: | :---: |
| HSLM-A1 | $\mathrm{L}=10 \mathrm{~m}, \mathrm{~b}=6 \mathrm{~m}$ | $\mathrm{~L}=10 \mathrm{~m}, \mathrm{~b}=10 \mathrm{~m}$ |
| $180.0 \mathrm{~km} / \mathrm{h}$ | $3.7 \%$ | $1.8 \%$ |
| $242.5 \mathrm{~km} / \mathrm{h}$ | $0.0 \%$ | $3.0 \%$ |

## D. 4 Function files

The "main programs" described earlier all gather calculations from several function files. Function files are used to mesh the geometry, calculate the travelling train load, calculate the dynamic response and for solving the eigenvalue-problem. Several of these function files have been made by the authors, and these are described in the following sub-chapters.

## D.4.1 Sdof Load

The function file SdofLoad.m is used in the "Mina program" for transforming a railway bridge into a SDOF system.

The function file calculates the influence lines for the critical point of the railway bridge (in the middle of the largest span) for a two-span bridge considering static displacement when a point load moves across the bridge. The file also gathers the influence from all point loads in a HSLM-A train load into an influence line with consideration to a chosen train configuration.

## D.4.1.1 Matlab code



```
end
distance((s-5):s) = [(N+2)*D+17-d (N+2)*D+17 (N+2)*D+20.525 ...
    (N+2)*D+23.525 (N+2)*D+34.525 (N+2)*D+37.525];
start = distance./v;
h = (sum(Li)+distance(s))/v/ntimes;
%%
defl_max = Li(2)^3/(48*E*I) - Li(2)^4/(64*E*I*sum(Li))*3/4;
F_sdof = zeros(1,ntimes+1);
for i=1:ntimes
    for j=1:length(start)
        a1 = v*(h*i-start(j));
        a2 = v*(h*i-start(j))-Li(1);
        b = Li(2)-a2;
        if a1<=0
                break
            end
        if b<=0
                continue
            end
        if v*(h*i-start(j)) <= Li(1)
                F_sdof(i+1) = F_sdof(i+1) + ...
                    P*(a1*Li(1)*Li(2)^2/(32*E*I*sum(Li))*(1-a1^2/Li(1)^2));
        else if v*(h*i-start(j)) <= Li(1)+Li(2)/2
            F_sdof(i+1) = F_sdof(i+1) - ...
                    P*(a2*Li(2)^2/(48*E*I)* (3-4*a2^2/Li (2)^2)-b*Li(2)^3/...
                    (32*E*I*sum(Li))*(1-b^2/Li(2)^2));
                else
                        F_sdof(i+1) = F_sdof(i+1) - ...
                    P* (b*Li (2)^2/(48*E*I)* (3-4*b^2/Li (2)^2)-b*Li (2)^3/...
                    (32*E*I*sum(Li))*(1-b^2/Li(2)^2));
                end
        end
    end
end
F_sdof = F_sdof./defl_max;
K_bridge = 1/defl_max;
```


## D.4.2 Train load 2D

The function file is called TrainTwoD. It creates the load matrix of a HSLM-A train load. A load matrix is a matrix where every column corresponds to the load on every degree of freedom and one time step. The load matrix hence has rows equal to the degrees of freedom and columns equal to the number of time steps.
The procedure behind the calculations in the file is as follows:

- The load parameters are decided based on which HSLM-A train load that the user has chosen.
- Based on the load parameters a vector containing the distance between all point loads with the first load as reference is created. The vector is called "distance".
- Based on the "distance" vector and the chosen train velocity a vector describing the time between loads is created. The vector is called "start" as it describes the time for which every point load starts affecting the bridge.
- When the "start" vector has been established a double loop is used for performing the calculations. The first loop corresponds to every time step and the second to every point load inside each time step. For one time step the program calculates where every point load is using the "start" vector and assembles the load on the corresponding elements vertical degrees of freedom.

The assembly of the load has been made on the vertical degrees of freedom, which means that an element affected by a point load is seen as simply supported. This solution will be correct if small element lengths are used as the moments become negligible.

## D.4.2.1 Matlab code



```
end
distance((s-5):s) = [(N+2)*D+17-d (N+2)*D+17 (N+2)*D+20.525 ...
    (N+2)*D+23.525 (N+2)*D+34.525 (N+2)*D+37.525];
start = distance./v;
load = P * ones(1,s);
h = (sum(Li)+distance(s))/v/ntimes;
for i=1:ntimes
    for j=1:length(start)
        s= v*(h*i-start(j));
        if s<=0
            break
        end
        if s>=sum(Li)
            continue
        end
        for k = 1:(length(Coord)-1)
            if s <= Coord(k+1)
                el = k;
                break
            end
        end
        a = s - Coord(Elnod(el,2));
        b = Coord(Elnod(el,3)) - s;
        L = Coord(Elnod(el,3)) - Coord(Elnod(el,2));
        P1 = load(j)*b/L;
        P2 = load(j)*a/L;
        F_train(Edof(el,[3 6]),i) = [-P1 -P2]';
    end
end
```


## D.4.3 Train load 3D

The function file is called TrainThreeD. It creates the load matrix considering a threedimensional geometry. The procedure behind the calculations is very similar to that in TrainTwoD, and the reader is referred to this description for a more thorough explanation of the calculation procedure.

The additional calculation that is required when considering a 3D geometry concerns the eccentricity of the load. The only addition is that when the "start" vector has been calculated an additional vector that describes the numbering of the elements that the train will pass through needs to be established before the double loop. The "Coord" vector was used for this purpose in the 2D load creation as it was known that the load would pass through all elements.

## D.4.3.1 Matlab code



```
function [F_train h] = ...
    TrainThreeD(Li, speed, ntimes, ndof, Ex, Ey, Edof,b,e,n2,n11, n12,A)
%--------------------------------------------------------------------------------
%Pogram that calculates the load matrix for a FE analysis corresponding to
%a HSLM-A train load.
%
    INDATA
    Li = Vector containing the length of each span. The function
                        file is compatible with arbitrary number of spans.
    speed = Train velocity
    ntimes = Number of time steps
    ndof = Number of degrees of freedom
    Ex = x-coordinates for every element. Each row corresponds
            to one element. The coordinates for one element are
            displaced in the following order:
                |
    Ey = y-coordinates for every element.
    Edof = Matrix connecting dofs and elements according to Calfem
    b = Width of the bridge
    e = Eccentricity of the load
    n2 = Number of elements perpendicular to the bridge length
    coordination
    n11 = Number of elements in the first span
    n12 = Number of elements in the second span
    A = Number of the HSLM-A train load considered.
    OUTPUT
    F_train = Matrix where every column corresponds to the load for a
    time step. The size is ndof*ntimes.
F_train = zeros(ndof,ntimes);
y_coord = b/2+e;
N = [18 17 17 16 15 14 13 13 12 11 11];
D = [llllllllllllllll
d = [ 2 3.5 2 3 2 2 2 2.5 2 2];
P}=[\begin{array}{lllllllllll}{170}&{200}&{180}&{190}&{170}&{180}&{190}&{190}&{210}&{210}\end{array}]*10^3
N = N(A);
D = D(A);
d = d(A);
P = P(A);
distance = zeros(1, (N+1)*2+12);
s = length(distance);
distance(1:6) = [l0 3 14 17 20.525 20.525+d];
for i=1:(N+1)
    distance([5+2*i 6+2*i]) = [18.7625+D*i-d/2 18.7625+D*i+d/2];
end
distance((s-5):s) = [(N+2)*D+17-d (N+2)*D+17 (N+2)*D+20.525 ...
    (N+2)*D+23.525 (N+2)*D+34.525 (N+2)*D+37.525];
start = distance./speed;
```

```
load = P * ones(1,s);
h = (sum(Li)+distance(s))/speed/ntimes;
%%
for i=1:n2
    if y_coord<=Ey(i,4)
        el_row_nr = i;
            break
    end
end
for i=1:(n11+n12)
    el_row(i) = el_row_nr + n2*(i-1);
end
%%
for i=1:ntimes
    for j=1:length(start)
        s= speed*(h*i-start(j));
        if }\textrm{s}<=
            break
        end
        if s>=sum(Li)
            continue
        end
        for j = 1:length(el_row)
            if s <= Ex(el_row(j),2)
                el = el_row(j);
                break
            end
        end
        a1 = s-Ex(el,1);
        a2 = Ex (el,2)-s;
        b1 = y_coord - Ey (el,1);
        b2= Ey(el,4) - y_coord;
        P1 = P*b2*a2/((Ey(el,3)-Ey(el,1))* (Ex(el,2)-Ex(el,1)));
        P2 = P*b2*a1/((Ey (el,3)-Ey (el,1))* (Ex (el,2) - Ex (el,1)));
        P3 = P*b1*a1/((Ey (el,3)-Ey(el,1))* (Ex(el,2)-Ex(el,1)));
        P4 = P*b1*a2/((Ey (el,3)-Ey (el,1))* (Ex (el,2)-Ex(el,1)));
        F_train(Edof(el,[4 9 14 19]) ,i) = F_train(Edof(el,[[4 9 14 19])
        ,i)+[-P1 -P2 -P3 -P4]';
    end
end
```


## D.4.4 Mesh in 2D

The function file is called MeshTwoD. Its purpose is to create a two-dimensional mesh of a continuous beam. Basically it creates a row of nodes placed with a spacing that depends on the requested span lengths and number of elements.

The program can create the mesh of a continuous beam with arbitrary span lengths and number of spans. It is however limited in the choice of spacing between nodes as the same number of elements is used for all spans.

Creating the mesh of a structure basically means establishing a number of matrices that describes coordinates and structure of the mesh. The created matrices are explained in the code below.

## D.4.4.1 Matlab code



```
end
non = length(Coord);
%% Creation of Elnod, Ex, Ey
for i=1:nel
    Elnod(i,:) = [i i i+1];
end
for i=1:length(Elnod)
    ex = [Coord(Elnod(i,2),1) Coord(Elnod(i,3),1)];
    Ex(i,:) = ex;
end
Ey = zeros(nel,2);
%% Creation of Ndof, ndof
for i=1:non
    for j=1:dof
        t(j) = (dof*i-dof)+j;
    end
    Ndof(i,:) = [i t];
end
ndof = max(max(Ndof));
%% Creation of Edof
for i=1:nel
    Edof(i,:) = [i Ndof(Elnod(i,2), 2:(dof+1)) Ndof(Elnod(i,3), 2:(dof+1))];
end
%% Creation of b
b = [1 2]';
for i=1:length(Li)
    b = [b; dof*neli*i*[1 1]'+[1 2]'];
end
```


## D.4.5 Mesh in 3D

The purpose with the 3D mesh file is to create a mesh for a rectangular domain which represents a one span bridge with simply supported boundary conditions. It will therefore create a rectangular mesh with the chosen mesh density of $n 11 \mathrm{xn} 2$ and a mesh size that depend on chosen size of the bridge and mesh density. It will assemble all necessary matrices needed to use the CALFEM-toolbox, this includes the following matrices:

- Edof - Describe the connection between elements
- Ex-Element wise x-coordinates
- Ey - Element wise y-coordinates
- Pi - Row vector containing x-dof (1st value) and y-dof (2nd value) for corner node $\mathrm{i}(\mathrm{i}=1,2,3,4)$.
- Bi - Matrix containing x-dofs (1st column) and y-dofs (2nd column) for nodes on boundary segment $\mathrm{i}(\mathrm{i}=1,2,3,4)$.

In excess of these matrices the function file also calculate and provide the main program with the number of elements (nel), number of nodes (non) and number of degrees of freedom (ndof) included in the model.

## D.4.5.1 Matlab Code

```
function[B1,B2,B3,B4,P1,P2,P3,P4,Ex,Ey, ndof,nel,non,Edof]= MeshThreeD
(11,b,n11,n2)
%------------------------------------------------------------------------------
%
Bridge Mesh 3D one span
%PURPOSE:
% Generates a mesh for one rectangular domain with corners 1, 2, 3,4 % and
boundaries 1, 2, 3 and 4
%
%-----------------------------------------------------------------------------
%
    P4 B3 P3
    B4 |
Input:
xcorner - x-coordinates for the four corners of the domain (1x4)
ycorner - x-coordinates for the four corners of the domain (1x4)
xcorner - x-coordinates for the four corners of the domain (1x4)
ycorner - x-coordinates for the four corners of the domain (1x4)
elemtype - Element type "tria3" - 3-node triangular element
    "quad4" - 4-node quadrilateral element
n1 - number of elements along the boundaries 1 and 3
n2 - number of elements along the boundaries 2 and 4
Output:
Edof - Connectivity matrix for mesh
Ex - Elementwise x-coordinates, cf. Calfem Toolbox
Ey - Elementwise y-coordinates, cf. Calfem Toolbox
% Bi - Matrix containing x-dofs (1st column) and y-dofs (2nd %
column) for nodes on boundary segment i (i=1,2,3,4)
%
```

```
% Pi - Row vector containing x-dof (1st value) and y-dof (2nd
value) for corner node i (i=1,2,3,4).
%
%----------------------------------------------------------------------------
%AUTHORS: PATRIK ERIKSSON, EMANUEL TROLIN 2010-02-05
%
```



```
% Specify the geometry of the bridge with respectively
xcorner = [0 l1 l1 0];
ycorner = [0 0 b b ];
% Specify the element type to use
elemtype = 'quad4';
% The program now calculates and define all geometrical parameters
nel = (n1+1)*(n2+1);
for k=1:nel
    ndof(k,:) = [lllllll}
end
B1 = zeros(n1+1,5);
B2 = zeros(n2+1,5);
B3 = B1;
B4 = B2;
    for i=1:n1
                for j=1:n2
                    sw=(i-1).* (n2+1) +j;
                    nw=sw+1;
                    se=sw+(n2+1);
            ne=se+1;
            xiw=((i-1)/n1);
            xie=(i/n1);
            etas=((j-1)/n2);
            etan=(j/n2);
            [xsw,ysw]=bilmap (xiw, etas,xcorner,ycorner);
            [xse,yse]=bilmap (xie,etas,xcorner,ycorner);
            [xne,yne]=bilmap (xie, etan, xcorner,ycorner);
            [xnw,ynw]=bilmap (xiw, etan,xcorner,ycorner);
            elno=(i-1)*n2+(j-1);
Edof(elno+1,:)=[elno+1, ndof(sw,:), ndof(se,:),ndof(ne,:),ndof(nw,:)];
            Ex(elno+1,:)=[xsw xse xne xnw];
            Ey(elno+1,:)=[ysw yse yne ynw];
            if j == 1
                B1([i i+1],:) = [ndof(sw,:);ndof(se,:)];
            end
            if i == n1
                        B2([j j+1],:) = [ndof(se,:);ndof(ne,:)];
            end
            if j == n2
                        B3([i i+1],:) = [ndof(nw,:);ndof(ne,:)];
            end
            if i == 1
                B4([j j+1],:) = [ndof(sw,:); ndof(nw,:)];
            end
        end
    end
P2=B1 (end,:);
P1=B1 (1,:);
P3=B3 (end,:);
P4=B3 (1,:);
function [x,y]=bilmap(xi,eta, xcorner,ycorner)
```

```
N=[(1-xi).*(1-eta) xi.*(1-eta) xi.*eta (1-xi).*eta];
x=N*xcorner';
y=N*ycorner';
nel = length(Edof(:,1));
ndof = max(max(Edof));
non = ndof/5;
```


## D.4.6 Newmark

Newmark- $\beta$ method is an unconditionally stable method if the parameters are chosen properly. The function file calculates the change in acceleration, velocity and displacement for each time step by first calculating the first initial acceleration and performs a LU factorization of the mass matrix. For each time step the function file calculates updated acceleration, velocity and displacements.

The time integration function file uses the Newmark- $\beta$ method with the parameters $\beta$ and $\gamma$ equal to 0.25 and 0.5 respectively to make it unconditionally stable.

More information about the Newmark- $\beta$ method can be found in section 2.3.1.1.

## D.4.6.1 Matlab Code

```
function [u,v,a,t] = NEWMARK(K,C,M,F,h,u0,v0,bet,gam)
%-----------------------------------------------------------------------------
%
% NEWMARK INTEGRATION
O
%PURPOSE:
%Function file that calculates the response in time using the Newmark
%integration method.
%
%------------------------------------------------------------------------------
%IN DATA:
%K-Stiffness matrix
%C-Damping matrix
%M-Mass matrix
%F-Trainload
%h-Time step
%u0-displacement vector
%v0-velocity vector
%---------------------------------------------------------------------------
%AUTHORS: PATRIK ERIKSSON, EMANUEL TROLIN 2010-02-05
%
```



```
%Calculate the initial acceleration from the equations of motion
a0 = inv(M)* (F (:, 1) -C*v0-K*u0);
%Calculate the LU factorization of M
LU = M + gam*h*C + bet*h^2*K;
%Initial Acceleration, Velocity and displacement
u = zeros(size(F));
v = u;
a = u;
u(:,1) = u0;
v(:,1) = v0;
a(:, 1) = a0;
%Loop for each time step
for i=1:length(F(1,:))-1
```

```
    RHS = -K*u(:,i) - (C+h*K)*v(:,i) - ...
        (h* (1-gam)*C+h^2/2*(1-2*bet)*K)*a(:,i) + F(:,i+1);
    %Solve for the second derivatives at the next time step
    a(:,i+1) = inv(LU) *RHS;
    %Evaluate the set of displacements and velocities
    u(:,i+1) = u(:,i) + h*v(:,i) +((1-2*bet)*a(:,i) +2*bet*a(:,i+1))*h^2/2;
    v(:,i+1) = v(:,i) + ((1-gam)*a(:,i) +gam*a(:,i+1))*h;
    %Continue to the next time step:
end
t = linspace (0, length(F(1,:)) *h, length(F(1,:)));
%-------------------------------END
```


## D.4.7 Inverse iteration

The largest difference between the 2D and 3D program, besides plate elements instead of beam elements, is that the Matlab FE 3D plate model uses another technique of solving the eigen-value problem. Instead of using the Calfem file eigen.m, which itself uses the built-in Matlab function file eigen, the eigenvalues and eigenvector is obtained by the inverse iteration method. The method is effective when only a few of the lowest eigenvalues and eigenvectors are desired (Craig and Kurdila, 2006). The benefit with using only a few eigenmodes is that the time required to perform the calculation is significantly reduced which is very important when the program is repeated many times.

## D.4.7.1 Matlab Code

```
function [X1,L]=Inverse_iteration(K,M,bc,BC2, ndof)
%----------------------------------------------------------------------------
%PURPOSE:
%Calculate the first eigenvalues (L) and eigenvectors (X1) by the vector
%iteration method inverse iteration.
%-----------------------------------------------------------------------------
%INDATA:
%K-Stiffness matrix
%M-Mass matrix
%bc-boundary condition for simply supported bridge (used to calculate the
%1st eigenvalue)
%BC2-boundary condition for simply supported bridge with locked mid row.
%(used to calculate the 2nd eigenvalue)
%ndof-number of degrees of freedom
%------------------------------------------------------------------------------
%AUTHORS: PATRIK ERIKSSON, EMANUEL TROLIN 2010-03-03
%
%-------------------------------------------------------------------------------
%%------Remove the locked dofs for each boundary condition--------------
    K1=K;
    M1=M;
    K2=K;
    M2=M;
    K1 (bc,:) = [];
    K1 (:,bc)=[];
    M1 (bc,:) = [];
    M1 (:,bc)=[];
    K2 (BC2, : ) = [];
    K2 (:,BC2) = [];
```

```
M2 (BC2,:) = [];
M2 (:, BC2) = [];
%%------Calculate the first eigen-value by Inverse Iteration----------
freedof = setdiff(1:ndof,bc);
unlockeddof=ndof-length(bc);
U0=ones(unlockeddof,1);
X1=zeros(ndof,2);
%Calculate the dynamic matrix
D=inv(K1) *M1;
%Calculate the Rayleigh quotient and shape U0 with 10 iterations
U(:,1)=U0;
for i=1:10
    V(:,i)=D*U(:,i);
    Lambda(i) =V(:,i)'*K1*V(:,i)/(V(:,i)'*M1*V(:,i));
    U(:,i+1)=Lambda(i) *V(:,i);
end
%Pick the first eigenvector
X1 (freedof,1)=U(:,11);
%Calculate the first eigenfrequency
Freq1=sqrt(Lambda(i))/(2*pi);
%%------Calculate the second eigen-value by Inverse Iteration---------
freedof2 = setdiff(1:ndof,BC2);
unlockeddof2=ndof-length(BC2);
U02=ones(unlockeddof2,1);
%Calculate the dynamic matrix
D2=inv(K2) *M2;
%Calculate the Rayleigh quotient and shape U0 with 10 iterations
U2(:,1)=U02;
for i=1:10
    V2(:,i)=D2*U2(:,i);
    Lambda2(i)=V2(:,i)'*K2*V2(:,i)/(V2(:,i)'*M2*V2(:,i));
    U2 (:,i+1) =Lambda2(i)*V2(:,i);
end
%Pick the second eigenvector
X1(freedof2,2)=U2(:,11);
%Calculate the second eigenfrequency
Freq2=sqrt(Lambda2(i))/(2*pi);
% Let L be a vector of 1st and 2nd eigenfrequencies
L=[Freq1 Freq2];
%%------Normalize the eigenvectors-------------------------------------------
fdof=[1:ndof]';
[nfdof,nfdof]=size(X1);
for j=1:nfdof;
    mnorm=sqrt(X1 (:,j)'*M(fdof,fdof)*X1 (:,j));
    X1(:, j) =X1 (:, j)/mnorm;
End
----------------------------END---------------------------------
```


## E. Appendix E

## Analytical solution for the acceleration response of a SDOF model

In this document a shorter derivation of the acceleration response for a SDOF model in the form of a mass-spring system subjected to a sinusoidal load is presented. The derivation is meant to show how the acceleration is affected by the mass and the stiffness of the spring.

Consider the SDOF model as defined by Figure E-1.


Figure E-1 SDOF model in the form of a mass-spring system
The following differential equation can be conducted for the SDOF model:

$$
\begin{equation*}
M \ddot{u}+K u=p_{0} \sin (\Omega t) \tag{E.1}
\end{equation*}
$$

The solution to such a simple model can analytically be found to be:

$$
\begin{align*}
& u(t)=\frac{p_{0}}{\left(1-r^{2}\right)} * \frac{1}{K} \cos (\Omega t)  \tag{E.2}\\
& \ddot{u}(t)=a(t)=-\frac{p_{0}}{\left(1-r^{2}\right)} * \frac{\Omega^{2}}{K} \cos (\Omega t) \tag{E.3}
\end{align*}
$$

Where:

$$
\begin{align*}
& r=\frac{\Omega}{\omega}  \tag{E.4}\\
& \omega=\text { eigenfrequency of the simple SDOF model }
\end{align*}
$$

The stiffness can be expressed as:

$$
\begin{equation*}
K=M * \omega^{2} \tag{E.5}
\end{equation*}
$$

(E.3), (E.4) and (E.5) give us that:

$$
\begin{equation*}
a(t)=-\frac{p_{0}}{\left(1-r^{2}\right)} * \frac{\mathrm{r}^{2}}{M} \cos (\Omega t) \tag{E.6}
\end{equation*}
$$

Since $r$ is a constant for any combination of load frequency and eigenfrequency we get from equation (E.6) that the acceleration amplitude has an inverse proportionality to the mass and is independent of the stiffness.

## F. Appendix F

Tabled values of eigenfrequency functions and design curves
In this appendix tabled values of important results in the report are presented. The tabled values are presented to give the interest reader an opportunity of producing own figures of the results.

## F. 1 Eigenfrequency functions for two-span bridge

| My | Eigenfrequency function |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | One | Two | Three | Four | Five |
|  | 6.2832 | 9.8155 | 25.1328 | 31.8086 | 56.5488 |
| 1.06 | 6.2502 | 9.8804 | 24.879 | 32.1756 | 55.7353 |
| 1.12 | 6.164 | 10.055 | 24.2858 | 33.0796 | 54.0489 |
| 1.18 | 6.0448 | 10.3095 | 23.5717 | 34.2573 | 52.2394 |
| 1.24 | 5.9092 | 10.6195 | 22.8514 | 35.5431 | 50.5779 |
| 1.3 | 5.7681 | 10.9664 | 22.1722 | 36.8217 | 49.1763 |
| 1.36 | 5.6279 | 11.3368 | 21.5526 | 37.9718 | 48.1303 |
| 1.42 | 5.4924 | 11.7201 | 20.9992 | 38.8333 | 47.5711 |
| 1.48 | 5.3634 | 12.1069 | 20.5148 | 39.2439 | 47.6327 |
| 1.54 | 5.2415 | 12.4882 | 20.1015 | 39.1771 | 48.313 |
| 1.6 | 5.127 | 12.8543 | 19.7624 | 38.7748 | 49.4402 |
| 1.66 | 5.0197 | 13.194 | 19.5024 | 38.2008 | 50.814 |
| 1.72 | 4.9193 | 13.4954 | 19.3281 | 37.5603 | 52.2815 |
| 1.78 | 4.8254 | 13.7463 | 19.2465 | 36.9083 | 53.7173 |
| 1.84 | 4.7375 | 13.937 | 19.2628 | 36.2721 | 54.9889 |
| 1.9 | 4.6552 | 14.0629 | 19.3772 | 35.6652 | 55.9438 |
| 1.96 | 4.578 | 14.1261 | 19.5835 | 35.0942 | 56.4583 |
| 2.02 | 4.5056 | 14.1346 | 19.8697 | 34.5624 | 56.5288 |
| 2.08 | 4.4376 | 14.0996 | 20.2206 | 34.0713 | 56.2697 |
| 2.14 | 4.3736 | 14.0323 | 20.6208 | 33.6221 | 55.8145 |
| 2.2 | 4.3134 | 13.9424 | 21.0561 | 33.2157 | 55.2586 |
| 2.26 | 4.2566 | 13.8375 | 21.5138 | 32.8538 | 54.6589 |
| 2.32 | 4.2029 | 13.7233 | 21.9825 | 32.5384 | 54.0474 |
| 2.38 | 4.1522 | 13.6039 | 22.4517 | 32.2729 | 53.4419 |
| 2.44 | 4.1041 | 13.4821 | 22.9109 | 32.0615 | 52.8526 |
| 2.5 | 4.0586 | 13.36 | 23.3495 | 31.9097 | 52.2853 |
| 2.56 | 4.0154 | 13.239 | 23.7566 | 31.824 | 51.7432 |
| 2.62 | 3.9744 | 13.1199 | 24.1216 | 31.8112 | 51.2282 |
| 2.68 | 3.9354 | 13.0034 | 24.435 | 31.877 | 50.7414 |
| 2.74 | 3.8982 | 12.89 | 24.6901 | 32.0249 | 50.2835 |

## F. 2 Eigenfrequency functions for three-span bridge

| Eigenfrequency functions for my $=0.5-1.95$ |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Oy | One | Two | Three | Four | Five | Six |
|  | 14.13717 | 18.117 | 26.45459 | 56.54905 | 64.44663 | 79.0665 | 127.2388 |
|  | 14.01267 | 19.07504 | 25.49564 | 55.50075 | 67.72738 | 76.9824 | 123.7507 |
|  | 13.6693 | 20.03279 | 25.20083 | 52.91896 | 70.96009 | 77.9157 | 116.3635 |
|  | 13.18384 | 20.97804 | 25.39332 | 49.91484 | 73.895 | 80.4455 | 109.2185 |
|  | 12.63619 | 21.89429 | 25.9178 | 47.10175 | 76.10762 | 83.7038 | 103.8394 |
|  | 12.08178 | 22.75925 | 26.662 | 44.68058 | 76.97044 | 87.2813 | 100.9446 |
| 0.8 | 11.55077 | 23.54341 | 27.55216 | 42.70102 | 76.12183 | 90.9396 | 100.7712 |
| 0.85 | 11.05657 | 24.20974 | 28.54045 | 41.17168 | 74.03129 | 94.4599 | 102.6927 |
| 0.9 | 10.60331 | 24.71713 | 29.59448 | 40.09387 | 71.44769 | 97.5436 | 105.8097 |
| 0.95 | 10.1905 | 25.03 | 30.69044 | 39.46398 | 68.82738 | 99.7387 | 109.5239 |
| 1 | 9.815544 | 25.13283 | 31.80867 | 39.26248 | 66.36893 | 100.5367 | 113.4927 |
| 1.05 | 9.475083 | 25.03972 | 32.93085 | 39.44487 | 64.14845 | 99.816 | 117.4733 |
| 1.1 | 9.165607 | 24.78938 | 34.03787 | 39.94621 | 62.19286 | 98.0317 | 121.214 |
| 1.15 | 8.88377 | 24.42957 | 35.1082 | 40.69602 | 60.51199 | 95.7483 | 124.3766 |
| 1.2 | 8.62651 | 24.00339 | 36.11662 | 41.63116 | 59.11245 | 93.327 | 126.5148 |
| 1.25 | 8.391089 | 23.54356 | 37.0337 | 42.70085 | 58.00347 | 90.9494 | 127.2532 |
| 1.3 | 8.175087 | 23.07245 | 37.82687 | 43.86596 | 57.19795 | 88.6999 | 126.6255 |
| 1.35 | 7.976383 | 22.60436 | 38.46414 | 45.09604 | 56.70958 | 86.6164 | 125.0562 |
| 1.4 | 7.793117 | 22.14798 | 38.92077 | 46.36624 | 56.54631 | 84.7152 | 122.9979 |
| 1.45 | 7.623664 | 21.70828 | 39.18675 | 47.65477 | 56.70282 | 83.0036 | 120.7492 |
| 1.5 | 7.466602 | 21.28783 | 39.27072 | 48.94078 | 57.15638 | 81.4864 | 118.4746 |
| 1.55 | 7.320686 | 20.88772 | 39.19715 | 50.20285 | 57.86941 | 80.1691 | 116.2597 |
| 1.6 | 7.18482 | 20.50808 | 38.99877 | 51.41775 | 58.7968 | 79.0596 | 114.1481 |
| 1.65 | 7.058044 | 20.1485 | 38.70869 | 52.5598 | 59.89347 | 78.1697 | 112.1619 |
| 1.7 | 6.939508 | 19.80824 | 38.35566 | 53.60109 | 61.1188 | 77.5143 | 110.312 |
| 1.75 | 6.828463 | 19.48636 | 37.96239 | 54.51306 | 62.43802 | 77.1102 | 108.6039 |
| 1.8 | 6.724242 | 19.18185 | 37.54582 | 55.26978 | 63.82153 | 76.9721 | 107.0412 |
| 1.85 | 6.626255 | 18.89368 | 37.1181 | 55.85261 | 65.2435 | 77.1084 | 105.6272 |
| 1.9 | 6.533973 | 18.62083 | 36.68778 | 56.25451 | 66.68033 | 77.5162 | 104.3662 |
| 1.95 | 6.446927 | 18.36231 | 36.26074 | 56.4819 | 68.10921 | 78.1798 | 103.2646 |
|  |  |  |  |  |  |  |  |

## F. 3 Design curves for a single-span railway bridges



## F. 4 Design curves for a two-span railway bridges

| $\mathrm{L}=16 \mathrm{~m}$ |  | $\mathrm{L}=18 \mathrm{~m}$ |  | $\mathrm{L}=20 \mathrm{~m}$ |  | $\mathrm{L}=22 \mathrm{~m}$ |  | $\mathrm{L}=24 \mathrm{~m}$ |  | $\mathrm{L}=26 \mathrm{~m}$ |  | $\mathrm{L}=28 \mathrm{~m}$ |  | $\mathrm{L}=30 \mathrm{~m}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\beta$ | $\tau$ | $\beta$ | $\tau$ | P | $\tau$ | P | $\tau$ | p | $\tau$ | $\beta$ | $\tau$ | $\beta$ | $\tau$ | P | $\tau$ |
| 0.00144 | 112 | 0.00182 | 111 | 0.00224 | 106 | 0.00272 | 103 | 0.00323 | 93 | 0.00379 | 82 | 0.00439 | 85 | 0.00505 | 102 |
| 0.00575 | 374 | 0.00727 | 420 | 0.00673 | 290 | 0.00815 | 257 | 0.00968 | 300 | ${ }^{0.00758}$ | 253 | 0.00877 | 210 | 0.01010 | 215 |
| 0.01006 | 649 | 0.01273 | 610 | 0.01121 | 493 | 0.01359 | 510 | 0.01613 | 517 | 0.01136 | 358 | 0.01316 | 305 | 0.01515 | 381 |
| 0.01437 | 980 | 0.01818 | 1184 | 0.01570 | 672 | 0.01902 | 817 | 0.02258 | 1050 | 0.01894 | 587 | 0.01754 | 523 | 0.02020 | 614 |
| 0.01868 | 1522 | 0.02364 | 1769 | 0.02018 | 1269 | 0.02446 | 1123 | 0.02903 | 1114 | 0.02273 | 815 | 0.02193 | 660 | 0.02525 | 614 |
| 0.02299 | 1780 | 0.02909 | 2436 | ${ }_{0}^{0.02466}$ | 1269 | 0.02989 | 1766 | 0.03548 | 1582 | 0.02652 | 996 | 0.02632 | 839 | 0.03030 | 688 |
| 0.02730 | 2474 | 0.03455 | 2605 | 0.02915 | 1895 | ${ }^{0.03533}$ | 1766 | 0.04194 | 2001 | 0.03409 | 1338 | 0.03070 | 839 | 0.03535 | 1032 |
| 0.03161 | 2943 | 0.04000 | 3658 | 0.03363 | 1963 | ${ }^{0.04076}$ | 2273 | 0.04839 | 2083 | 0.03788 | 1654 | 0.03509 | 1149 | 0.04040 | 1032 |
| 0.03592 | 3200 | 0.04545 | 3716 | 0.03812 | 2500 | 0.04620 | 2325 | 0.05484 | 2960 | 0.04167 | 1761 | 0.03947 | 1382 | 0.04545 | 1093 |
| 0.04023 | 3780 5065 | ${ }^{0.05091}$ | 5038 | ${ }^{0.04260}$ | 3579 3579 | ${ }^{0.05163}$ | 3471 | 0.06129 | 4042 | ${ }^{0.04924}$ | 1979 | 0.04386 | 1403 | ${ }^{0.05051}$ | ${ }^{1566}$ |
| 0.04454 | 5065 | 0.05636 | 5321 | 0.04709 | 3579 | ${ }^{0.05707}$ | 3978 | 0.06774 | 4042 | ${ }^{0.05303}$ | 1979 | 0.04825 | 1469 | 0.05556 | 1720 |
| 0.04885 | ${ }^{6263}$ | 0.06182 | 6706 | 0.05157 | 5360 | ${ }^{0.06250}$ | 4688 | 0.07419 | 4042 | 0.05682 | 2808 | 0.05263 | 1965 | 0.06061 | 2306 |
| 0.05316 | 6762 | 0.06727 | 6706 | 0.05605 | 5360 | ${ }^{0.06793}$ | 4688 | 0.08065 | 4042 | ${ }^{0.06439}$ | 3714 | 0.05702 | 2303 | ${ }^{0.06566}$ | 2306 |
| 0.05747 | 6882 | 0.07273 | 6706 | 0.06054 | 5360 | 0.07337 | 4688 | 0.08710 | 5115 | 0.06818 | 3714 | 0.06140 | 2490 | 0.07071 | 2306 |
| 0.06178 | 7934 | 0.07818 | 6706 | 0.06502 | 5360 | 0.07880 | 5415 | 0.09355 | 6116 | 0.07197 | 3714 | 0.06579 | 2551 | 0.07576 | 2306 |
| 0.06609 | 7934 | 0.08364 | 7741 | 0.06951 | 5360 | 0.08424 | 6390 | 0.10000 | 6389 | 0.07955 | 3787 | 0.07018 | 2551 | 0.08081 | 2733 |
| 0.07040 | 7934 | 0.08909 | 10458 | 0.07399 | 5360 | 0.08967 | 7492 | 0.10645 | 6389 | 0.08333 | 3965 | 0.07456 | 2609 | 0.08586 | 3570 |
| 0.07471 | 7934 | 0.0945 | 12737 | 0.07848 | 5360 | ${ }^{0.09511}$ | 8295 | 0.11290 | 6389 | 0.08712 | 4632 | 0.07895 | 3167 | 0.09091 | 3570 |
| 0.07902 | 9994 | 0.10000 | 13512 | 0.08296 | 7538 | 0.10054 | 8584 | 0.11935 | 6637 | 0.09470 | 5270 | 0.08333 | 3797 | 0.09596 | 5161 |
| 0.08333 | 10541 | 0.10545 | 13512 | 0.08744 | 7538 | 0.10598 | 8584 | 0.12581 | 7577 | 0.09848 | 5454 | 0.08772 | 4068 | 0.10101 | 5161 |
| 0.08764 | 12526 | 0.11091 | 13512 | 0.09193 | 8299 | 0.11141 | 8584 | 0.13226 | 8268 | ${ }^{0.10227}$ | 5454 | 0.09211 | 4437 | 0.10606 | ${ }_{5161}$ |
| 0.09195 | 12526 | 0.11636 | 13512 | 0.09641 | 11920 | 0.11685 | 8584 | 0.13871 | 9643 | 0.10985 | 5637 | 0.09649 | 4604 | 0.11111 | 5161 |
| 0.09626 | 13074 | 0.12182 | 13512 | 0.10090 | 11920 | 0.12228 | 8584 | 0.14516 | 9704 | 0.11364 | 5637 | 0.10088 | 5648 | 0.11616 | 5161 |
| 0.10057 | 13074 | 0.12727 | 13512 | 0.10538 | 11920 | 0.12772 | 9109 | 0.15161 | 12547 | 0.11742 | 5637 | 0.10526 | 5648 | 0.12121 | 5161 |
| 0.10489 | 13074 | 0.13273 | 14149 | 0.10987 | 11920 | 0.13315 | 10500 | 0.15806 | 13387 | 0.12500 | 6276 | 0.10965 | 5648 | 0.12626 | 5161 |
| 0.10920 | 13074 | 0.13818 | 14149 | 0.11435 | 11920 | 0.13859 | 10989 | 0.16452 | 13387 | 0.12879 | 6276 | 0.11404 | 5648 | 0.13131 | 5456 |
| 0.11351 | 13074 | 0.14364 | 14149 | 0.11883 | 11920 | 0.14402 | 13478 | 0.17097 | 13387 | 0.13258 | 6380 | 0.11842 | 5648 | 0.13636 | 5630 |
| 0.11782 | 15282 | 0.14909 | 16914 | 0.12332 | 11920 | 0.14946 | 13478 | 0.17742 | 17463 | 0.14015 | 8410 | 0.12281 | 5671 | 0.14141 | 6203 |
| 0.12213 | 16589 | 0.15455 | 22502 | 0.12780 | 11920 | 0.15489 | 13478 | 0.18387 | 17763 | 0.14394 | 8410 | 0.12719 | 5671 | 0.14646 | 8404 |
| 0.12644 | 16589 | 0.16000 | 22502 | 0.13229 | 11920 | 0.16033 | 14243 | 0.19032 | 17463 | 0.14773 | 9796 | 0.13158 | 5821 | 0.15152 | 8404 |
| 0.13075 | 16589 | 0.16545 | 22502 | 0.13677 | 11920 | 0.16576 | 14593 | 0.19677 | 17463 | 0.15530 | 12803 | 0.13596 | 6887 | 0.15657 | 8404 |
| 0.13506 | 18208 | 0.17091 | 22502 | 0.14126 | 11920 | 0.17120 | 14593 | 0.20323 | 17463 | 0.15909 | 12803 | 0.14035 | 6887 | 0.16162 | 8404 |
| 0.13937 | 20272 | 0.17636 | 22846 | 0.14574 | 13787 | 0.17663 | 14593 | 0.20968 | 19029 | 0.16288 | 12803 | 0.14474 | 7630 | 0.16667 | 8404 |
| 0.14368 | 20272 | 0.18182 | 23207 | 0.15022 | 14880 | 0.18207 | 14814 | 0.21613 | 19029 | 0.17045 | 12803 | 0.14912 | 9836 | 0.17172 | 8404 |
| 0.14799 | 20272 | 0.18727 | 25302 | 0.15471 | 15127 | 0.18750 | 18373 | 0.22258 | 19029 | 0.17424 | 12803 | 0.15351 | 9836 | 0.17677 | 11344 |
| 0.15230 | 26148 | 0.19273 | 25302 | 0.15919 | 15127 | 0.19293 | 20222 | 0.22903 | 19029 | 0.17803 | 12803 | 0.15789 | 9836 | 0.18182 | 11344 |
| 0.15661 | 28214 | 0.19818 | 25302 | 0.16368 | 15129 | 0.19837 | 20222 | 0.23548 | 19029 | 0.18561 | 12803 | 0.16228 | 10971 | 0.18687 | 11344 |
| 0.16092 | 29520 | 0.20364 | 25302 | 0.16816 | 17040 | $0^{0.20380}$ | 20222 | 0.24194 | 19029 | 0.18939 | 14553 | 0.16667 | 10971 | 0.19192 | 11344 |
| ${ }^{0.16523}$ | 29520 | 0.20909 | 27019 | 0.17265 | 17957 | ${ }^{0.20924}$ | 20222 | 0.24839 | 19029 | 0.19318 | 14921 | 0.17105 | 10971 | 0.19697 | 11344 |
| 0.16954 | 29520 | 0.21455 | 29238 | 0.17713 | 22082 | 0.21467 | 20222 | 0.25484 | 19753 | 0.20076 | 14921 | 0.17544 | 10971 | 0.20202 | 11344 |
| 0.17385 | 29520 | 0.22000 | 29238 | 0.18161 | 22082 | 0.22011 | 20222 | 0.26129 | 20869 | 0.20455 | 14921 | 0.17982 | 10971 | 0.20707 | 11344 |
| 0.17816 | 29520 | 0.22545 | 29238 | 0.18610 | 22082 | 0.22554 | 20222 | 0.26774 | 20869 | 0.20833 | 14921 | 0.18421 | 10971 | 0.21212 | 11344 |
| 0.18247 | 29520 | 0.23091 | 29238 | 0.19058 | 22082 | 0.23098 | 20222 | 0.27419 | 20869 | 0.21591 | 14921 | 0.18860 | 11157 | 0.21717 | 11344 |
| 0.18678 | 29520 | 0.23636 | 29238 | 0.19507 | 22560 | 0.23641 | 20222 | 0.28065 | 20869 | 0.21970 | 14921 | 0.19298 | 11157 | 0.22222 | 11344 |
| 0.19109 | 38217 | 0.24182 | 29238 | 0.19955 | 22560 | 0.24185 | 20222 | 0.28710 | 20869 | 0.22348 | 14921 | 0.19737 | 11157 | 0.22727 | 11344 |
| 0.19540 | 46465 | 0.24727 | 29238 | 0.20404 | 22560 | 0.24728 | 20222 | 0.29355 | 22427 | ${ }_{0} 0.23106$ | 14921 | 0.20175 | 11157 | 0.23232 | 11344 |
| 0.19971 | 46465 | 0.25273 | 31314 | 0.20852 | 24532 | 0.25272 | 20222 | 0.30000 | 23532 | 0.23485 | 14921 | 0.20614 | 12311 | 0.23737 | 11344 |
| 0.20402 | 46465 | 0.25818 | 35808 | 0.21300 | 24541 | 0.25815 | 21253 | 0.30645 | 23532 | 0.23864 | 14921 | 0.21053 | 15078 | 0.24242 | 11344 |
| 0.20833 | 46465 | 0.26364 | 37156 | 0.21749 | 24541 | 0.26359 | 29883 | 0.31290 | 23532 | 0.24621 | 15646 | 0.21491 | 15078 | 0.24747 | 11344 |
| 0.21264 | 46465 | 0.26909 | 38099 | 0.22197 | 24541 | 0.26902 | 32790 | 0.31935 | 23532 | 0.25000 | 18062 | 0.21930 | 15078 | 0.25253 | 11344 |
| 0.21695 | 46465 | 0.27455 | 38099 | 0.22646 | 24541 | 0.27446 | 34874 | 0.32581 | 23532 | 0.25379 | 18062 | 0.22368 | 15078 | 0.25758 | 12598 |
| 0.22126 | 46465 | 0.28000 | 38099 | 0.23094 | 24541 | 0.27989 | 34874 | ${ }^{0.33226}$ | 23532 | 0.26136 | 19243 | 0.22807 | 15078 | ${ }^{0.26263}$ | 13877 |
| 0.22557 | 46465 | 0.28545 | 38099 | 0.23543 | 24541 | 0.28533 | 34874 | 0.33871 | 23532 | 0.26515 | 22268 | 0.23246 | 15078 | 0.26768 | 13877 |
| 0.22989 | 46465 | 0.29091 | 38099 | 0.23991 | 28061 | ${ }^{0.29076}$ | 34874 | 0.34516 | 23532 | 0.26894 | 22564 | 0.23684 | 15078 | 0.27273 | 13877 |
| 0.23420 | 46465 | ${ }^{0.29636}$ | 38099 | 0.24439 | 30678 | ${ }^{0.29620}$ | 34874 | 0.35161 | 23532 | 0.27752 | 23310 | 0.24123 | 15078 | 0.27778 | 13877 |
| 0.23851 | 46465 | 0.30182 | 38844 | 0.24888 | 31478 | 0.30163 | 34874 | 0.35806 | 23532 | 0.28030 | 23310 | 0.24561 | 15078 | 0.28283 | 13877 |
| 0.24282 | 46465 | 0.30727 | 41011 | 0.25336 | 31736 | 0.30707 | 34874 | 0.36452 | 27268 | 0.28409 | 23310 | 0.25000 | 15078 | 0.28788 | 13877 |
| 0.24713 | 46465 | 0.31273 | 42745 | 0.25785 | 31736 | ${ }^{0.31250}$ | 34874 | 0.37097 | 32094 | 0.29167 | 23310 | 0.25439 | 15078 | 0.29293 | 13877 |
| 0.25144 | 46465 | 0.31818 | 42745 | 0.26233 | 31736 | ${ }_{0}^{0.31793}$ | 34874 | 0.37742 | 32094 | 0.29545 | 23310 | 0.25877 | 15078 | 0.29798 | 13877 |
| 0.25575 | 46465 | 0.32364 | 42745 | 0.26682 | 31736 | 0.32337 | 34874 | 0.38387 | 32094 | 0.29924 | 23310 | 0.26316 | 15078 | 0.30303 | 13877 |
| 0.26006 | 47095 | 0.32909 | 42745 | 0.27130 | 31736 | ${ }^{0.32880}$ | 34874 | 0.39032 | 32094 | 0.30682 | 23310 | 0.26754 | 15778 | 0.30808 | 13877 |
| 0.26437 | 52510 | 0.33455 | 42745 | 0.27578 | 31736 | ${ }^{0.33424}$ | 34874 | 0.39677 | 32094 | 0.31061 | 23310 | 0.27193 | 19054 | 0.31313 | 13877 |
| 0.26868 | 54718 | 0.34000 | 42745 | 0.28027 | 31736 | 0.33967 | 34874 | 0.40323 | 32094 | 0.31439 | 23310 | 0.27632 | 19054 | 0.31818 | 14975 |
| 0.27299 | 57812 | 0.34545 | 42745 | 0.28475 | 31736 | 0.34511 | 34874 | 0.40968 | 32094 | 0.32197 | 23310 | 0.28070 | 19054 | 0.32323 | 17290 |
| 0.27730 | 57812 | 0.35091 | 42745 | 0.28924 | 31736 | 0.35054 | 34874 | 0.41613 | 32094 | ${ }^{0.32576}$ | 23310 | 0.28509 | 19054 | 0.32828 | 19995 |
| 0.28161 | 57812 | ${ }^{0.35636}$ | 42745 | 0.29372 | 32409 | ${ }^{0.35598}$ | 34874 | ${ }^{0.42258}$ | 32094 | ${ }^{0.32955}$ | 24716 | 0.28947 | 19054 | ${ }^{0.33333}$ | 19995 |
| 0.28592 | 57812 | 0.36182 | 42745 | 0.29821 | 32409 | 0.36141 | 34874 | 0.42903 | 32094 | 0.33712 | 24716 | 0.29386 | 19054 | 0.33838 | 19995 |
| 0.29023 | 57812 | 0.36727 | 45019 | 0.30269 | 32409 | 0.36685 | 34874 | 0.43548 | 32094 | 0.34091 | 24716 | 0.29825 | 19054 | 0.34343 | 19995 |
| 0.29454 | 57812 | 0.37273 | 48486 | 0.30717 | 32409 | 0.37228 | 34874 | 0.44194 | 32094 | 0.34470 | 24716 | 0.30263 | 19054 | 0.34848 | 19995 |
| 0.29885 | 57812 | 0.37818 | 54905 | 0.31166 | 32409 | ${ }^{0.37772}$ | 34874 | 0.44839 | 32094 | 0.35227 | 24716 | 0.30702 | 19054 | 0.35354 | 19995 |
| ${ }^{0.30316}$ | 57812 | ${ }^{0.38364}$ | 60766 | 0.31614 | 32409 | ${ }^{0.38315}$ | 34874 | 0.45484 | 32094 | ${ }^{0.35606}$ | 24716 | 0.3140 | 19054 | ${ }^{0.35859}$ | 19995 |
| 0.30747 | 57812 | 0.38809 | 63645 | 0.32063 | 32409 | 0.38859 | 34874 | 0.46129 | 32094 | 0.35985 | 24716 | 0.31579 | 19054 | 0.36364 | 19995 |
| 0.31178 | 57812 | 0.39455 | 64262 | 0.32511 | 32409 | 0.39402 | 34874 | 0.46774 | 33017 | 0.36742 | 24716 | 0.32018 | 19054 | 0.36869 | 20332 |
| ${ }_{0}^{0.31609}$ | 57812 | ${ }^{0.40000}$ | 64457 | ${ }^{0.32960}$ | 32409 | ${ }^{0.39946}$ | 34874 | ${ }^{0.47419}$ | 39313 | ${ }^{0.37121}$ | 24716 | ${ }^{0.32456}$ | 19054 | ${ }_{0}^{0.37374}$ | ${ }_{20593}^{20593}$ |
| 0.32040 | 57812 | ${ }^{0.40545}$ | 69046 | 0.33408 | 32409 | ${ }^{0.40489}$ | 34874 | 0.48065 | 48241 | ${ }^{0.37500}$ | 24716 | ${ }^{0.32895}$ | 19054 | 0.37879 | 20593 |
| 0.32471 | 57812 | 0.41091 | 69956 | 0.33857 | 32409 | 0.41033 | 34874 | 0.48710 | 59063 | 0.38258 | 24716 | 0.33333 | 19054 | 0.38384 | 20593 |
| 0.32902 | 62602 | 0.41636 | 69956 | 0.34305 | 32409 | 0.41576 | 34874 | 0.49355 | 71913 | ${ }^{0.38636}$ | 24716 | 0.33772 | 19054 | 0.38889 | 20593 |
| 0.33333 | 63900 | 0.42182 | 72354 | 0.34753 | 32409 | ${ }^{0.42120}$ | 34874 | 0.50000 | 78358 | ${ }^{0.39015}$ | 24716 | 0.34211 | 19054 | 0.39394 | 20593 |
| 0.33764 | 63900 | 0.42727 | 75925 | ${ }^{0.35202}$ | 32409 | ${ }^{0.42663}$ | 34874 |  |  | ${ }^{0.39773}$ | 24716 | ${ }^{0.34649}$ | 19054 | ${ }_{0}^{0.398999}$ | ${ }_{20593}$ |
| 0.34195 | 63900 | ${ }^{0.43273}$ | 75925 | ${ }^{0.35650}$ | 32409 | ${ }^{0.43207}$ | 34874 |  |  | ${ }^{0.40152}$ | 28022 | 0.35088 | 19054 | 0.40404 | 20593 |
| 0.34626 | 63900 | 0.43818 | 75925 | 0.36099 | 32409 | 0.43750 | 35815 |  |  | 0.40530 | 32505 | 0.35526 | 19054 | 0.40909 | 20593 |
| ${ }^{0.35057}$ | ${ }^{63900}$ | ${ }^{0.44364}$ | 75925 | ${ }_{0}^{0.36547}$ | 32409 | ${ }^{0.44293}$ | 43037 |  |  | ${ }^{0.412888}$ | ${ }^{34721}$ | ${ }_{0}^{0.35965}$ | 19054 | ${ }^{0.41414}$ | ${ }_{20593}^{20593}$ |
| ${ }^{0.35459}$ | ${ }^{63900}$ | ${ }^{0.44909}$ | 75925 | ${ }_{0}^{0.36996}$ | 32409 | ${ }^{0.44837}$ | 44221 |  |  | ${ }^{0.41667}$ | 34721 | 0.36404 | 19054 | 0.41919 | ${ }_{2}^{20593}$ |
| ${ }^{0.35920}$ | ${ }_{6}^{63900}$ | ${ }^{0.44455}$ | 75925 75925 | ${ }_{0}^{0.37474}$ | 32409 3209 | 0.45380 0.45924 | ${ }_{5}^{47305}$ |  |  | 0.42045 0.2083 | ${ }_{3}^{34721}$ | ${ }^{0.36842}$ | 19054 | ${ }^{0.42424}$ | ${ }_{20593}$ |
| 0.36351 | ${ }^{63900}$ | 0.46000 | 75925 | 0.37892 | 32409 | ${ }^{0.45924}$ | 54875 |  |  | ${ }^{0.42803}$ | 36668 |  | 19054 | ${ }^{0.42929}$ | ${ }_{2}^{20593}$ |
| 0.36782 0.37213 | 70900 74013 | 0.46545 0.47091 | 75925 75925 | 0.38341 0.38789 | 32409 32409 | 0.46467 0.47011 | 60273 64108 |  |  | 0.43182 0.43561 | 36668 <br> 36688 | 0.37719 0.38158 |  | 0.43434 | 20593 24021 |
| 0.37213 0.37644 | 74013 74013 | 0.47091 0.47636 | 75925 75925 | 0.38789 0.3938 | 32409 33287 | 0.47011 0.47554 | 64108 73180 |  |  | 0.43561 0.44318 | 36668 36688 | 0.38158 0.3896 | 19054 19054 | 0.43939 0.4444 | 24021 26728 |
| 0.38075 | 74013 | 0.48182 | 75925 | 0.39686 | 36496 | 0.48098 | 79552 |  |  | 0.44697 | 36668 | 0.39035 | 19054 | 0.44949 | 26728 |
| 0.38506 | 74013 | 0.48727 | 75925 | 0.40135 | 39282 | 0.48641 | 79552 |  |  | 0.45076 | 36668 | 0.39474 | 19054 | 0.45455 | 26728 |
| 0.38937 | 74013 | 0.49273 | 75925 | 0.40583 | 41999 | 0.49185 | 82370 |  |  | 0.45833 | 36668 | 0.39912 | 19217 | 0.45960 | 27164 |
| 0.39368 0.39799 | 74013 74013 | 0.49818 0.50000 | 75925 75925 | 0.41031 0.41480 | 43985 47018 | 0.49728 0.50000 | 96649 100133 |  |  | 0.46212 0.46591 | 36668 <br> 36688 | 0.40351 0.40789 | 20443 | ${ }^{0.46465}$ | ${ }_{3}^{29258}$ |
| 0.39799 0.4023 | 74013 74013 | 0.50000 | 75925 | 0.41480 0.41928 | 47018 53735 | 0.50000 | 100133 |  |  | 0.46591 0.47348 | 36668 36688 | 0.40789 0.41228 | 21715 24126 | 0.46970 0.47775 | 31445 36286 |
| 0.40661 | 74013 |  |  | 0.42377 | 60390 |  |  |  |  | 0.47727 | 36668 | 0.41667 | 24126 | 0.47980 | 43584 |
| 0.41092 | 74013 |  |  | 0.42825 | 60991 |  |  |  |  | 0.48106 | 36668 | 0.42105 | 28584 | 0.48485 | 43656 |
| 0.41523 0.41954 | 74013 |  |  | 0.43274 0.43722 | ${ }_{6}^{60991}$ |  |  |  |  | 0.48864 0.49242 | 36668 <br> 36688 | 0.42544 0.42982 | 31699 31699 | 0.48990 0.49495 | 43656 46564 |
| 0.41954 0.42885 | 74013 74013 |  |  | 0.43722 0.44170 | 67721 74991 |  |  |  |  | 0.49242 0.49621 | 36668 <br> 36688 | 0.42982 0.43421 | 31699 31699 | 0.49495 0.5000 | 46564 49599 |
| 0.42816 | 74013 |  |  | 0.44619 | 75777 |  |  |  |  | 0.50000 | 36668 | 0.43860 | 34510 |  |  |
| 0.43247 | 74013 |  |  | 0.45067 | 75777 |  |  |  |  |  |  | 0.44298 | 39887 |  |  |
| 0.43678 0.44109 | 74013 74013 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 0.44109 0.4540 | 74013 74013 |  |  | 0.45964 0.46413 | 75777 75777 |  |  |  |  |  |  | 0.45175 0.45614 | 39887 <br> 39887 |  |  |
| 0.44971 | 74013 |  |  | 0.46861 | 75777 |  |  |  |  |  |  | 0.46053 | 39887 |  |  |
| 0.44402 0.45833 | 74013 74013 |  |  | 0.47309 0.47758 | 78739 85157 |  |  |  |  |  |  | 0.44491 0.46930 | 39887 39887 |  |  |
| $\begin{aligned} & 0.45833 \\ & 0.46264 \end{aligned}$ | $\begin{aligned} & 74013 \\ & 74013 \end{aligned}$ |  |  | 0.47758 0.48206 | $\begin{aligned} & 85157 \\ & 86462 \end{aligned}$ |  |  |  |  |  |  | 0.46930 0.47368 | $\begin{aligned} & 39887 \\ & 39887 \end{aligned}$ |  |  |
| 0.46264 0.4695 | 74013 76391 |  |  | 0.48206 0.4855 | ${ }_{86462}^{86462}$ |  |  |  |  |  |  | 0.47368 0.47807 | 39887 39887 |  |  |
| 0.47126 | 82072 |  |  | 0.49103 | 86462 |  |  |  |  |  |  | 0.48246 | 39887 |  |  |
| 0.47557 0.47989 | 87548 93846 |  |  | 0.49327 0.49552 | 86462 86462 |  |  |  |  |  |  | 0.48684 0.49123 | 39887 39887 |  |  |
| 0.47989 0.48420 | 93846 98911 |  |  |  | 86462 89424 |  |  |  |  |  |  | $\begin{aligned} & 0.49123 \\ & 0.49561 \end{aligned}$ | $\begin{aligned} & 39887 \\ & 39887 \end{aligned}$ |  |  |
| ${ }_{0}^{0.48851}$ | 106099 |  |  |  |  |  |  |  |  |  |  |  | 39887 |  |  |
| ${ }_{0}^{0.49282}$ | 116435 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 0.49713 0.50000 | $\begin{aligned} & 128341 \\ & 134908 \end{aligned}$ |  |  |  |  |  |  |  |  |  |  |  |  |  |  |

F. 5 Design curves for a three-span railway bridge

| $\mathrm{L}=30 \mathrm{~m}$ |  | $\mathrm{L}=31.5 \mathrm{~m}$ |  | $\mathrm{L}=33 \mathrm{~m}$ |  | $\mathrm{L}=34.5 \mathrm{~m}$ |  | $\mathrm{L}=36 \mathrm{~m}$ |  | $\mathrm{L}=37.5 \mathrm{~m}$ |  | $\mathrm{L}=39 \mathrm{~m}$ |  | $\mathrm{L}=40.5 \mathrm{~m}$ |  | $\mathrm{L}=42 \mathrm{~m}$ |  | $\mathrm{L}=43.5 \mathrm{~m}$ |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\beta$ | $\tau$ | $\beta$ | $\tau$ | $\beta$ | $\tau$ | $\beta$ |  | $\beta$ | $\tau$ | $\beta$ |  | $\beta$ | $\tau$ | $\beta$ | $\tau$ | $\beta$ | г | $\beta$ | $\tau$ |
| 0.00242 | 116.2585 | 0.002472 | 102.9108 | 0.002713 | 96.62865 | 0.002965 | 97.47153 | 0.00329 | 108.8065 | 0.003503 | 118.6441 | 0.003789 | 106.3135 | 0.004886 | 95,48339 | 0.004395 | 108.4476 | 0.004714 | ${ }^{89} 98916$ |
| 0.008968 | 523.5175 | 0.009888 | 446.4143 | 0.010852 | 431.9652 | 0.01861 | 445.38 | 0.00986 | 275.7452 | 0.01051 | 355.9948 |  | 354.9664 | 0.012259 |  | 0.013184 | 22835 | 0.14142 | 323.7149 |
| 0.015695 | 654.173 | 0.017303 | 675.5032 | 0.018991 | 617.6315 | 0.020756 | 768.5169 | 0.016143 | 507.3154 | 0.017516 | 636.0863 | 0.018946 | 57.4226 | 0.020431 | 570.9336 | 0.021973 | 705.8436 | 0.02357 | 540.0452 |
| 0.022421 | 1222.764 | 0.024719 | 1083.903 | 0.027129 | 1175.115 | 0.029652 | 1102.548 | 0.0226 | 81.9627 | 0.024523 | 87.7022 | 0.026524 | 1030.499 | 0.02860 | ${ }_{921.3637}$ | 0.03762 | 798.656 | 0.03298 | 975.885 |
| 0.02914 | 1657.914 | 0.032135 | 1511.036 | 0.035268 | 1801.994 | 0.038547 | 1772.889 | 0.029058 | 927.6 | 0.031529 | 933.5443 | 0.034102 | 1364.993 | 0.036776 | 1223.927 | 0.039551 | 1400.427 | 0.042426 | 1507.089 |
| 0.035874 | 2291.726 | 0.039551 | 2493.666 | 0.043407 | 1868.109 | 0.04743 | 2170.185 | 0.03515 |  | 0.038536 | 1427.013 | 0.041681 | 1662.552 | 0.04948 | 1540.826 | 0.04834 | 1506.934 | 0.051854 | 1943.059 |
| ${ }^{0.0426}$ | 2791.472 | 0.046966 | 2493.666 | 0.051546 | 2612.694 | 0.056338 | 2447.445 | 0.041972 | 1715.338 | 0.045543 | 1913.929 | 0.049259 | 2456.931 | 0.053121 | 1936.714 | 0.057129 | 1816.966 | 0.061282 | 2240.409 |
| 0.099326 | 3428.53 | 0.054382 | 3259.332 | 0.059685 | 3161.418 | 0.065234 | 2963.162 | 0.048429 | 2004.547 | 0.052549 | 2264.75 | 0.056837 | 2456.931 | 0.061293 | 2346.848 | 0.065918 | 2188.127 | 0.07071 | ${ }^{2456.256}$ |
| 0.056052 | 3519.524 | 0.06178 | 3814.478 | 0.067823 | 3530.805 | 0.074129 | 3732.116 | 0.054887 | 2601.103 | 0.059556 | 2433.026 | 0.064415 | 2456.931 | 0.06946 | 2746.571 | 0.074707 | 2508.274 | 0.080138 | 2468.999 |
| 0.06279 | 4752.455 | 0.069214 | 3835.656 | 0.075962 | 4338.856 | 0.08322 | 4051.167 | 0.061344 | 2828.286 | 0.066562 | 2744.046 | 0.071994 | 2886.47 | 0.077638 | 2978.072 | 0.083496 | 2770.46 | 0.089566 | 3443.74 |
| 0.069505 | 4752.455 | 0.07629 | 4666.247 | 0.084101 | 4338.856 | 0.09192 | 4743.905 | 0.067801 | 2888.286 | 0.073569 | 2841.485 | 0.079572 | 3037.914 | 0.085811 | 3115.145 | 0.09285 | 3403.181 | 0.08899 | 3443.74 |
| 0.076231 | 5412.912 | 0.084045 | 5644.819 | 0.09224 | 5061.093 | 0.100816 | ${ }^{5808.86}$ | 0.074258 | 3762.7 | 0.080575 | 3167.63 | 0.08715 | 3400.239 | 0.093983 | 3927.128 | 0.101074 | 4127.702 | 0.108422 | ${ }^{3804.579}$ |
| 0.082958 | 5412.912 | 0.091461 | 5644.819 | 0.100379 | 5929.631 | 0.109712 | 5808.86 | 0.080716 | 3762.7 | 0.087582 | 4217.125 | 0.094729 | 3665.302 | 0.102156 | 3927.128 | 0.109863 | 5132.154 | 0.11785 | 4722.025 |
| 0.089884 | 6227.773 | 0.09887 | 7431.03 | 0.108518 | 6142.925 | 0.118607 | 6431.733 | 0.087173 | 4088.297 | 0.094588 | ${ }^{4303.01}$ | 0.102307 | 4825.444 | 0.110328 | 4624.145 | 0.118652 | 5132.154 | 0.127278 | ${ }^{6148.335}$ |
| 0.09641 | 6707.287 | 0.106292 | 7774.616 | 0.116656 | 8455.796 | 0.127503 | 8493.045 | 0.09363 | 4510.725 | 0.101595 | 4812.522 | 0.109885 | 4825.444 | 0.1185 | 5550.994 | 0.127741 | 5354.236 | 0.136706 | ${ }^{6148.335}$ |
| 0.103136 | 7543.619 | 0.113708 | 7774.616 | 0.124795 | 845.796 | 0.136398 | 9187.112 | 0.10087 | 4988.699 | 0.108602 | 5229.263 | 0.117463 | 5835.043 | 0.126673 | 6922.597 | ${ }^{0.13623}$ | 5941.833 | 0.146134 | ${ }^{6148.335}$ |
| 0.109863 | 8976.76 | 0.121124 | 9734.305 | 0.132934 | 8455.796 | 0.145294 | 9187.112 | 0.106445 | 5219.374 | 0.115608 | 5986.295 | 0.125042 | 6499.988 | 0.138445 | 6922.597 | 0.145019 | 6790.933 | 0.15562 | 6914.119 |
| 0.116589 | 10746.86 | 0.12854 | 10232.45 | 0.141073 | 10771.72 | 0.154189 | 9187.112 | 0.113002 | 5959.318 | 0.122615 | 7557.29 | 0.13262 | 6499.988 | 0.143018 | 6922.597 | 0.153808 | 7723.014 | 0.16499 | ${ }^{8364.992}$ |
| 0.123315 | 10923.15 | 0.135955 | 10232.45 | 0.149212 | 10771.72 | 0.163385 | 12372.66 | 0.111959 | ${ }^{6849.97}$ | 0.129621 | 7557.29 | 0.140198 | 7567.581 | ${ }^{0.15119}$ | ${ }^{7963.4588}$ | 0.162597 | 10904.21 | 0.174418 | ${ }^{8573.748}$ |
| 0.130042 | 10933.15 | 0.143371 | 10232.45 | 0.15735 | 10971.57 | 0.17198 | 12372.66 | 0.125916 | 7234.196 | 0.136628 | 7557.29 | 0.147777 | 7857.357 | 0.159363 | 8936.625 | 0.1711386 | 11308.06 | 0.183846 | ${ }^{8573.748}$ |
| 0.136768 | 12948.06 | 0.150787 | 11249.89 | 0.165489 | 13062.57 | 0.188876 | 14612.01 | 0.132374 | 7513.355 | 0.143634 | 7557.29 | 0.155355 | ${ }^{92393.27}$ | 0.167535 | 8936.625 | 0.180175 | 11308.06 | 0.19374 | ${ }^{8803.226}$ |
| 0.143494 | 12948.06 | 0.158823 | 13453.88 | 0.173628 | 13062.57 | 0.18971 | 15229.77 | 0.138831 | 7935.919 | 0.150641 | 837.664 | 0.162933 | 10273.79 | 0.175708 | 9677.676 | 0.188964 | 11987.79 | 0.202702 | ${ }^{11215.63}$ |
| 0.15022 | 13672.11 | 0.165618 | 13453.88 | 0.181767 | 13979.36 | 0.198667 | 16133.51 | 0.145288 | 8297.098 | 0.157647 | 10132.89 | 0.170511 | 11170.66 | 0.18388 | 9677.676 | 0.197753 | 13306.52 | 0.21213 | ${ }^{11605.98}$ |
| 0.15697 | 13672.11 | 0.173034 | 13649.88 | 0.189906 | 13979.36 | 0.20756 | 16770.51 | 0.151745 | 9127.746 | 0.164654 | 10132.89 | 0.17809 | 11170.66 | 0.192052 | 10800.42 | 0.206542 | 13306.52 | 0.221558 | 11605.98 |
| 0.163673 | 15637.24 | 0.18845 | 16129.84 | 0.198845 | 14525.99 | 0.216458 | 18784.41 | 0.158203 | 10697.7 | ${ }^{0.17166}$ | 1024032 | 0.185668 | 11170.66 | 0.200225 | 12557.53 | 0.215331 | 14475.32 | 0.230987 | ${ }^{11605.98}$ |
| 0.17039 | 16452.53 | 0.187865 | 16588.22 | 0.206183 | 18237.1 | 0.225353 | 18784.41 | 0.16466 | 11577.87 | 0.178667 | 13558.28 | 0.193246 | 14287.51 | 0.208397 | 12557.53 | 0.22412 | 14475.32 | 0.240415 | ${ }^{12481.29}$ |
| 0.177126 | 16452.53 | 0.195281 | 20353.94 | 0.214322 | 18237.1 | 0.234249 | 18784.41 | 0.171117 | 11577.87 | 0.185674 | ${ }^{131558.28}$ | 0.200825 | 14288751 | ${ }^{0.221657}$ | 12579.15 | 0.232909 | ${ }^{14475.32}$ | ${ }^{0.2498843}$ | ${ }^{14666.03}$ |
| 0.183852 | 16452.53 | 0.202697 | 20353.94 | 0.222461 | 19092.43 | 0.243145 | 18784.41 | 0.177574 | 11577.87 | ${ }^{0.19268}$ | 13158.28 | 0.208403 | 14287.51 | 0.224742 | 12579.15 | 0.241698 | 14475.32 | 0.25971 | ${ }^{14666.03}$ |
| 0.190578 | 16668.3 | 0.21011 | 20353.94 | 0.2306 | 20877.56 | 0.25204 | 18784.41 | 0.184031 | 11577.87 | 0.199687 | 13318.13 | 0.215981 | ${ }^{146848.6}$ | 0.232915 | 12579.15 | 0.250487 | 14475.32 | 0.26869 | ${ }^{22362.31}$ |
| 0.197304 | 16668.3 | 0.217528 | 20353.94 | 0.238739 | 20877.56 | 0.260936 | 18784.41 | 0.19048 | 11577.87 | 0.206693 | 16362.41 | 0.223559 | 17192.99 | 0.241087 | 13987.01 | 0.259276 | 14475.32 | 0.278127 | ${ }^{22362.31}$ |
| 0.204031 | 19545.53 | 0.224944 | 20353.94 | 0.246877 | 24229.67 | 0.269831 | 19400.01 | 0.19646 | ${ }^{11731.29}$ | ${ }^{0.2137}$ | 16362.41 | 0.231138 | 17602.16 | ${ }^{0.24926}$ | 15844.07 | 0.268865 | 14594.87 | 0.287555 | ${ }^{223623.31}$ |
| 0.210757 | 24180.35 | 0.23236 | 20353.94 | 0.255016 | 24299.67 | 0.278727 | 19400.01 | 0.203403 | ${ }^{14617.36}$ | 0.220706 | 16362.41 | 0.238716 | 17602.16 | 0.257732 | 15844.07 | 0.276854 | 14594.87 | 0.296983 | ${ }^{22362.31}$ |
| 0.217783 | 25446.55 | 0.239776 | 20353.94 | 0.263155 | 27536.41 | 0.287622 | 19400.01 | 0.20986 | 14617.36 | 0.227713 | 16362.41 | 0.246294 | 17602.16 | 0.265604 | 15844.07 | 0.285643 | 14594.87 | 0.306411 | ${ }^{223623231}$ |
| 0.22421 | 27474.51 | 0.247191 | 20353.94 | 0.271129 | 33125.83 | 0.296518 | 19400.01 | 0.216318 | 14617.36 | 0.234719 | 1632.41 | 0.253873 | 17602.16 | 0.273777 | ${ }^{18231.83}$ | 0.29433 | 14594.87 | 0.315839 | ${ }^{22362.31}$ |
| 0.230936 | 27754.45 | 0.254607 | ${ }^{21620.76}$ | 0.279433 | 33125.83 | 0.305413 | 19400.01 | 0.222775 | 17421.04 | 0.241726 | 16362.41 | 0.261451 | ${ }^{17602.16}$ | 0.281949 | ${ }^{18231.83}$ | 0.303221 | 14594.87 | ${ }^{0.3325267}$ | ${ }^{22362231}$ |
| 0.237662 | 27754.45 | 0.262023 | 26515.68 | 0.287571 | 33125.83 | 0.314339 | 19400.01 | 0.22932 | 18645.28 | 0.288733 | 16362.41 | 0.269029 | 22042.27 | 0.290122 | 18231.83 | ${ }^{0.31201}$ | 14594.87 | 0.334695 | ${ }^{223623231}$ |
| 0.244388 | 27754.45 | 0.269439 | 28155.82 | 0.29571 | 33125.83 | 0.323204 | 19490.01 | 0.235689 | 18645.28 | 0.255739 | 16362.41 | 0.276607 | 22042.27 | 0.298294 | 18231.83 | 0.322799 | 14594.87 | 0.344123 | ${ }^{223623.31}$ |
| 0.25115 | 27754.45 | 0.278854 | 28155.82 | 0.303849 | 33125.83 | ${ }^{0.3321}$ | 22609.67 | 0.242147 | 18645.28 | 0.262746 | 17897.95 | 0.284186 | 22042.27 | 0.306467 | ${ }^{18231.83}$ | 0.329588 | 14594.87 | 0.353551 | ${ }^{223623231}$ |
| 0.257841 | 27754.45 | 0.28427 | 28155.82 | 0.311988 | 33125.83 | 0.340999 | ${ }^{23332.09}$ | 0.248804 | 18645.28 | 0.269752 | 199488.88 | 0.291764 | 22042.27 | 0.314639 | ${ }^{18231.83}$ | 0.338377 | 14594.87 | 0.362979 | ${ }^{223623.31}$ |
| 0.264567 | 27754.45 | 0.291686 | 28155.82 | 0.320127 | 33125.83 | 0.3448971 | 23332.09 | 0.255061 | 18645.28 | 0.276759 | 19948888 | 0.299342 | 22042.27 | 0.322812 | ${ }^{18231.83}$ | 0.347166 | 17397.4 | 0.372407 | ${ }^{22362231}$ |
| 0.271294 | 27754.45 | 0.299102 | 28155.82 | 0.328266 | 33125.83 | 0.358786 | 23332.09 | 0.261518 | 19602.21 | 0.283765 | 19488.88 | 0.30692 | 22042.27 | 0.330984 | 20875.27 | 0.355955 | 17729.59 | 0.381835 | ${ }^{22362.31}$ |
| 0.27802 | 27754.45 | 0.306517 | 2815.82 | 0.336404 | 33125.83 | 0.367882 | 28125.47 | 0.267976 | 23523.44 | 0.290772 | 19488.88 | 0.314499 | 22042.27 | 0.339156 | 23303.24 | 0.36474 | 17729.59 | 0.391263 | ${ }^{223623.31}$ |
| ${ }^{0.288746}$ | 27754.45 | 0.313933 | 28155.82 <br> 28558 <br> 1858 | 0.034543 | 33125.83 <br> 3312853 | ${ }^{0.376577}$ | 30958.59 <br> 310856 | ${ }^{0.2774333}$ | ${ }_{2} 2514.444$ | 0.397778 | 19488.88 | 0.322077 | ${ }^{22042} 2$ | 0.347329 | ${ }^{233033.24}$ | ${ }^{0.3573533}$ | 1772.59 | ${ }^{0.400691}$ | ${ }^{22362.31}$ |
| 0.291473 | 27754.45 | 0.321349 | 28155.82 | 0.352682 | 33125.83 | 0.385473 | ${ }^{31085.66}$ | ${ }^{0.28889}$ | 25143.44 | 0.304785 | 194888.88 | 0.329655 | 22042.27 | 0.355501 | 23303.24 | 0.382322 | ${ }^{18269.29}$ | 0.410119 | ${ }^{223623.31}$ |
| 0.298199 | 27754.45 | 0.328765 | 28155.82 | 0.368821 | 33125.83 | 0.394369 | 31085.66 | 0.28734 | 25143.44 | 0.311791 | 19488.88 | 0.337234 | 22042.27 | 0.363674 | 23303.24 | 0.391111 | 18269.29 | 0.419547 | ${ }^{223623.31}$ |
| 0.309925 | 27754.45 | 0.33618 | 34050.29 | ${ }^{0.36896}$ | 33200.38 | 0.403264 | 31085.66 | 0.293805 | 25143.44 | 0.318798 | 19488.88 | 0.344812 | 22042.27 | 0.371846 | 23303.24 | ${ }^{0.3999}$ | 18703.69 | 0.428975 | ${ }^{223623231}$ |
| 0.311651 | 27754.45 | 0.343596 | 34812.38 | 0.377098 | 33200.38 | ${ }^{0.41216}$ | ${ }^{31085.66}$ | 0.300262 | 25143.44 | 0.325805 | 19488888 | ${ }^{0.35239}$ | 22042.27 | 0.380019 | 23303.24 | 0.408689 | 19115.63 | 0.438803 | ${ }^{223622.31}$ |
| 0.318378 | 27754.45 | 0.351012 | 34812.38 | 0.385237 | 33200.38 | 0.421055 | 31085.66 | 0.306719 | 25143.44 | 0.332811 | 19488.88 | 0.359968 | 22042.27 | 0.388191 | 23303.24 | 0.417478 | 19115.63 | 0.447831 | ${ }^{223623.31}$ |
| 0.325104 | 27754.45 | 0.358828 | 34812.38 | 0.393376 | 33200.38 | 0.429951 | 31949.53 | 0.313176 | 25143.44 | 0.339818 | 19488.88 | 0.367547 | 22042.27 | 0.396364 | 23303.24 | 0.426267 | 19115.63 | 0.457259 | 28716.89 |
| 0.33183 | 28830.66 | 0.365843 | 34812.38 | 0.401515 | 33200.38 | 0.438846 | 34561.99 | 0.319634 | 25143.44 | 0.346824 | 19488.88 | 0.375125 | 22042.27 | 0.404536 | 23303.24 | 0.435056 | 20301.64 | 0.466887 | ${ }^{34801.26}$ |
| 0.38857 | 32315.05 | 0.373259 | 34812.38 | 0.409654 | 33200.38 | 0.44742 | 37975.76 | 0.326091 | 25143.44 | 0.353831 | 19488.88 | 0.382703 | 22042.27 | 0.412708 | 23303.24 | 0.443845 | 25929.95 | 0.47615 | ${ }^{370288.25}$ |
| 0.345283 | 32315.05 | 0.388675 | 34812.38 | 0.417793 | 38855.72 | 0.45637 | 45216.3 | 0.332548 | 25143.44 | 0.360837 | 19488.88 | 0.390282 | 22042.27 | 0.420881 | 23303.24 | 0.452634 | 29360.69 | 0.48543 | ${ }^{38909.85}$ |
| 0.352009 | 32315.05 | 0.388891 | 34812.38 | 0.425931 | 43889.87 | 0.465533 | 55433.18 | 0.339005 | 25143.44 | 0.367844 | 19488.88 | 0.39786 | 22042.27 | 0.429053 | 27985.99 | 0.461423 | 35823 | 0.494971 | ${ }^{40331.52}$ |
| 0.358735 | 32315.05 | 0.395506 | 34812.38 | 0.43407 | 50018.16 | 0.474428 | 61349.02 | 0.345463 | 25143.44 | 0.37785 | 19488.88 | 0.405438 | 22042.27 | 0.437226 | 31112.28 | 0.470212 | 39834.45 | 0.504399 | 40331.52 |
| 0.365462 | 32315.05 | 0.402922 | 34812.38 | 0.442209 | 50018.16 | 0.483324 | ${ }^{61549.3}$ | 0.35192 | 25143.44 | 0.381857 | 19488888 | 0.413016 | 26070.63 | 0.445398 | ${ }^{33233.88}$ | 0.479001 | 40736.92 | 0.513827 | ${ }^{40331.52}$ |
| 0.372188 | 32315.05 | 0.410338 | 43913.34 | 0.450348 | 50018.16 | 0.492219 | 61599.3 | 0.358377 | 25143.44 | 0.388864 | 19626.84 | 0.420595 | 28300.98 | 0.453571 | 33233.88 | ${ }^{0.48779}$ | 40736.92 | 0.523255 | 40331.52 |
| 0.378914 | 32315.05 | 0.417753 | 47752.75 | 0.458887 | 50018.16 | 0.501115 | 62518.61 | 0.364834 | 25143.44 | 0.39587 | 23290.49 | 0.428173 | 33191.58 | 0.461743 | 33444.56 | 0.496579 | 40736.92 | 0.532683 | 40331.52 |
| 0.385641 | 32599.66 | 0.425169 | 47752.75 | 0.466625 | 50018.16 | 0.51001 | 62518.61 | 0.371292 | 25143.44 | 0.402877 | 26025.97 | 0.435751 | 36301.74 | 0.46999 | 33844.56 | 0.505368 | 40736.92 | 0.542111 | ${ }^{40331.52}$ |
| 0.392367 | 40679.86 | 0.432585 | 51593.12 | 0.474764 | ${ }^{521295.8}$ | 0.518906 | ${ }^{62518.61}$ | 0.377749 | 25143.44 | 0.409883 | 30988.15 | ${ }^{0.44333}$ | 37165.44 | 0.478088 | 33344.56 | 0.514157 | 40736.92 | 0.551539 | ${ }^{40331.52}$ |
| ${ }^{0.3999933}$ | ${ }^{45272.688}$ | 0.400001 | ${ }_{5}^{5538279}$ | 0.482903 | ${ }_{5}^{56123.05}$ | 0.527801 | ${ }^{62518.61}$ | ${ }^{0.384206}$ | 25143.44 | ${ }^{0.41689}$ | ${ }^{34156.16}$ | 0.455908 | 37165.44 | ${ }^{0.48626}$ | ${ }^{33444.56}$ | 0.522946 | ${ }^{40736.92}$ | ${ }^{0.556967}$ | ${ }^{43500.84}$ |
| 0.405819 | 45272.68 | 0.447416 | 55622.98 | 0.491042 | 56123.05 | 0.536697 | 62518.61 | 0.390663 | 28023.93 | 0.423896 | 35070.47 | 0.458886 | 37165.44 | 0.494433 | 33444.56 | 0.531735 | 40736.92 | 0.575395 | ${ }^{48143.4}$ |
| - 0.412546 | 45805.24 | 0.454832 | 55622.98 | 0.499181 | 56123.05 | 0.545593 | ${ }_{6}^{62518.61}$ | 0.397121 0.403578 | 31590.29 <br> 3214983 <br> 1 | 0.430903 | 35070.47 | 0.466064 | 37165.44 | 0.502605 | 33444.56 | 0.540524 | 44584.07 | 0.579823 |  |
| - 0.419272 | 5 | 0.4.4629864 | 55622.98 <br> 5562.98 | 0.50732 0.515458 | 5 | ${ }_{0}^{0.55433888}$ | ${ }_{62518.61}^{62518}$ | 0.403578 0.410035 | ${ }_{32762.22}^{3219.83}$ | - $\begin{aligned} & 0.437999 \\ & 0.44916\end{aligned}$ | ${ }^{35070.47}$ 3507.47 | 0.473643 0.481221 | 37165.44 <br> 371654 | ${ }_{0}^{0.510778} \begin{gathered}\text { 0.51895 }\end{gathered}$ | ${ }_{\substack{354823.7 \\ 42888}}$ | ${ }^{0.545313} \mathbf{0 . 5 8 1 0 2}$ | 46640.69 48608.16 | ${ }_{0}^{0.5598857}$ | 48143.4 48143.4 |
| 0.432725 | 52466.85 | 0.477079 | 55622.98 | 0.523597 | 56123.05 | 0.57279 | 62518.61 | 0.416492 | 32762.22 | 0.451922 | 35070.47 | 0.488799 | 37165.44 | 0.527123 | 42586.88 | 0.566891 | 49310.28 | 0.688107 | 59825.12 |
| 0.439451 | 52466.85 | 0.484495 | 55622.98 | 0.531736 | 56123.05 | 0.581175 | 62518.61 | 0.42295 | 32776.22 | 0.458929 | 35070.47 | 0.496378 | 37165.44 | 0.535295 | ${ }^{42586.88}$ | 0.57568 | 49310.28 | 0.617535 | ${ }^{61448.62}$ |
| 0.446177 | S2466.85 | ${ }_{0}^{0.4991911}$ | 55622.98 5522298 | ${ }^{0.539875}$ | 56123.05 <br> 57720.45 | 0.59007 0.559966 | ${ }^{62518.61}$ | ${ }^{0.429407}$ | 32762.22 <br> 3276222 | ${ }^{0.465936}$ | 35070.47 | ${ }^{0.503956}$ | 37680.23 3896232 | ${ }_{0}^{0.543468}$ | 44727.49 <br> 582624 | 0.584469 | 49310.28 4931028 | ${ }^{0.0626633}$ | ${ }^{61448.62}$ |
| $\begin{gathered} 0.452903 \\ 0.45963 \end{gathered}$ | $\begin{aligned} & 52466.85 \\ & 52466.85 \end{aligned}$ | 0.499327 0.506742 | $\begin{aligned} & 55622.98 \\ & 5562.98 \\ & 5 \end{aligned}$ | 0.548014 0.556152 | 57420.45 64870.4 | 0.598966 0.607861 | ${ }_{85152.56}^{7471.76}$ | 0.0.435864 | ${ }_{32762.22}^{32722.22}$ | 0.477942 0.479949 | ${ }^{40064.61} 4$ | 0.515354 0.519112 | 38962.32 40883.97 | $\begin{array}{r} 0.55164 \\ 0.559812 \end{array}$ | $\begin{aligned} & 58262.42 \\ & 6084.01 \\ & 608 \end{aligned}$ | $\begin{aligned} & 0.593258 \\ & 0.602047 \end{aligned}$ | 49310.28 49310.28 | 0.636391 0.645819 | 61448.62 61448.62 |
| 0.466356 | 52466.85 | 0.514158 | 55622.98 | 0.564291 | 64870.4 | 0.616757 | 87659.57 | 0.48879 | 33182.32 | 0.486955 | 5218.37 | 0.526691 | 40883.97 | 0.567985 | 60845.01 | 0.610837 | 56678.67 | 0.65547 | 61488.62 |
| 0.473082 | 55653.06 | 0.521574 | 55622.98 | ${ }^{0.57243}$ | ${ }^{64870.4}$ | 0.625652 | 89098.57 | 0.4555236 | 40191.34 | 0.493962 | 52323.58 | 0.534269 | 443664.52 | 0.576157 | 68845.01 | 0.619626 | 56678.67 | 0.664675 | ${ }^{61448.62}$ |
| 0.479889 | 56327.11 | ${ }^{0.52899}$ | 58736.74 | 0.580569 | 72659.87 | 0.634458 | 99686.48 | 0.461693 | 48711.71 | 0.500968 | 52323.58 | 0.541847 | 47307.48 | ${ }^{0.58433}$ | 60845.01 | 0.628415 | 56678.67 | 0.674103 | ${ }^{61448.62}$ |
| 0.486535 | 56327.11 | 0.536405 | 66265.03 | 0.588708 | 76499.19 | 0.643343 | 103262 | 0.46815 | 53031.69 | 0.507975 | 52323.58 | 0.594926 | 60619.78 | 0.592502 | 60845.01 | 0.637204 | 56678.67 | 0.683532 | ${ }^{61448.62}$ |
| 0.493261 | 56377.11 | 0.543821 | 77551.23 | 0.596846 | 82511.67 | 0.652339 | 103262 | 0.477468 | 53352.91 | 0.514981 | 52323.58 | 0.557004 | 67123.35 | 0.600675 | 68845.01 | 0.645993 | 56678.67 | ${ }^{0.69296}$ | ${ }^{61448.62}$ |
| 0.499987 | 60775.34 | 0.551237 | 88137.16 | 0.604885 | 95198.11 | 0.661234 | 106844.2 | 0.481065 | 53409.14 | 0.522988 | 52323.58 <br> 5232388 | ${ }^{0.564582}$ | ${ }_{6}^{67123.35}$ | ${ }^{0.068847}$ | 60845.01 <br> 6085501 | ${ }^{0.654782}$ | 5667.67 56787 | ${ }^{0.702388}$ | ${ }^{61448.62}$ |
| ${ }^{0.505714}$ | ${ }_{60775534}$ | ${ }_{0}^{0.558653}$ | 94640.41 | 0.613124 | 99100.82 | ${ }^{0.67713}$ | ${ }_{1}^{113015.1}$ | ${ }^{0.4875322}$ | ${ }^{601006.01}$ | 0.528995 | 52323.58 <br> 523238 | ${ }^{0.57216}$ | ${ }_{6}^{67123.35}$ | ${ }^{0.61702}$ |  | 0.663571 | ${ }_{56677.67}^{56787}$ | ${ }^{0.711816}$ |  |
| 0.51344 | 60775.34 | 0.566068 | 94640.41 | 0.621263 | 99100.82 | 0.679926 | 113015.1 | 0.493979 | 61486.79 | 0.536001 | 52323.58 | 0.579739 | 67123.35 | 0.625192 | 60845.01 | 0.67236 | ${ }_{56678.67}$ | 0.721244 | ${ }^{65299.02}$ |
| ${ }^{0.520166}$ | ${ }^{6293997}$ | 0.573884 | 94640.41 | 0.629402 | 107684.9 | 0.687921 | 124782 | 0.500437 | 61915.18 | 0.543008 | 52323.58 | 0.587317 | 67123.35 | 0.633364 | ${ }^{624455.78}$ | 0.681149 | ${ }_{56678.67}$ | ${ }^{0.730672}$ | ${ }^{68413.86}$ |
| 0.526893 0.533619 | 69211.43 <br> 7757238 | 0.5809 <br> 0.58816 | 94640.41 | 0.0.637411 | 124621.1 <br> 128861 | ${ }^{0.696817}$ | 141040.5 <br> 1612077 |  | 61915.18 6191518 |  |  |  |  |  |  |  |  |  |  |
| 0.533619 0.540395 | 77572.38 87511.93 | 0.588316 0.59571 | 94640.41 <br> 99640.41 | 0.0.645679 | 128861 12861 | 0.705712 0.714608 | 161027.7 164999.4 | 0.513351 0.519888 | 61915.18 61915.18 | 0.5577211 0.564027 | 58210.61 <br> 58210.61 | 0.60274 0.610052 | 67123.35 67123.35 | $0.649709$ $0.657882$ | $\begin{array}{r} 69555.24 \\ 69555.24 \end{array}$ | $\begin{aligned} & 0.698727 \\ & 0.787516 \end{aligned}$ | $58894.42$ $69456.9$ | 0.749528 0.758956 | 84811.85 <br> 89374.52 |
| 0.540345 0.547071 | 8751193 <br> 10146.6 | ${ }^{0.595731}$ | 994404.4. ${ }^{\text {97432 }}$ | ${ }^{0.6563198}$ | ${ }_{1268861}^{12681}$ | 0.174488 0.723503 | 166499.4 166999 | 0.519888 0.526265 | 61915.18 619518 | ${ }_{0}^{0.567027}$ | 58210.61 <br> 58210.61 | ${ }^{0.610072}$ 0.6173 | 67123.35 | ${ }_{0}^{0.6568584}$ | ${ }_{695555.24}^{6955}$ | - $\begin{aligned} & 0.777516 \\ & 0.716305\end{aligned}$ | 77755.94 7 | -0.7588384 | 89377.52 <br> 8934.52 |
| 0.553798 | 107199.4 | 0.610563 | 108799.4 | 0.6770996 | 126861 | 0.732399 | 166499.4 | 0.532723 | 61915.18 | ${ }^{0.57804}$ | 58210.61 | 0.625208 | 67123.35 | 0.674227 | ${ }_{69555.24}$ | 0.725094 | ${ }^{87711.76}$ | 0.777812 | 89374.52 |
| 0.560524 0.5725 | 1088952.1 1089521 1 | ${ }_{0}^{0.617979}$ | 116197.2 126736 1 | - 0.678235 | ${ }^{126861}$ | ${ }_{0}^{0.771294}$ | 166499.4 <br> 166994 <br> 1 | 0.53918 0.555637 | ${ }_{6}^{61915.18} 6$ | 0.585047 <br> 0.59253 | 58210.61 <br> 5821061 | ${ }_{0}^{0.637878} 0$ | 67123.35 <br> 6772355 | ${ }_{0}^{0.688399}$ | ${ }_{6}^{69555524}$ | -0.738833 | ${ }_{\substack{87862.99 \\ 963473}}$ | ${ }^{0.78724}$ | 89374.52 89374.52 |
| ${ }^{0.556725}$ | 108952.1 | 0.625394 | 126738.6 | 0.686373 | 126881 | 0.75019 | 166499.4 | 0.545637 | 61915.18 | 0.592053 | 58220.61 | 0.640365 | 67123.35 | 0.690572 | 69555.24 | 0.742672 | 9634773 | 0.796668 | 89374.52 |
| 0.573977 | 108952.1 | 0.63281 | 147560 | 0.694512 | 129142.7 | 0.759885 | 166499.4 | 0.552094 | 61915.18 | ${ }^{0.59906}$ | 58220.61 | 0.647943 | 67123.35 | 0.698744 | 75172.51 | 0.751461 | 108937.5 |  |  |
| 0.588703 | 108852.1 | 0.640226 | 158414.1 <br> 15841 | 0.702651 |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 0.587429 0.594155 | 108952.1 108952.1 | ${ }^{0.647641}$ | 158414.1 15814.1 | 0.71079 0.718929 | 155056.2 155056.2 | 0.776876 0.78572 | 166499.4 166999.4 | 0.555099 | ${ }_{6}^{61915.18} 6$ | 0.613073 | 58210.61 583508 | 0.6631 0.67078 | 69899.2 79488.05 | $\begin{aligned} & 0.715089 \\ & 0.723261 \end{aligned}$ | 97615.24 109830 | $\begin{aligned} & 0.769039 \\ & 0.777898 \end{aligned}$ | 109836.2 <br> 109836.2 |  |  |
| 0.600882 | 108952.1 | 0.662473 | ${ }_{158844.1}$ | 0.727768 | ${ }_{1} 1555556.2$ | 0.794667 | 166499.4 | ${ }_{0} .5577923$ | 61915.18 | 0.627086 | ${ }_{59071.87}$ | 0.678256 | 79488.05 | 0.731434 | 113524 | 0.786617 | 109836.2 |  |  |
| ${ }^{0.067608}$ | 121524.1 <br> 135389 | 0.668889 | 1588414.1 | ${ }^{0.735206}$ | 155055.2 <br> 15556.2 | 0.797633 | 166499.4 |  |  |  |  |  |  |  |  |  |  |  |  |
| 0.614334 0.621061 | 136389.6 138038.8 | ${ }_{\substack{0.677304 \\ 0.68472}}^{0}$ | 158414.1 158414.1 | ${ }^{0.774345}$ | 155055.2 155056.2 |  |  | 0.590838 0.597295 | ${ }_{6}^{61915.18} 6$ | 0.0.6481099 | 71592.4 73883.05 | 0.693413 <br> 0.700991 | 89230.65 98716.85 | $0.747779$ | $\begin{array}{r} 113524 \\ 122399.1 \end{array}$ | 0.7998 | 109836.2 |  |  |
| 0.627787 | 138838.8 | 0.692136 | 158414.1 | 0.759623 | ${ }_{1} 555056.2$ |  |  | ${ }^{0.603752}$ | 61915.18 | 0.655112 | 85337.46 | ${ }_{0} 0.70857$ | 100870.7 | 0.764124 | 122399.1 |  |  |  |  |
| 0.634513 0.64124 | 148821.6 15418.7 | ${ }_{0}^{0.699552} 0$ |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |
| 0.64124 0.647966 | ${ }_{1}^{155418.7} 1$ | 0.706967 0.71483 | 158414.1 158414.1 | 0.7759 0.78439 | 155055.2 155056.2 |  |  | 0.616667 0.623124 | 649901.05 73509.42 | 0.669125 0.676132 | ${ }_{\substack{96276.6 \\ 9923282}}$ | $\begin{aligned} & 0.723726 \\ & 0.731304 \end{aligned}$ | $\begin{array}{r} 114500.9 \\ 121536 \end{array}$ | 0.780469 <br> 0.788641 | $\begin{aligned} & 122399.1 \\ & 122399.1 \end{aligned}$ |  |  |  |  |
| 0.654692 | 155418.7 | 0.721799 | 158414.1 | 0.792178 | 155556.2 |  |  | 0.629581 | 84067.03 | 0.683139 | 99232.82 | 0.738883 | 121536 | 0.796813 | 122399.1 |  |  |  |  |
| 0.661418 0.668145 | 155418.7 154487 1 | ${ }_{\substack{0.729215 \\ 0.7363}}$ | 1588414.1 158414.1 | 0.79764 | 155056.2 |  |  | 0.636039 0.642996 | 89669.47 <br> 9159306 | - 0.699145 | 99332828 <br> 9923828 | $0.746461$ | 121536 <br> 129337 |  |  |  |  |  |  |
| 0.668145 0.677871 | 155418.7 15418.7 | 0.73663 <br> 0.74046 | 158414.1 <br> 158414.1 |  |  |  |  | 0.642496 0.648953 | ${ }^{9119930.06}$ | 0 | 920322828.5 1002 | 0.754039 <br> 0.761618 | 129133.7 132346.8 |  |  |  |  |  |  |
| 0.681597 | 155418.7 | 0.751462 | 158814.1 |  |  |  |  | 0.65541 | 97415.31 | 0.711165 | 107835.4 | 0.769196 | 133346.8 |  |  |  |  |  |  |
| 0.688324 0.69505 | 155418.7 15448.7 | 0.758878 0.76623 | 1588414.1 158414.1 |  |  |  |  | 0.661868 0.668325 | ${ }_{1}^{103748989} 1$ | 0.718171 0.725178 | 119936.5 <br> 122255.2 | $0.776774$ $0.784352$ | 137391.1 |  |  |  |  |  |  |
| 0.69505 0.701776 | 155418.7 15418.7 | 0.766293 0.773799 | 158414.1 <br> 158414.1 |  |  |  |  | 0.668325 0.674782 | ${ }^{107592.1}$ 10752.1 | ${ }^{0.725178}$ | 122255.2 130757.9 |  | 140939.1 140939.1 |  |  |  |  |  |  |
| 0.788502 | 155418.7 | 0.781125 | 158414.1 |  |  |  |  | 0.681239 | 107592.1 | 0.739191 | 130757.9 | 0.799509 | 140939.1 |  |  |  |  |  |  |
| 0.77529 0.721955 | $\begin{array}{r}155418.7 \\ 162525 \\ \\ \hline\end{array}$ | 0.788541 0.795956 | 1588414.1 158414.1 1 |  |  |  |  | 0.687697 0.694154 | 107592.1 1142117 | 0.746198 | 130757.9 1307579 |  |  |  |  |  |  |  |  |
| 0.721955 0.728881 | 162525 168687.3 | ${ }^{0.79595668}$ | 158414.1 158414.1 |  |  |  |  | 0.69454 0.700611 | 112411.7 <br> 126961.5 | ${ }^{0.753204}$ | $\begin{aligned} & 130757.9 \\ & 130757.9 \end{aligned}$ |  |  |  |  |  |  |  |  |
| 0.735408 | 175731.3 |  |  |  |  |  |  | 0.707068 | 142556.4 | 0.767217 | 130757.9 |  |  |  |  |  |  |  |  |
| 0.742134 | ${ }_{2221491.5}^{2035.7}$ |  |  |  |  |  |  | 0.713526 0.719983 0 | 146314.7 <br> 146314.7 | 0.774224 | 143894.2 <br> 150955 <br> 1 |  |  |  |  |  |  |  |  |
| 0.778886 0.755886 | 221491.5 <br> 222950.1 <br> 2 |  |  |  |  |  |  | 0.719983 <br> 0.72644 | ${ }_{14633414.7}^{14847}$ | ${ }^{0.7888237}$ | $\begin{aligned} & 150975.4 \\ & 150975.4 \end{aligned}$ |  |  |  |  |  |  |  |  |
| 0.762313 | ${ }^{222950.1}$ |  |  |  |  |  |  | 0.773887 | 146314.7 | 0.795243 | 150975.4 |  |  |  |  |  |  |  |  |
| 0.769393 0.775756 | 222950.1 222950.1 |  |  |  |  |  |  | ${ }^{0.7393555}$ | 146314.7 <br> 163314 | ${ }^{0.798747}$ | 150975.4 |  |  |  |  |  |  |  |  |
| 0.775765 0.782992 | 222950.1 <br> 222950.1 |  |  |  |  |  |  | $\begin{aligned} & 0.745812 \\ & 0.752269 \end{aligned}$ | $\begin{aligned} & 146314.7 \\ & 146314.7 \end{aligned}$ |  |  |  |  |  |  |  |  |  |  |
| - 0.7898248 | ${ }^{22229550.1}$ |  |  |  |  |  |  | $\begin{aligned} & 0.752269 \\ & 0.758726 \end{aligned}$ | 146314.7 146314.7 |  |  |  |  |  |  |  |  |  |  |
| 0.79594 |  |  |  |  |  |  |  | ${ }^{0.7655184}$ | 1468314.7 |  |  |  |  |  |  |  |  |  |  |
|  |  |  |  |  |  |  |  | 0.771641 | 146314.7 |  |  |  |  |  |  |  |  |  |  |
|  |  |  |  |  |  |  |  | 0.778098 | 146314.7 |  |  |  |  |  |  |  |  |  |  |
|  |  |  |  |  |  |  |  | 0.78455 | 16314, |  |  |  |  |  |  |  |  |  |  |
|  |  |  |  |  |  |  |  | 0.79747 | 146314.7 |  |  |  |  |  |  |  |  |  |  |

